



# *Overview of results from ALICE*

- Particle production and collective flow
- Correlations and Fluctuations
- Summary and Outlook

UJ 22.03.2021 (virtual ☹)

Jacek Otwinowski (IFJ PAN, Krakow)  
On behalf of the ALICE Collaboration

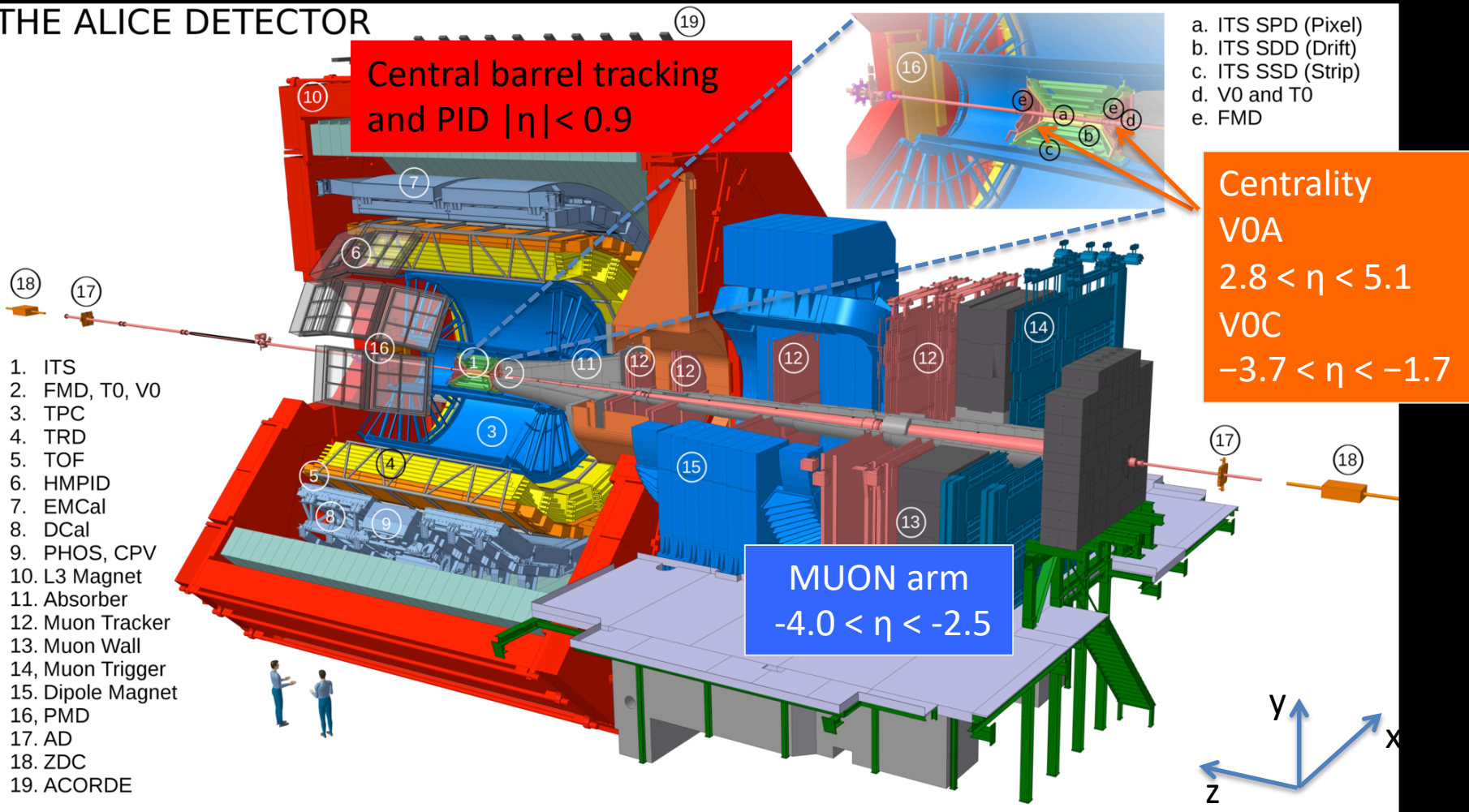


# A Large Ion Collider Experiment



- Excellent particle identification capabilities over a wide  $p_T$  range 0.1-20 GeV/c
- Good momentum resolution  $\sim 1-5\%$  for  $p_T = 0.1-50$  GeV/c

## THE ALICE DETECTOR



# ALICE at work since 2009

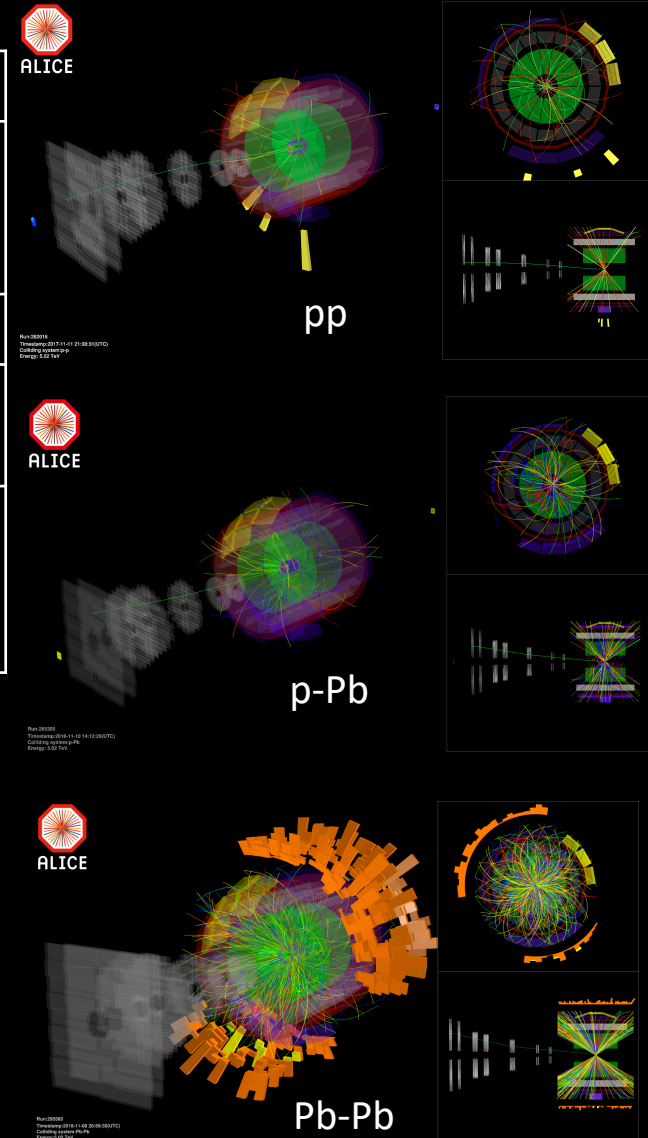


System	Year	$\sqrt{s_{NN}}$ (TeV)	$L_{int}$
Pb-Pb	2010-2011	2.76	$\sim 75 \mu\text{b}^{-1}$
	2015	5.02	$\sim 250 \mu\text{b}^{-1}$
	2018	5.02	$\sim 0.9 \text{nb}^{-1}$
Xe-Xe	2017	5.44	$\sim 0.3 \mu\text{b}^{-1}$
p-Pb	2013	5.02	$\sim 15 \text{nb}^{-1}$
	2016	5.02, 8.16	$\sim 3 \text{nb}^{-1}, \sim 25 \text{nb}^{-1}$
pp	2009-2013	0.9, 2.76, 7, 8	$\sim 200 \mu\text{b}^{-1}, \sim 100 \mu\text{b}^{-1}, \sim 1.5 \text{pb}^{-1}, \sim 2.5 \text{pb}^{-1}$
	2015-2018	5.02, 13	$\sim 1.3 \text{pb}^{-1}, \sim 59 \text{pb}^{-1}$

- Energy and system dependence studies of particle production are possible
- Large statistics of pp, p-Pb and Pb-Pb collisions at the same  $\sqrt{s_{NN}}$

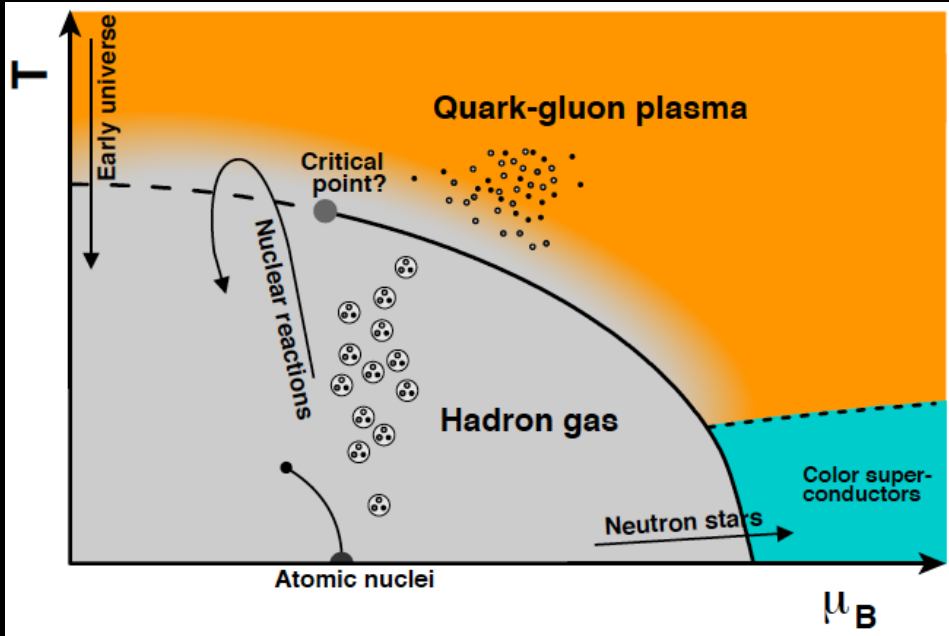
→ **Precise comparison studies**

**Study of the quark-gluon plasma**

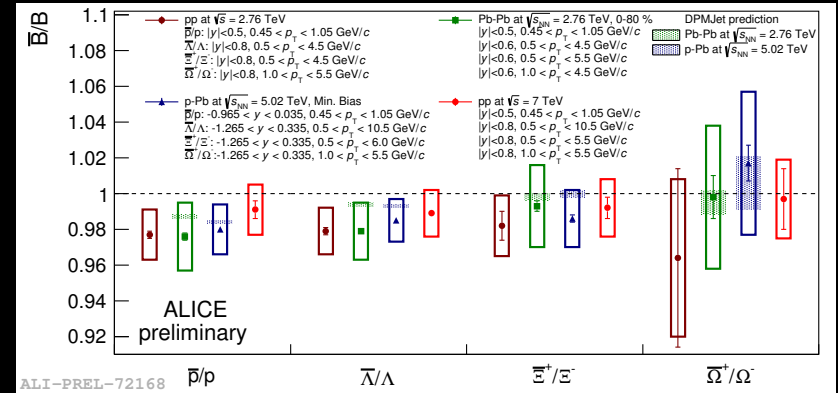


# ALICE case

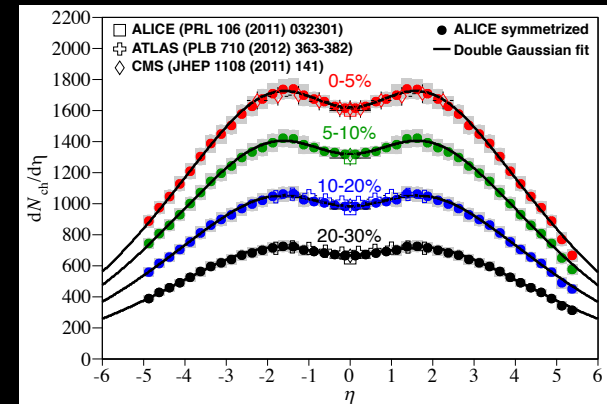
ALICE probes nuclear matter at  $\mu_B \sim 0$  and high temperature



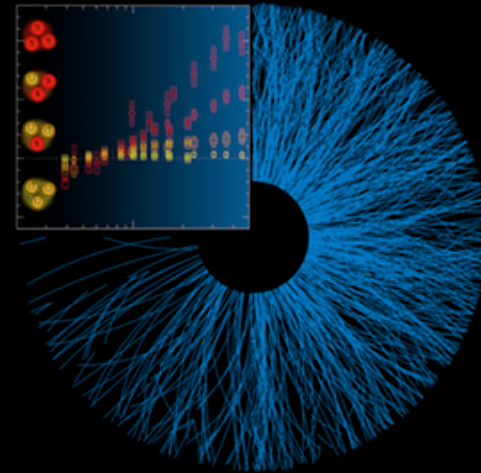
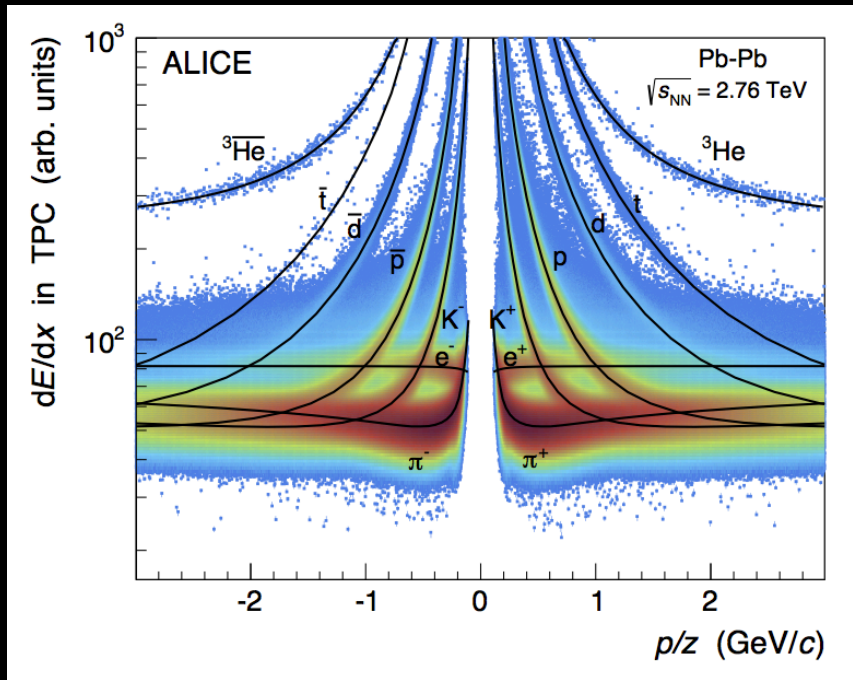
Lattice QCD: Hadron gas to quark-gluon plasma  
 ( $\mu_B = 0$ ,  $T_C = 156.5 \pm 1.5$  MeV,  
 $\epsilon_C = (0.42 \pm 0.06)$  GeV/fm<sup>3</sup>, smooth crossover)  
 A. Bazavov et al. Phys. Lett. B 795 (2019) 15



Phys. Lett. B 726 (2013) 610-622



Central collisions:  
 $\epsilon \sim 14$  GeV/fm<sup>3</sup>  $\gg \epsilon_C$

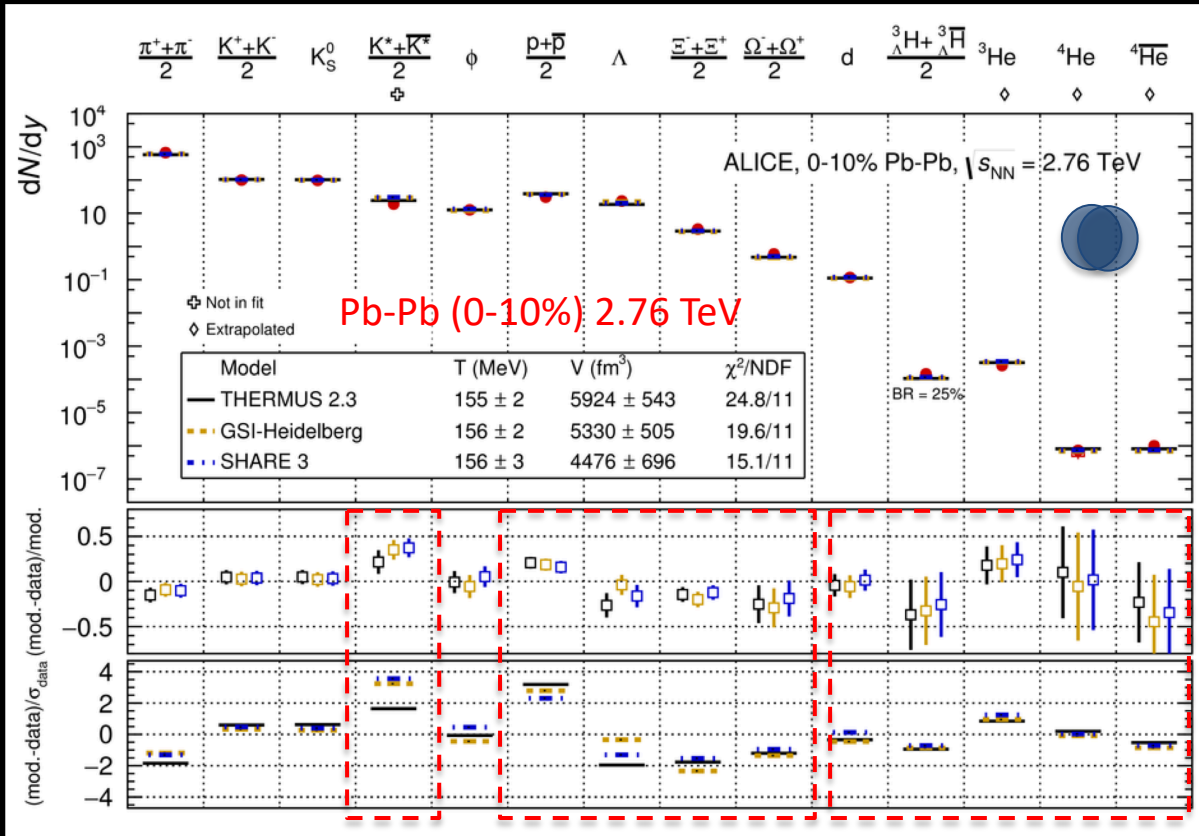


# PARTICLE PRODUCTION

# Hadron yields vs thermal models in Pb-Pb at $\sqrt{s_{NN}} = 2.76$ TeV



Nucl. Phys. A 971 (2018) 1



- Production of (most) hadrons well described at freezeout temperature  $T_{ch} \sim 156$  MeV
- $K^{*0}$  resonance (not included in the fit): production overestimated by thermal models
- Tension for protons and multi-strange baryons
- Light nuclei production described by thermal models (binding energy  $\ll T_{ch}$ )?

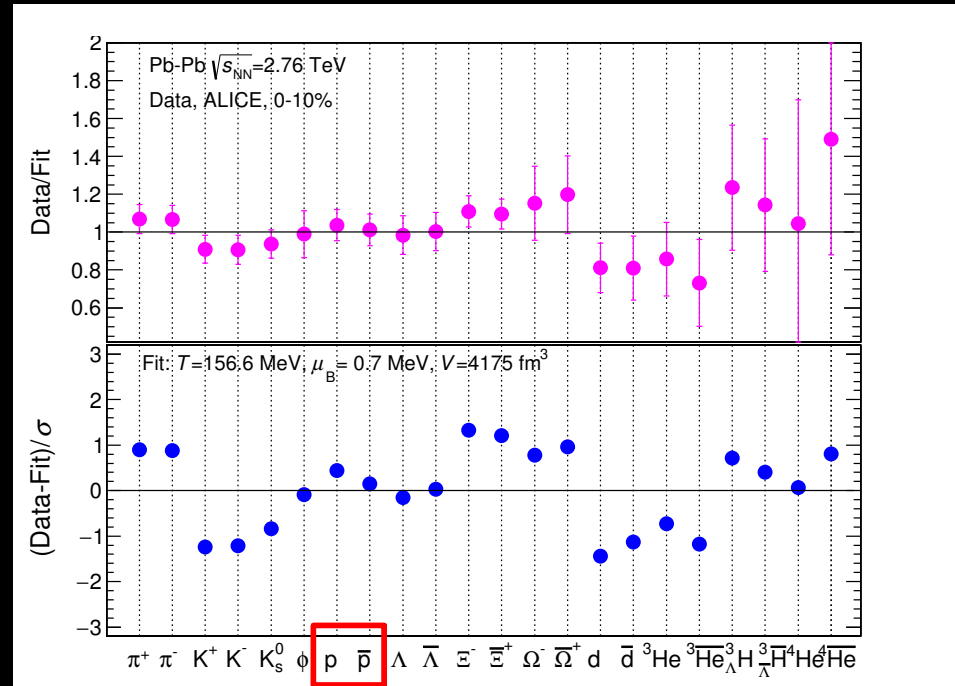
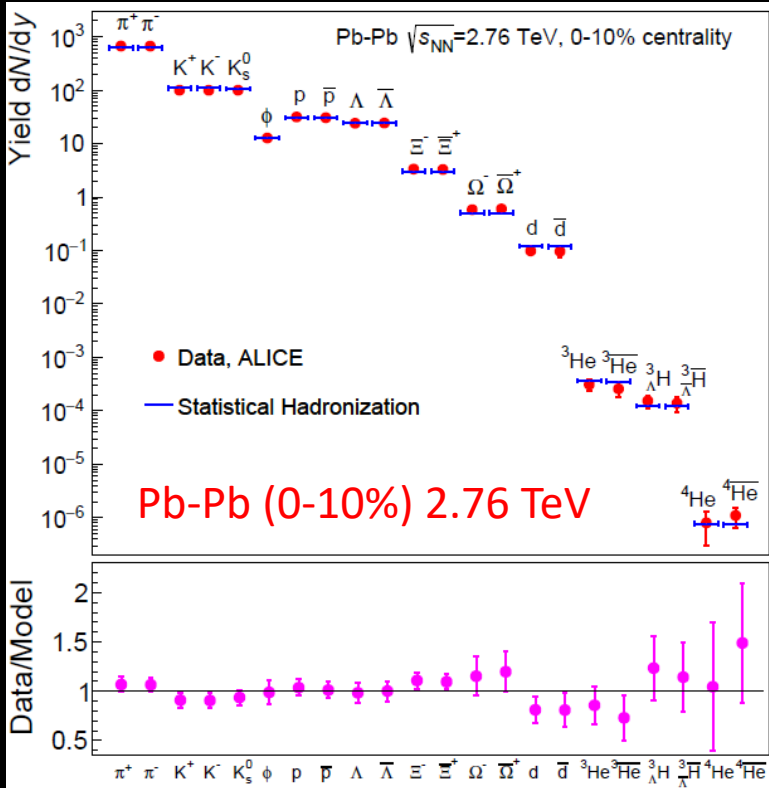
THERMUS: Wheaton et al., Comput. Phys. Commun, 180 (2009) 84

GSI-Heidelberg: Andronic et al., Phys. Lett. B 673 142

SHARE: Petran et al., Comp. Phys. Commun. 195 (2014) 2056

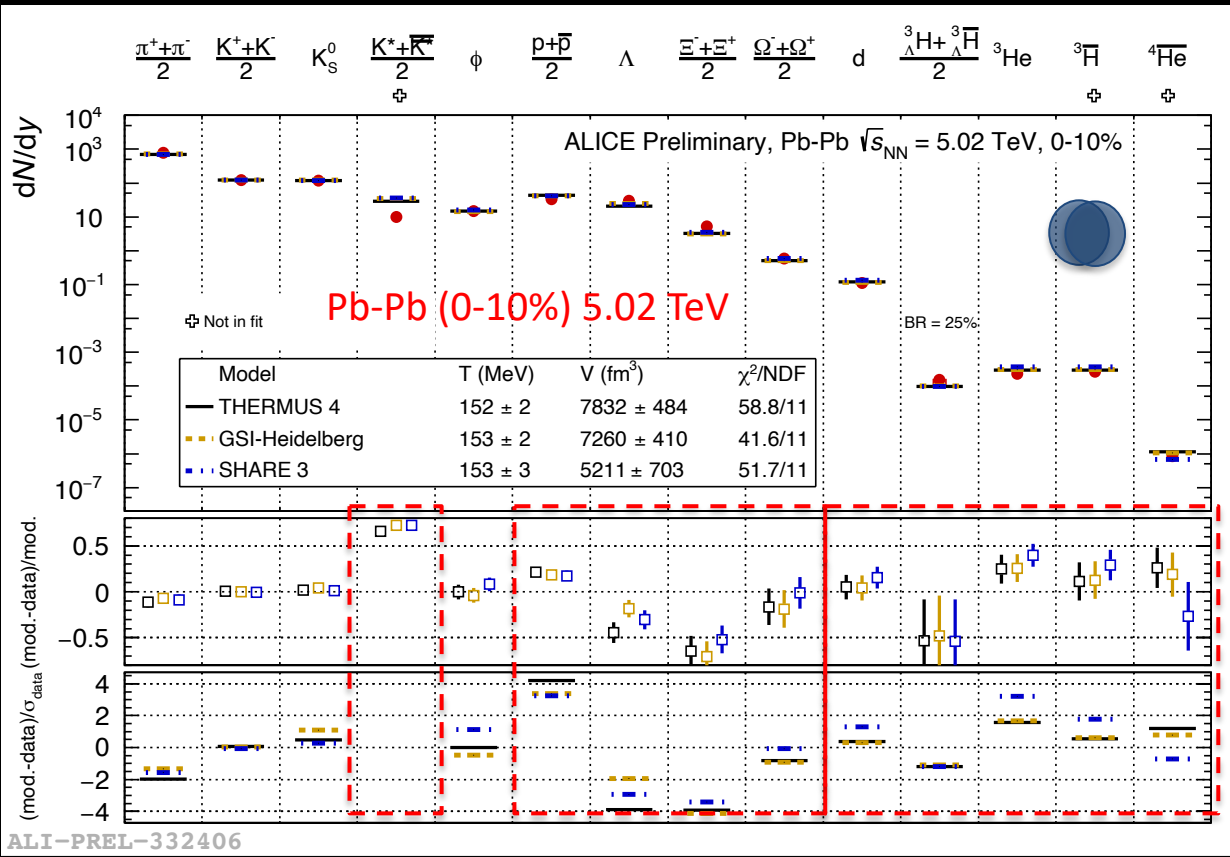
# Hadron yields vs improved thermal model in Pb-Pb at $\sqrt{s_{NN}} = 2.76$ TeV

A. Andronic et al. Phys. Lett. B792 (2019) 304



- Resonant and non-resonant pion-nucleon and multi-pion-nucleon interactions included in the model (S-matrix approach)
- Production of hadrons well described at freezeout temperature  $T_{ch} = 156.6 \pm 1.7$  MeV  $\rightarrow$  consistent with critical temperature  $T_c$  from the lattice QCD calculations

# Particle yields vs thermal models in Pb-Pb at $\sqrt{s}_{NN} = 5.02$ TeV

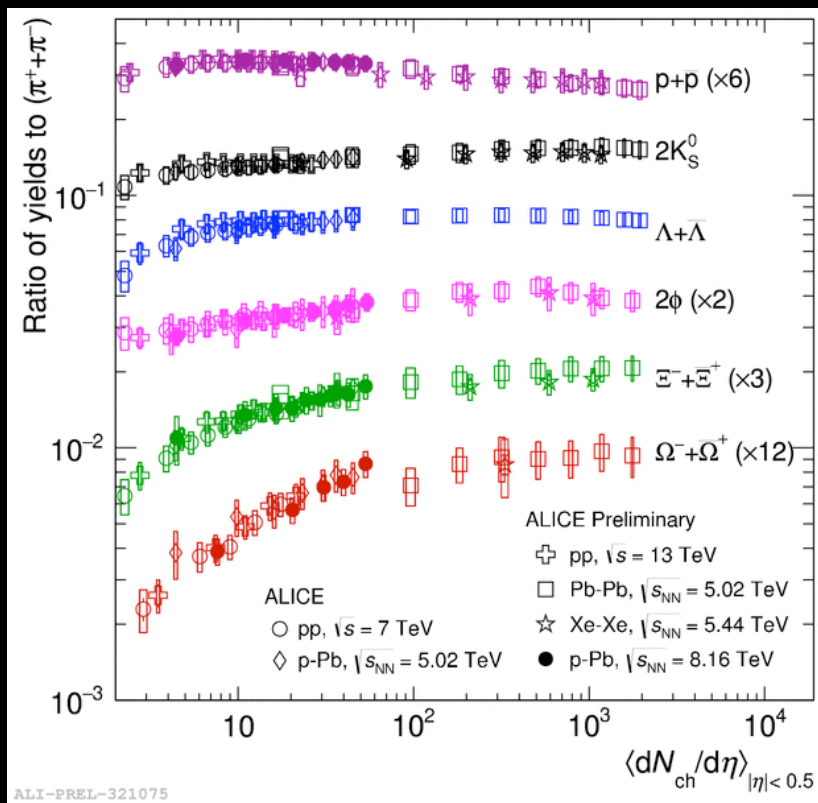


- Production of (most) hadrons well described at freeze-out temperature  $T_{ch} \sim 153$  MeV (different than at 2.76 TeV)
- $K^{*0}$  resonance (not included in the fit): production overestimated by thermal models
- Tension for protons and multi-strange baryons (new approach is currently studied)
- Light nuclei production described by thermal models (binding energy  $\ll T_{ch}$ )?

THERMUS: Wheaton et al., Comput. Phys. Commun, 180 (2009) 84  
 GSI-Heidelberg: Andronic et al., Phys. Lett. B 673 142  
 SHARE: Petran et al., Comp. Phys. Commun. 195 (2014) 2056



# Relative particle production in pp, p-Pb, Pb-Pb and Xe-Xe



- Smooth evolution from pp to Pb-Pb
- No significant energy and system dependence is observed at similar multiplicity
- Relative strangeness production increases with multiplicity and strangeness content in small systems

→ Hadron production is driven by the characteristics of final state

pp 7 TeV: [Nature Physics 13 \(2017\) 535](#)

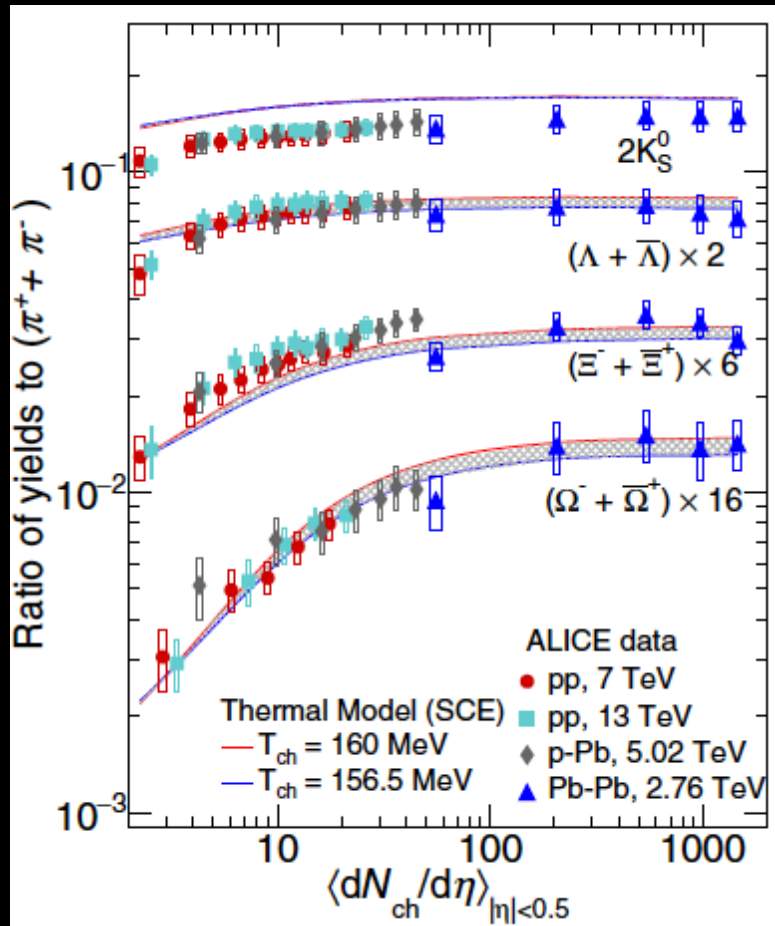
pp 7 and 13 TeV: [arXiv:2005.11120](#)

p-Pb 5.02 TeV: [Phys. Lett. B728 \(2014\) 25](#), [Phys. Lett. B758 \(2016\) 389](#)

Xe-Xe 5.44 TeV: [arXiv:2101.03100](#)

pp and Pb-Pb 5.02 TeV: [Phys. Rev. C 101, 044907 \(2020\)](#)

# Relative particle production in pp, p-Pb, Pb-Pb and Xe-Xe



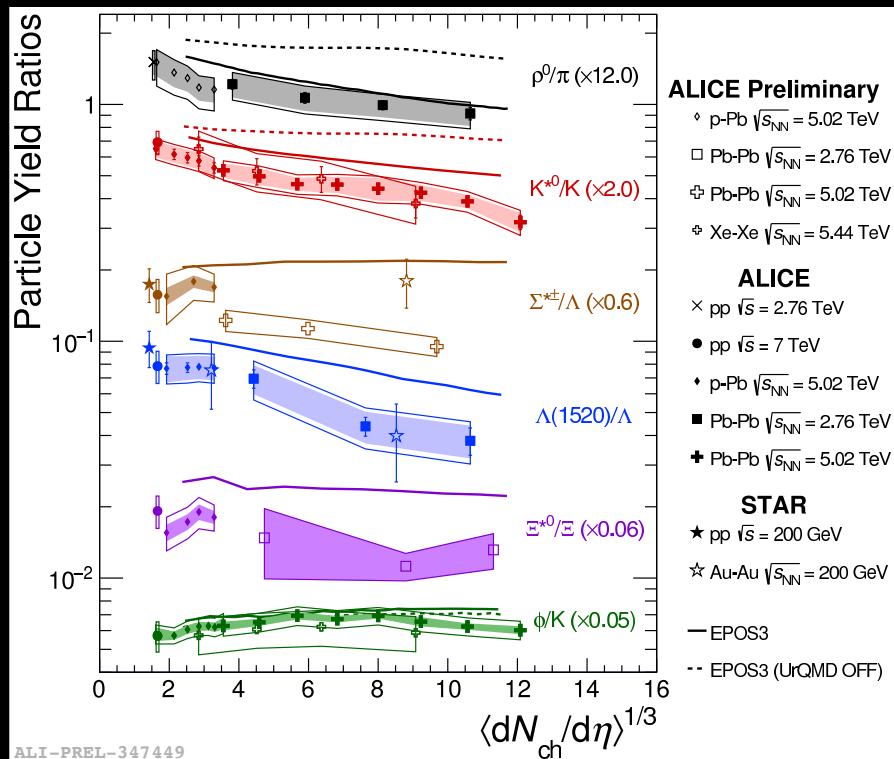
- Smooth evolution from pp to Pb-Pb
- No significant energy and system dependence is observed at similar multiplicity
- Relative strangeness production increases with multiplicity and strangeness content in small systems
- Thermal model (SCE) including interaction between hadrons (S-matrix approach) and exact strangeness conservation J. Cleymans et al. Phys. Rev. C103 014904 (2021)

pp 7 TeV and Pb-Pb 2.76 TeV: [Nature Physics 13 \(2017\) 535](#)  
pp 7 and 13 TeV: [arXiv:2005.11120](#)

# Relative resonance production in pp, p-Pb, Pb-Pb and Xe-Xe collisions



Resonance	$\rho^0$	$K^{*0}$	$\Sigma^{*\pm}$	$\Lambda(1520)$	$\Xi^{*0}$	$\phi$
Lifetime (fm/c)	1.3	4.16	5.5	12.6	21.7	46.2



- Relative suppression of  $\rho^0$ ,  $K^{*0}$ ,  $\Lambda(1520)$  with increasing multiplicity
- $\Sigma^*/\Lambda$ ,  $\Xi^*/\Xi$ ,  $\phi/K$  is independent of multiplicity
- Similar trend seen in all collision systems
- EPOS3 + UrQMD describes the trend of data

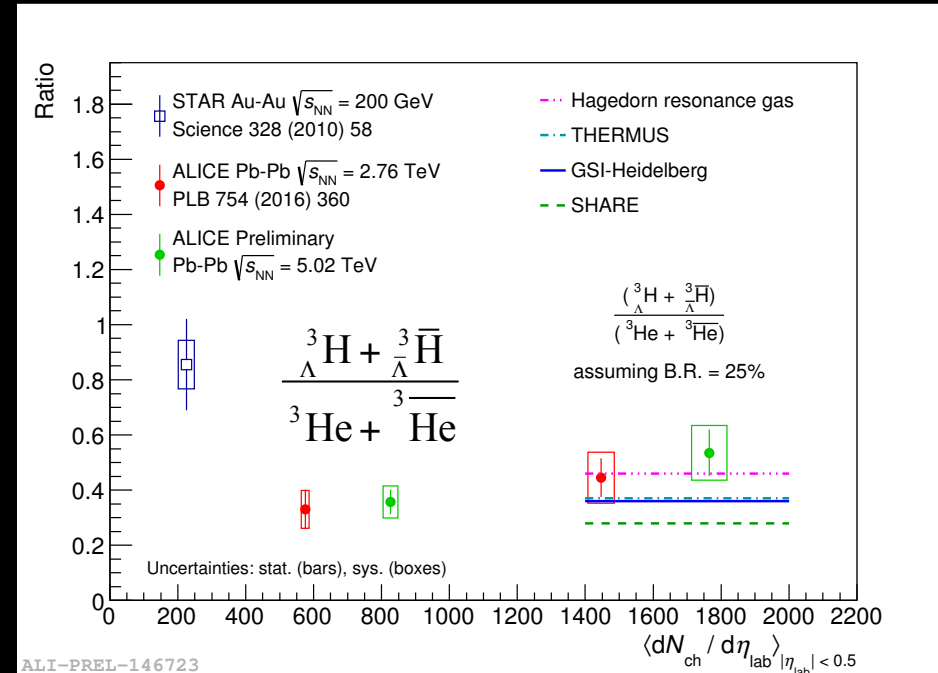
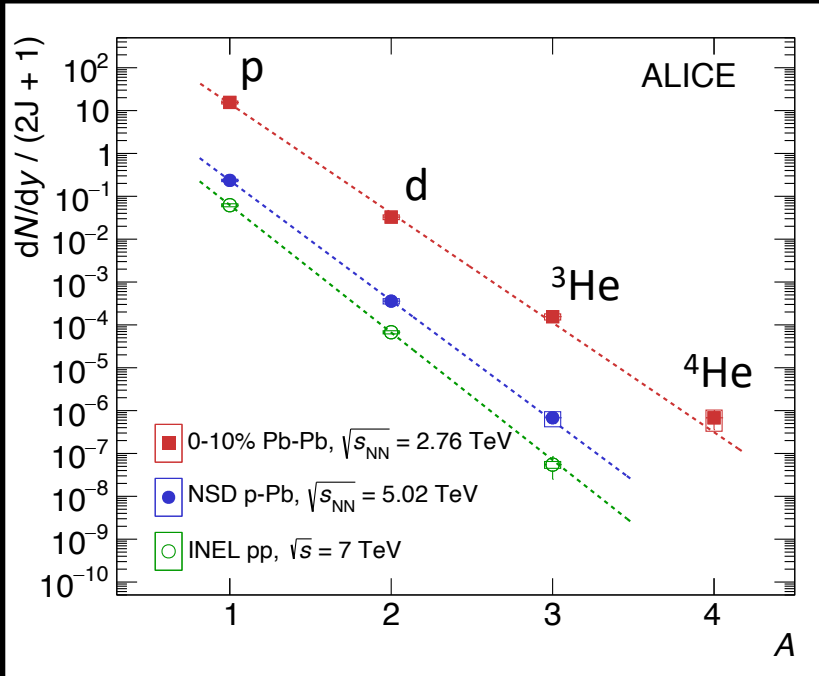
$\rho^0$ : Phys. Rev. C99, (2019) 064901  
 $K^{*0}$ : Phys. Rev. C95, (2017) 064606  
 $\Lambda(1520)$ : Phys. Rev. C99, (2019) 024905  
 STAR, Phys.Rev.C78 (2008) 044906  
 STAR, Phys. Rev. Lett. 97 (2006) 132301  
 $\phi$ : Phys. Rev. C91 (2015) 024609  
 $\Sigma^*$ ,  $\Xi^*$ : Eur. Phys. J. C 77 (2017) 389

→ Dominance of rescattering over (re)combination in hadronic phase

# Nuclei production in pp, p-Pb and Pb-Pb



p-Pb: *Phys. Lett. B* 800 (2019) 135043

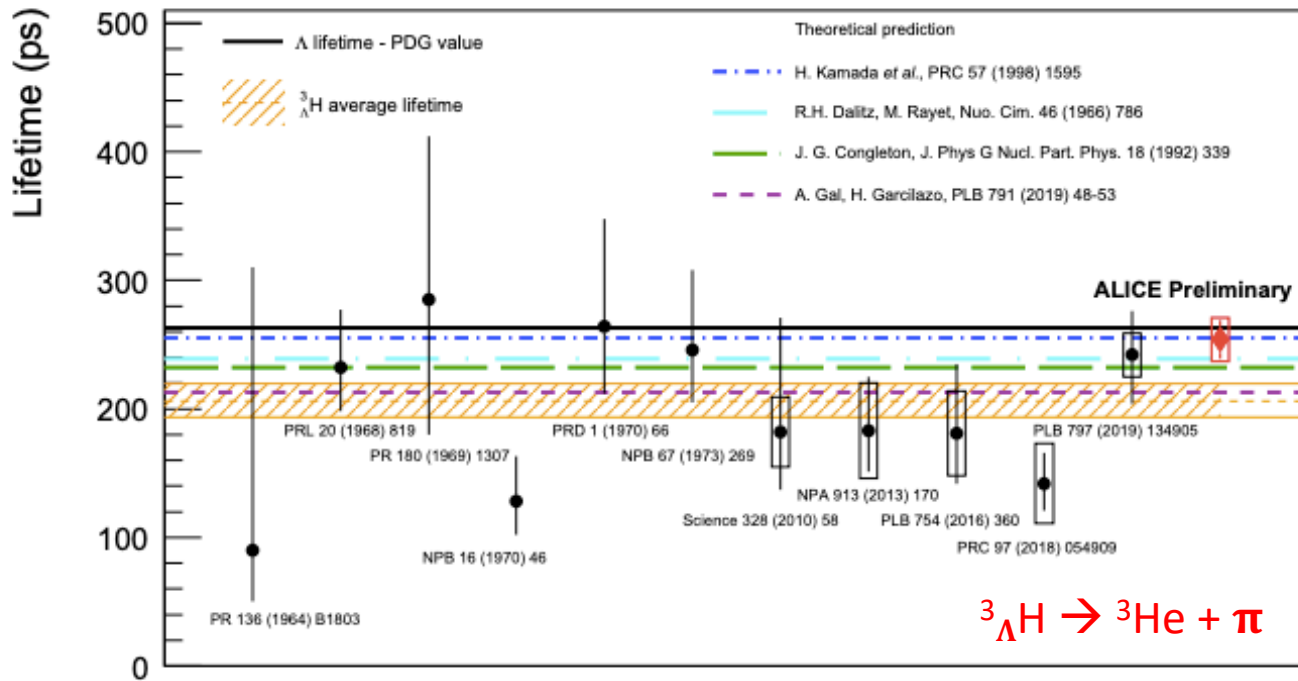


- Exponential decrease in nuclei rate described by thermal model
- Nuclei production rate decreases by factor of  $\sim 300$  (Pb-Pb),  $\sim 600$  (p-Pb) and  $\sim 1000$  (pp) for each additional nucleon
- ${}^3_{\Lambda}\text{H}$  production consistent with thermal model (binding energy  $\sim 0.13$  MeV  $\ll T_{ch}$ )  
→ Production mechanisms: thermal vs coalescence?

pp: *Phys. Rev. C* 97 (2018) 024615

Pb-Pb: *Nucl. Phys. A* 971 (2018) 1

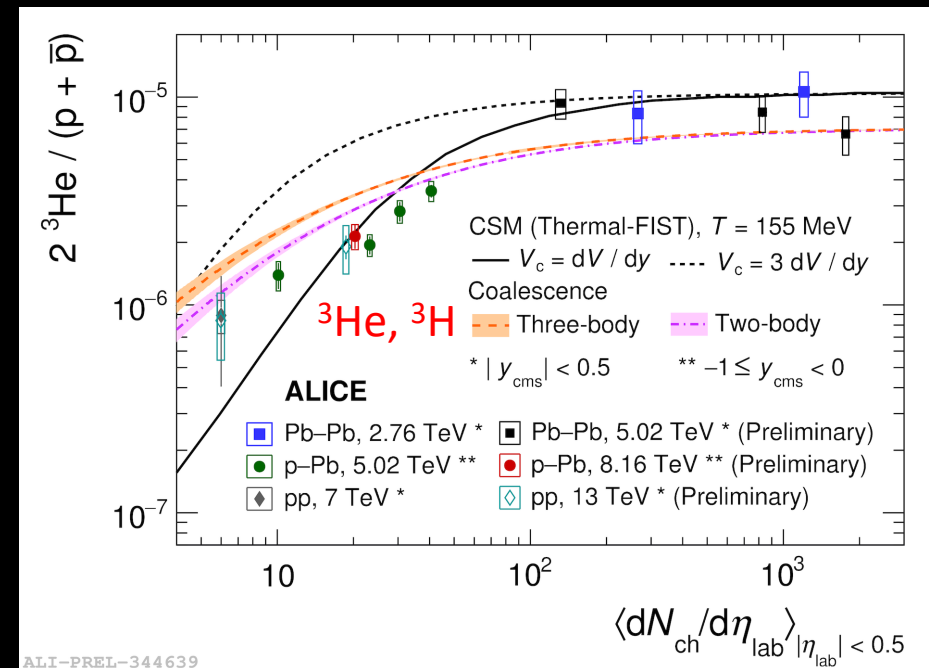
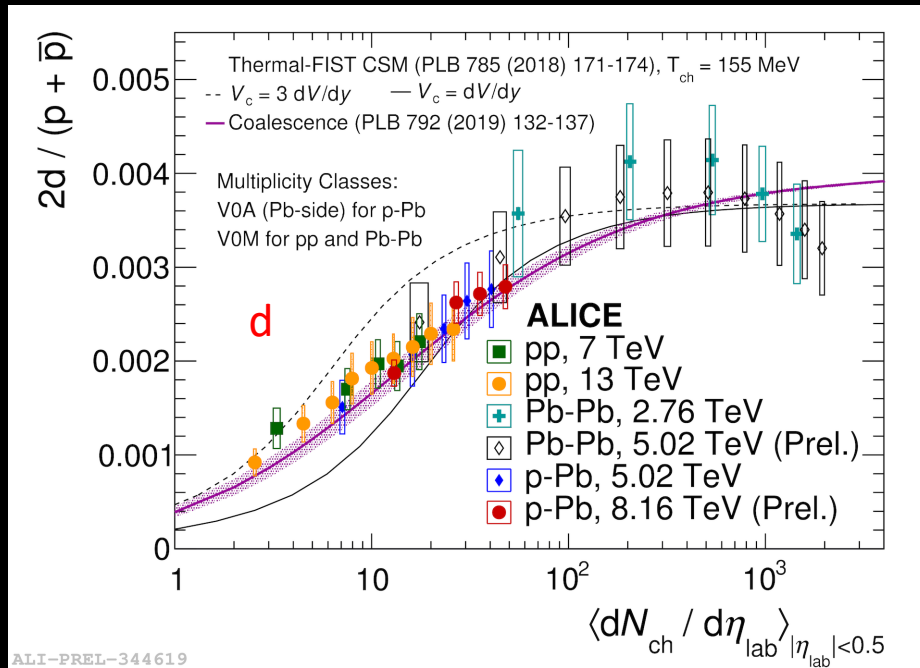
# Hypertriton lifetime



ALI-PREL-342050

- The most precise measurement of hypertriton lifetime
- Two body  ${}^3_{\Lambda}\text{H}$  decay channel has been used
- ${}^3_{\Lambda}\text{H}$  lifetime consistent with the free  $\Lambda$  lifetime (not included in the average)  
→ binding energy  $\sim 0.13$  MeV
- Large discrepancy between recent World results

# Formation of light nuclei



- Increase from pp to peripheral Pb-Pb
- No centrality dependence in high multiplicity Pb-Pb
- Data described qualitatively by coalescence and thermal models

→ Production mechanisms: thermal vs coalescence?

Thermal- FIST CSM: V. Vovchenko et al., Phys. Lett. B 785 (2018) 171

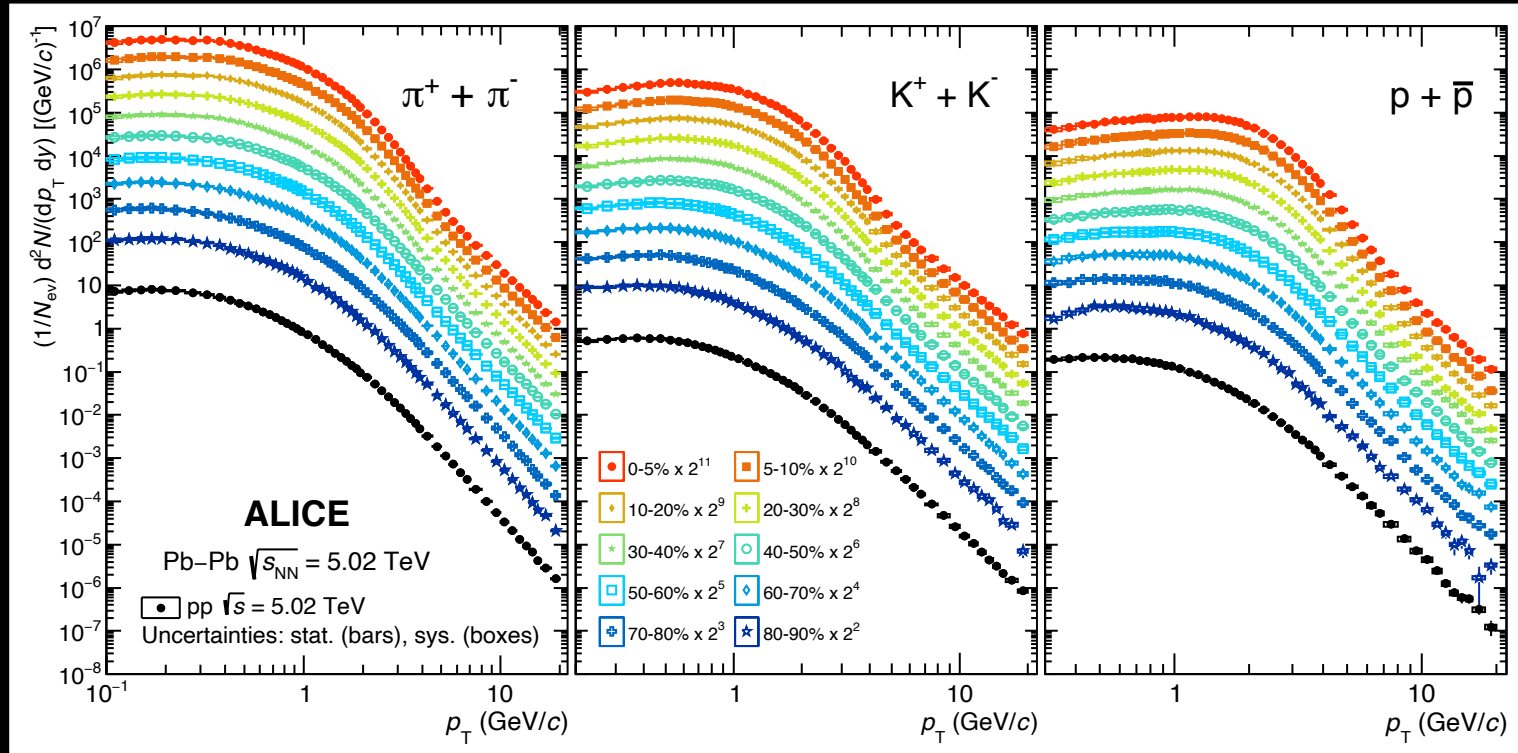
Coalescence: K.-J. Sun et al., Phys. Lett. B 792 (2019) 132

# SPECTRA

# Transverse momentum spectra of charged $\pi$ , K and p in Pb-Pb



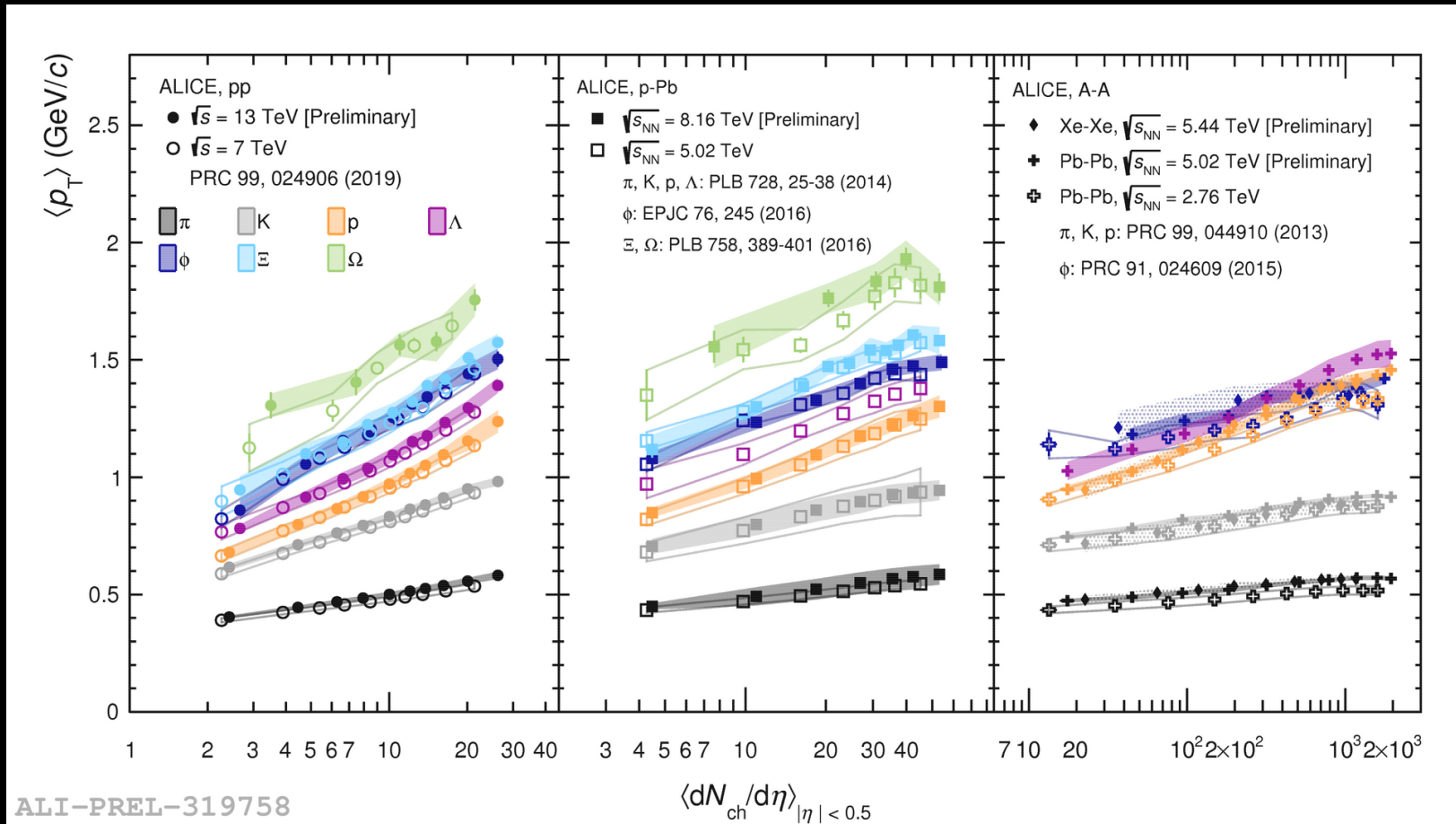
Phys. Rev. C 101, 044907 (2020)



- Particle identification with different analysis techniques: ITS, TPC, TOF, HMPID and topological identification of decaying charged kaons
- Mass dependent hardening of the spectra with increasing centrality  
→ Collective radial expansion



# $\langle p_T \rangle$ vs centrality of hadrons in pp, p-Pb, Pb-Pb and Xe-Xe



- $\langle p_T \rangle$  increases with increasing centrality and mass
- Larger increase in smaller systems  
 → Collective radial expansion

pp 13 TeV: [arXiv:2005.11120](https://arxiv.org/abs/2005.11120)  
 Xe-Xe 5.44 TeV: [arXiv:2101.03100](https://arxiv.org/abs/2101.03100)  
 Pb-Pb 5.02 TeV: [PRC 101, 044907 \(2020\)](https://arxiv.org/abs/2004.04907)

# Blast-Wave fit to hadron $p_T$ spectra

Phys. Rev. C 101, 044907 (2020)

$$E \frac{d^3 N}{d p^3} \propto \int_0^R m_T I_0 \left( \frac{p_T \sinh(\rho)}{T_{kin}} \right) K_1 \left( \frac{m_T \cosh(\rho)}{T_{kin}} \right) r dr$$

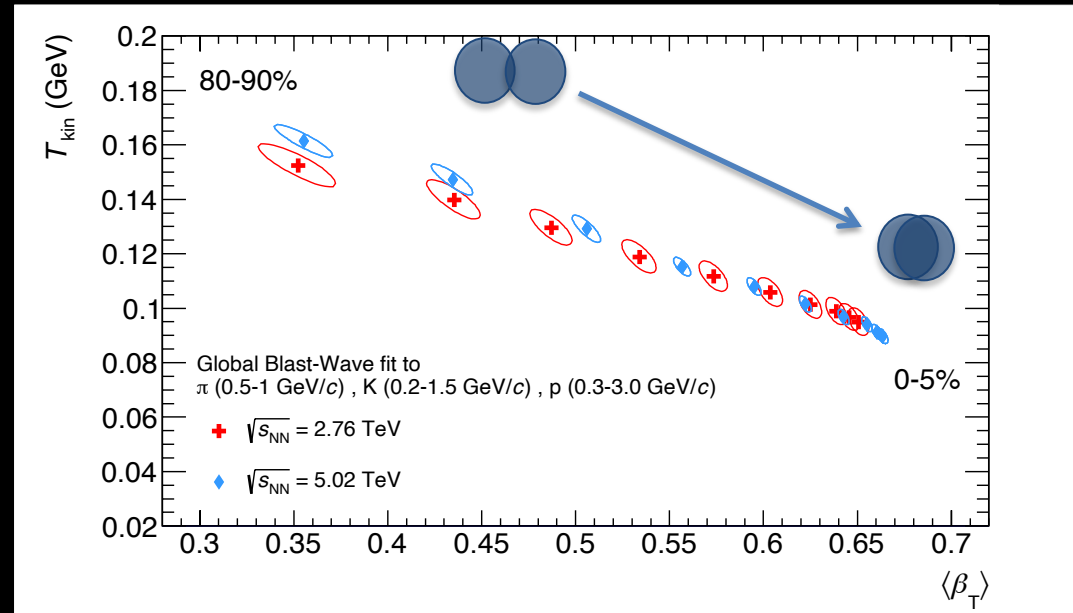
$$m_T = \sqrt{m^2 + p_T^2} \quad \rho = \tanh^{-1}(\beta_T) \quad \beta_T = \beta_s \left( \frac{r}{R} \right)^n$$

Schnedermann, Sollfrank and Heinz Phys. Rev. C 48, 2462

Simplified hydrodynamic model with 3 parameters:

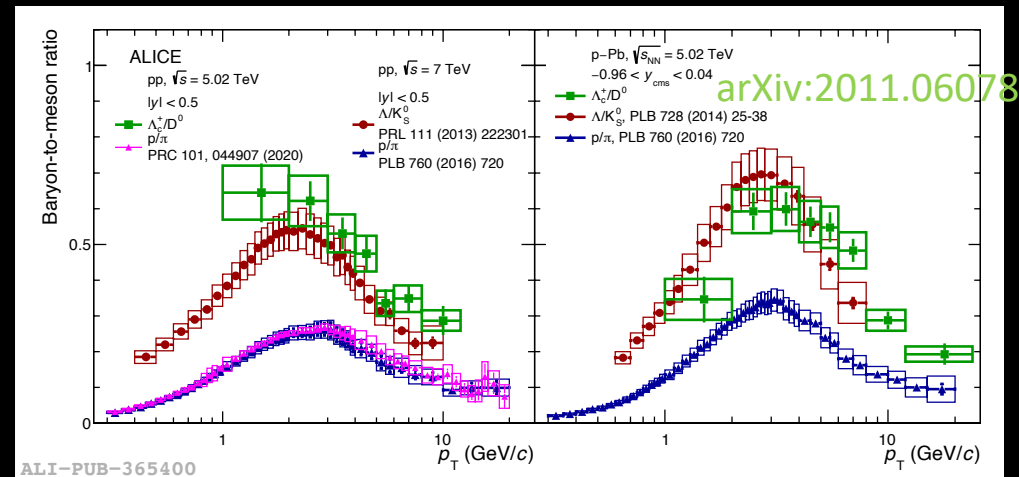
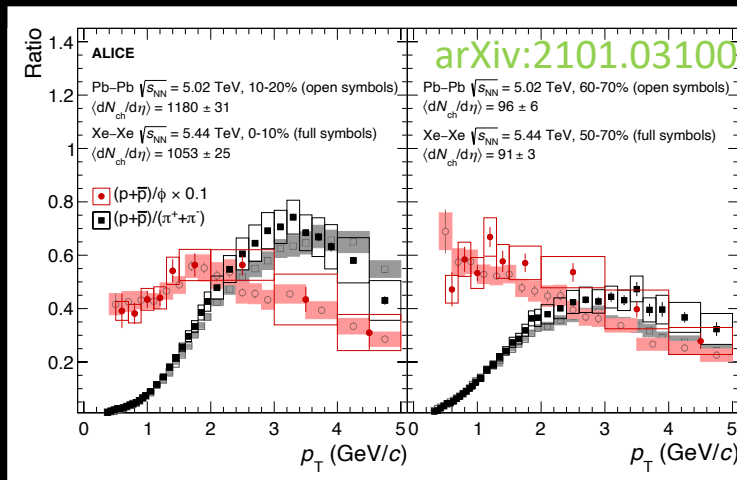
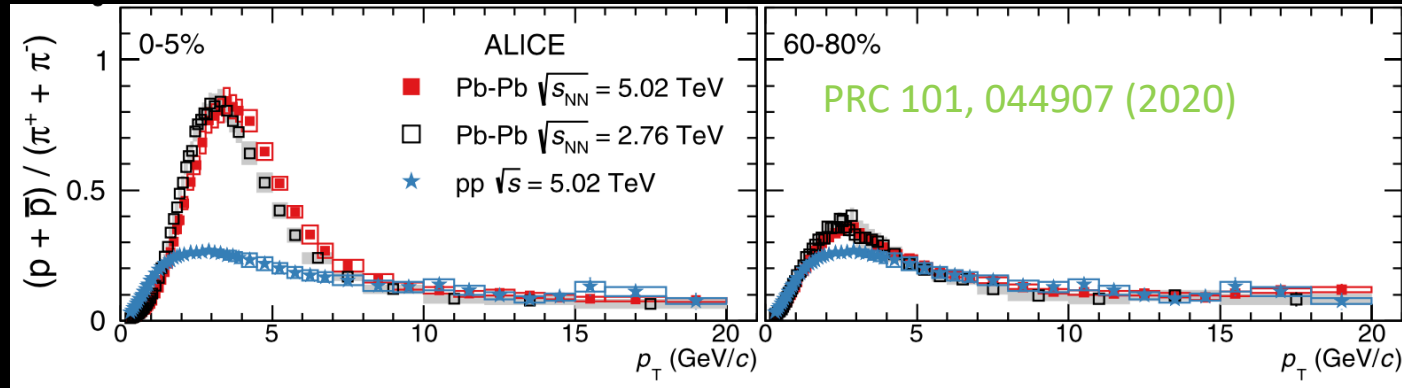
- $\beta_T$  - radial expansion velocity
- $T_{kin}$  - kinetic freeze-out temperature
- $n$  - velocity profile

Simultaneous fit to the  $\pi$ , K, p spectra



- $\langle \beta_T \rangle$  reaches  $\sim 0.65c$  in central Pb-Pb collisions
- $T_{kin}$  decreases with collision centrality

# Baryon-to-meson ratios

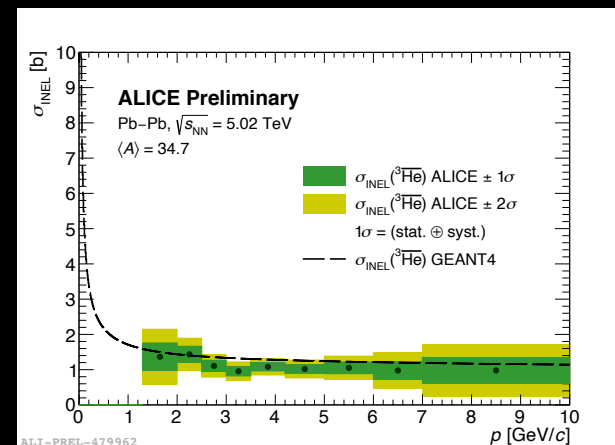
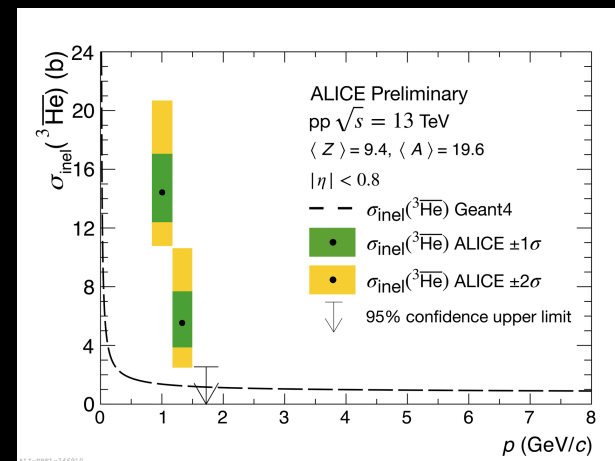
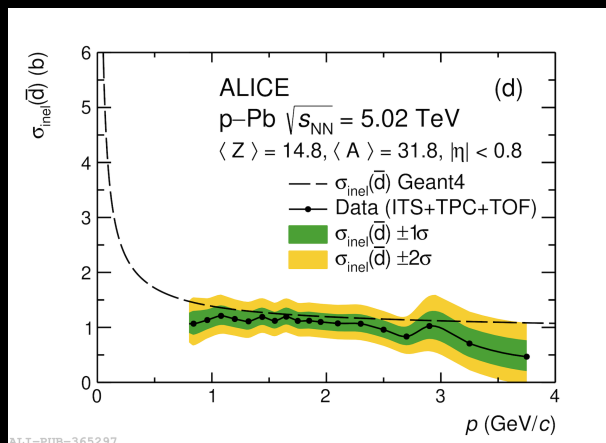
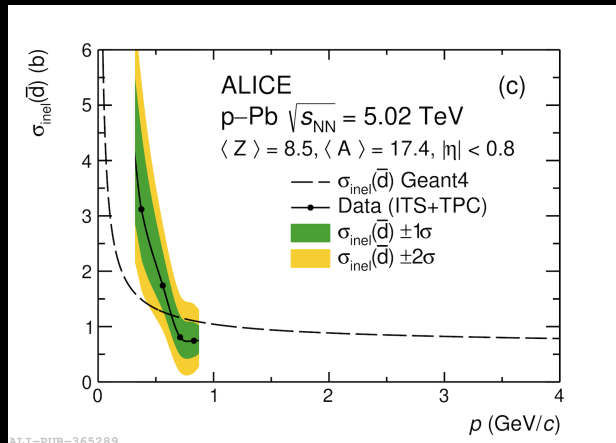


- Similar trends independent of collision system for the same multiplicity
- $\Lambda_C/D_0$  show similarities with those for light-flavor  $p/\pi$  and  $\Lambda/K^0_s$   
 → hint for the common production mechanism of light- and heavy-flavor baryons (coalescence vs. fragmentation)

# Antinuclei interaction cross sections



Phys. Rev. Lett. 125 (2020) 162001

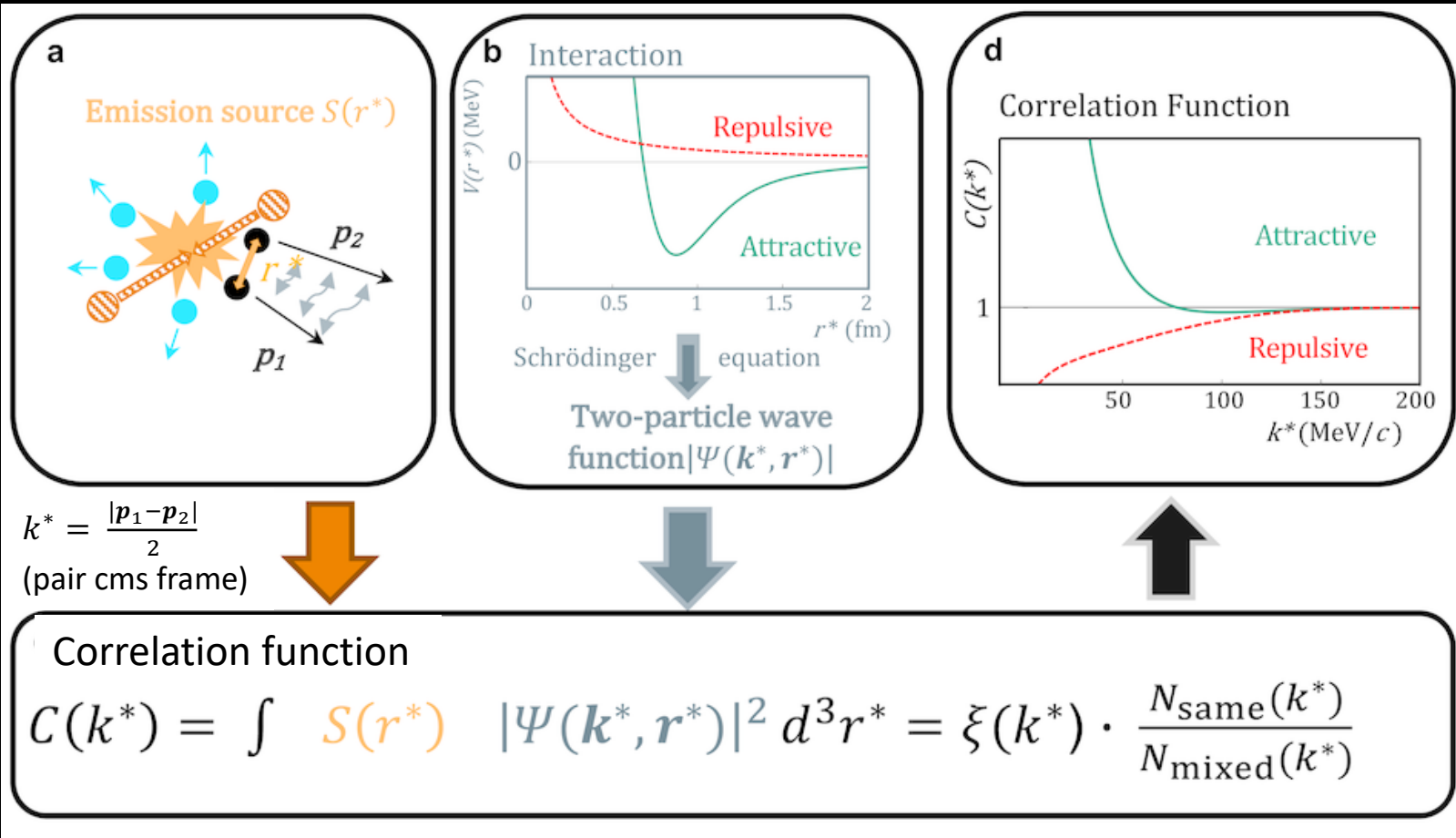


- Studies relevant for understating anti-nuclei absorption in the galactic medium (dark matter searches)
- **First measurement of anti-deuteron and anti-<sup>3</sup>He INEL cross section at low momentum**
- ALICE detector material used as target to study anti-nuclei absorption



# FEMTOSCOPY CORRELATIONS

# Femtoscscopy correlations to study stable and unstable hadron interactions

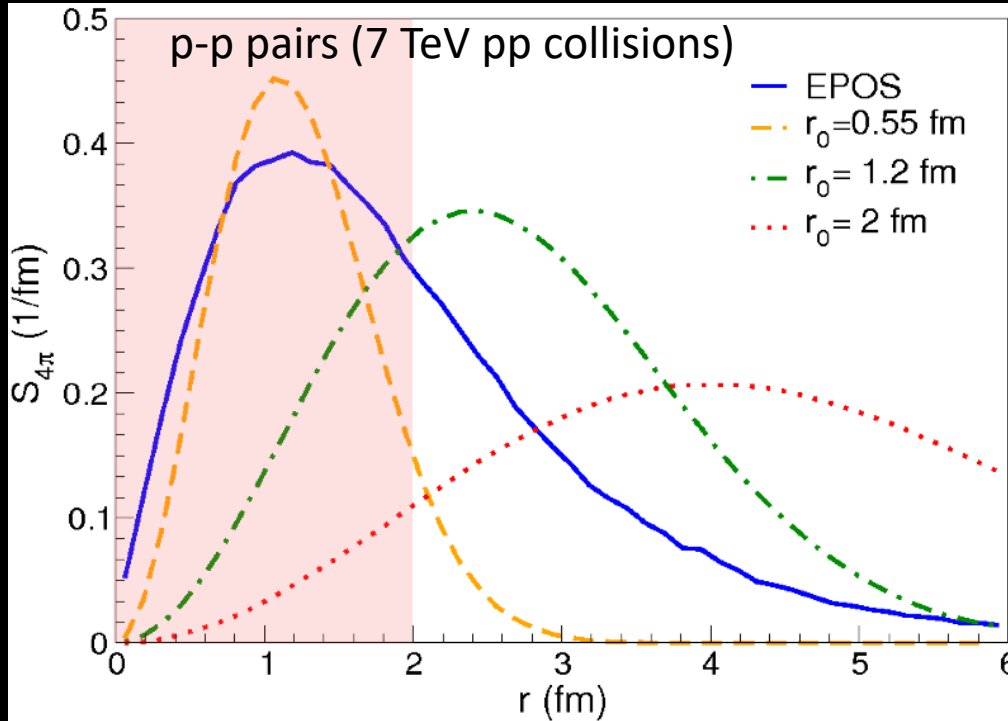


- Origin of correlations: quantum interference, resonances, conservation laws or final-state interactions
- Final-state interactions dominate at small  $k^*$

Pioneering by HADES Coll.  
Phys. Rev. C 94, 025201 (2016)

# Femtoscscopy in small systems

CATS, D. Mihaylov et al., Eur. Phys. J. C78 (2018) 394



Gaussian emission source:

$$S(r^*) = \frac{1}{(4\pi r_0^2)^{3/2}} \exp\left(-\frac{r^{*2}}{4r_0^2}\right)$$

Pair emission probability:

$$S_{4\pi}(r^*) = 4\pi r^2 S(r^*)$$

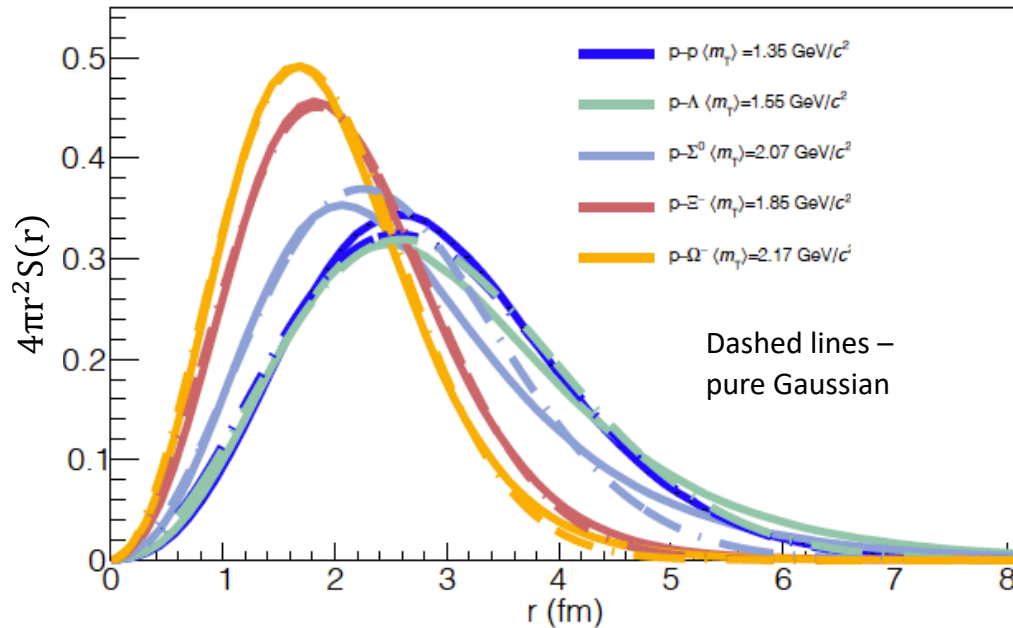
EPOS model predicts non-Gaussian emission source

T. Pierog et al. PRC 92 (2015) 034906

- A lot of pairs is emitted from the source at distance below  $\sim 2$  fm (typical range of strong interaction)
- Small particle emission source in pp and p-Pb collisions is essential to study the strong interaction!

# Common baryon source

Source using a Gaussian core plus resonances



ALICE PLB 811 (2020) 135849

Pair	$r_{\text{Core}}$ [fm]	$r_{\text{Eff}}$ [fm]
p-p	1.1	1.2
p- $\Lambda$	1.0	1.3
p- $\Sigma^0$	0.87	1.02
p- $\Xi^-$	0.93	1.02
p- $\Omega^-$	0.86	0.95

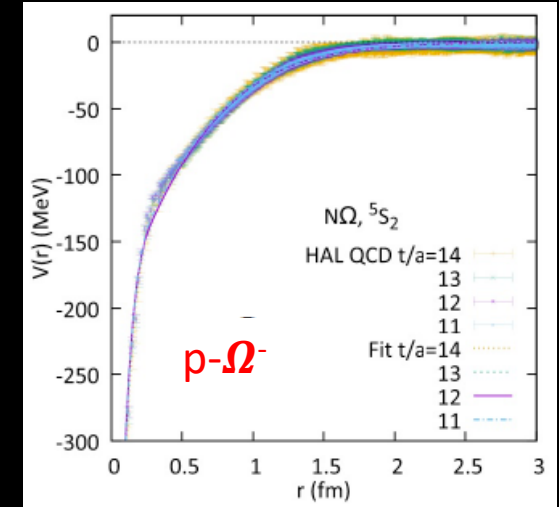
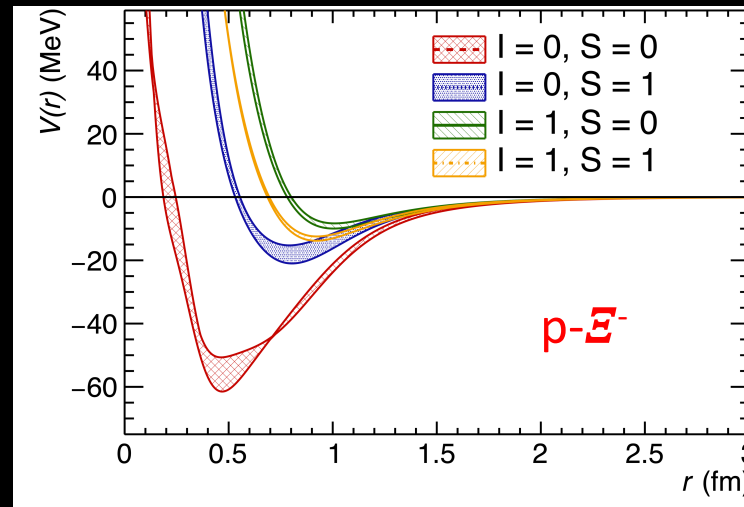
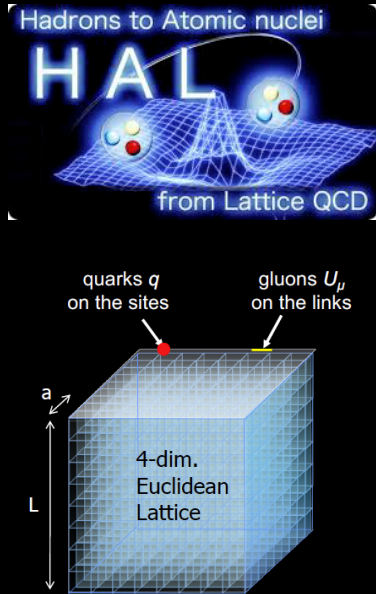
$$S(r^*) = G(r^*, r_{\text{core}}(m_T)) = \frac{1}{(4\pi r_{\text{core}}^2)^{3/2}} \exp\left(-\frac{r^{*2}}{4r_{\text{core}}^2}\right) \otimes E(r^*, M_{\text{res}}, \tau_{\text{res}}, p_{\text{res}})$$

- Emission source for heavier pairs using p-p correlation function plus resonances
- Gaussian source with  $r_{\text{eff}} = 1.02 \pm 0.05$  ( $0.95 \pm 0.06$ ) fm used for the p- $\Xi^-$  (p- $\Omega^-$ ) emission



# The strong interaction for $p\text{-}\Xi$ and $p\text{-}\Omega$ on lattice

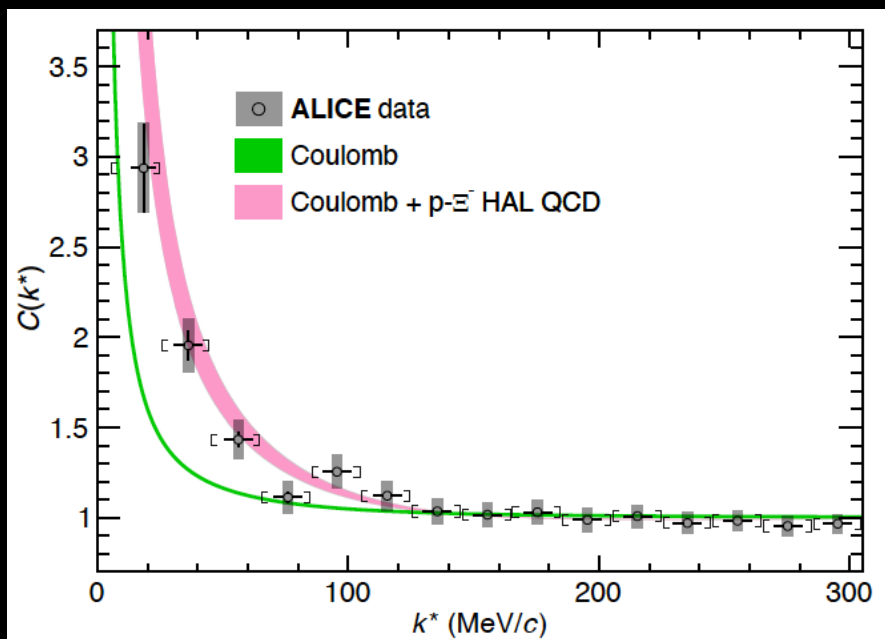
HAL QCD Coll., NPA 998 (2020) 121737, PLB 792 (2019) 284



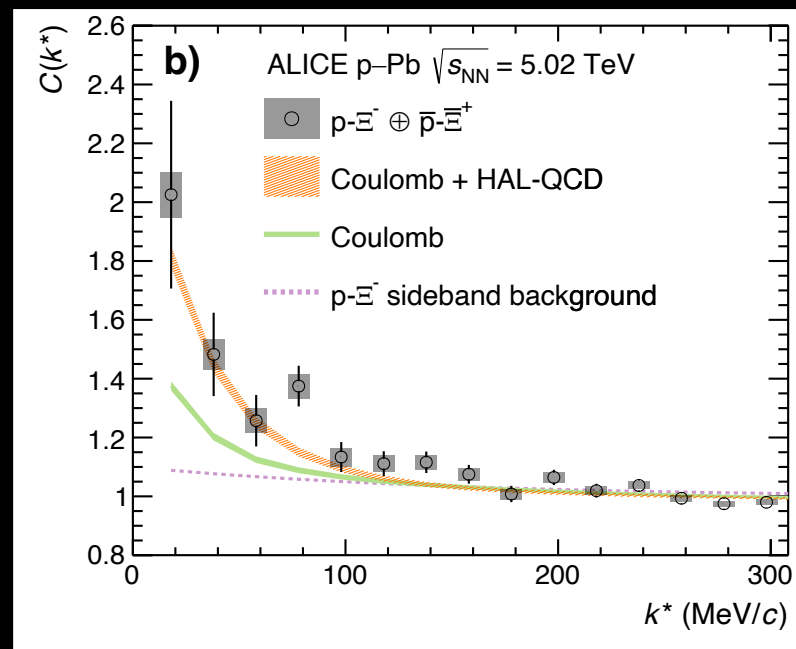
- $p\text{-}\Xi$  interaction in four channels: isospin ( $I = 0, 1$ ) and spin ( $S=0, 1$ )
  - **Attractive with repulsive core at small distances**
- $p\text{-}\Omega$  interaction in  ${}^5S_2$  ( $I=1/2, S=2$ ) channel
  - **Attractive in the whole range**
  - **After inclusion Coulomb interaction prediction of bound state with binding energy  $\sim 2.5$  MeV**
- $p\text{-}\Omega$  interaction in  ${}^3S_1$  channel does not include yet inelastic channels (e.g.  $p\Omega \rightarrow \Lambda\Xi$ )

# $p$ - $\Xi$ correlation function in pp and p-Pb

Nature 588 (2020) 232



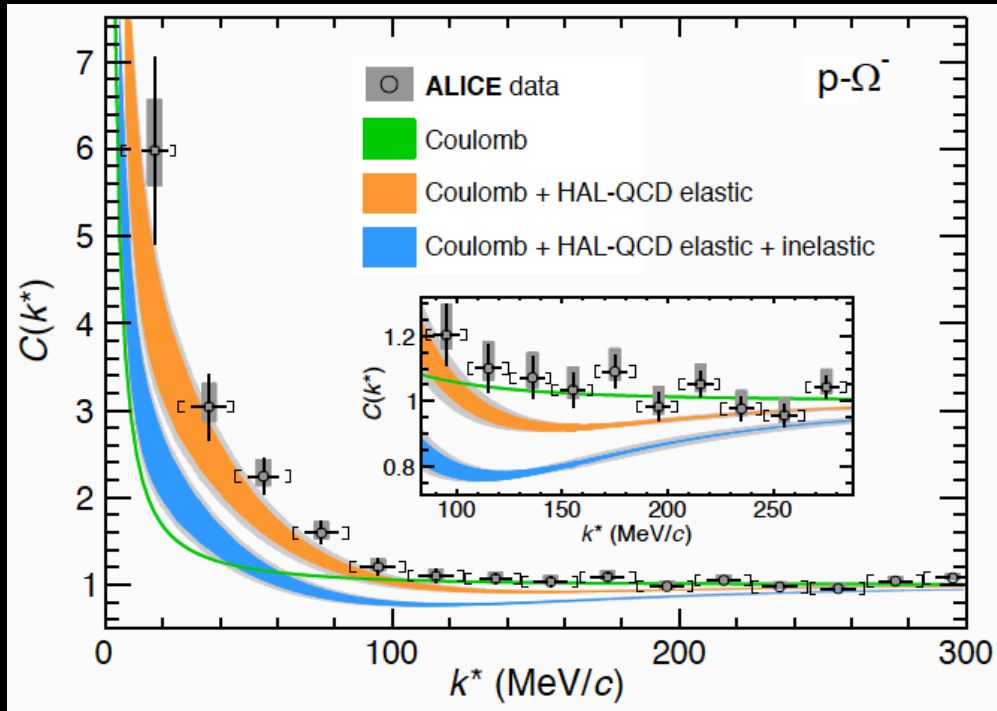
Phys. Rev. Lett. 123, 112002 (2019)



- $p$ - $\Xi$  interaction is attractive
- No indication of bound state in data
- $p$ - $\Xi$  interaction stronger than Coulomb → observation of strong interaction
- Coulomb + HAL QCD in agreement with  $p$ - $\Xi$  measurements

# p- $\Omega$ correlation function in pp at 13 TeV

Nature 588 (2020) 232

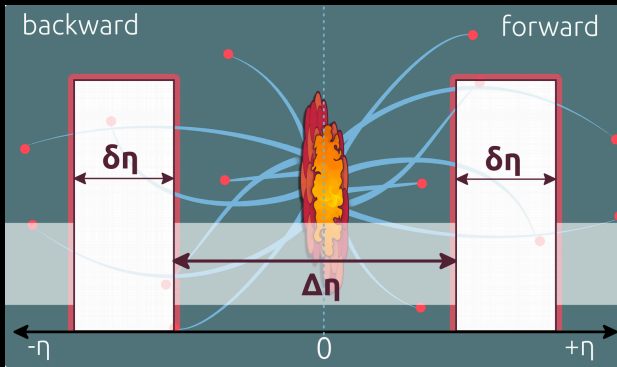


- p- $\Omega$  interaction is attractive
- No indication of bound state in data
- p- $\Omega$  interaction stronger than Coulomb  $\rightarrow$  observation of strong interaction
- Calculations underestimate p- $\Omega$  measurements for both cases tested for missing p- $\Omega$  inelastic channels in  $^3S_1$  state
  - Inelastic channel dominated by absorption
  - Neglecting inelastic channel

# CORRELATIONS AND FLUCTUATIONS

# Forward-backward correlations with strongly intensive quantity $\Sigma$

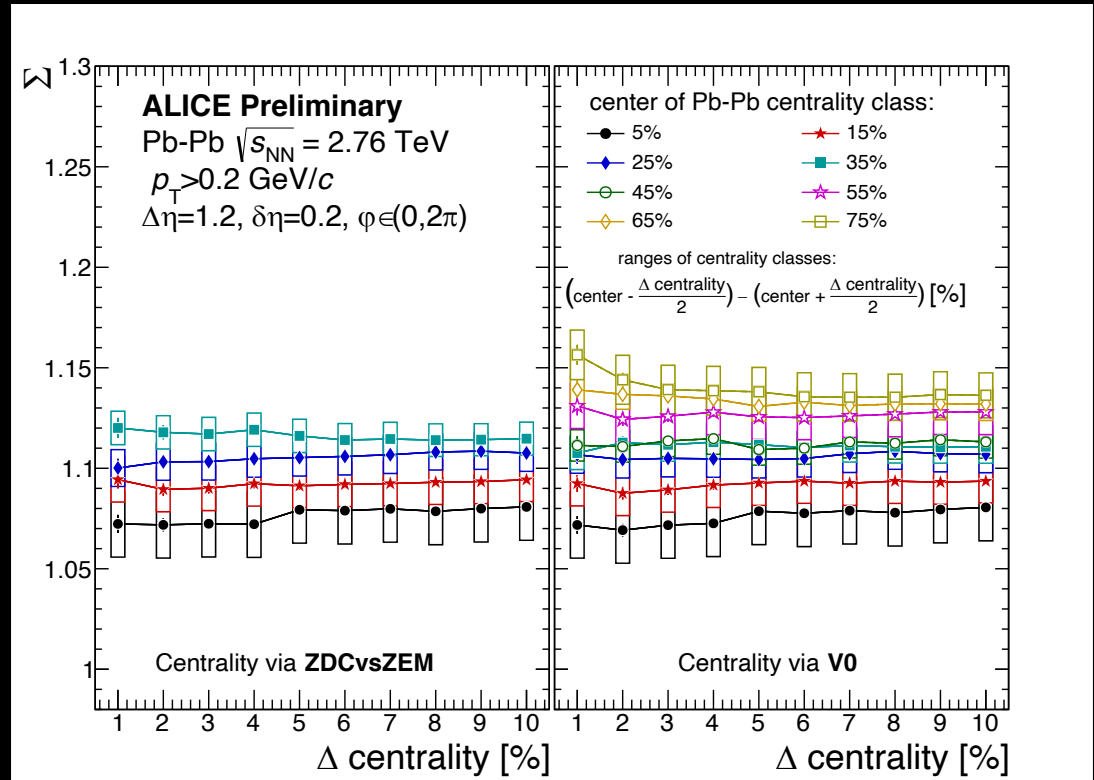
$\Sigma$  provides information on the early collision dynamics



$$\Sigma = \frac{1}{\langle n_B \rangle + \langle n_F \rangle} [\langle n_F \rangle \omega_B + \langle n_B \rangle \omega_F - 2 \text{Cov}(n_F, n_B)]$$

$$\omega_{B(F)} = \frac{\text{Var}(n_{B(F)})}{\langle n_{B(F)} \rangle}$$

M. Gaździcki and M. I. Gorenstein,  
PRC84 (2011) 014904



- $\Sigma$  is independent from event centrality estimator and width of the centrality interval

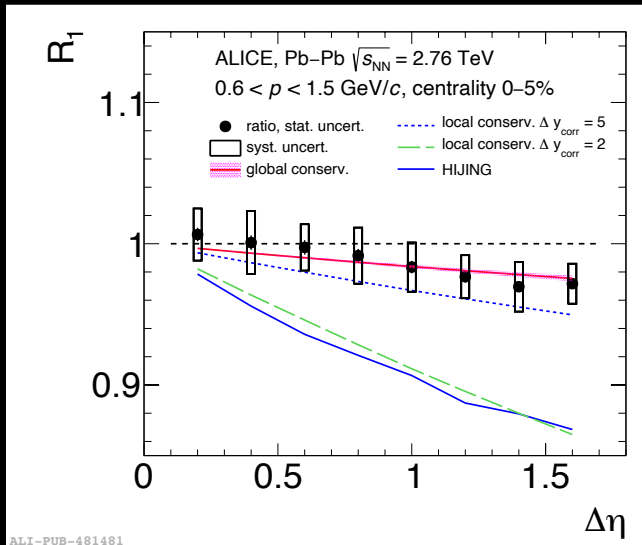
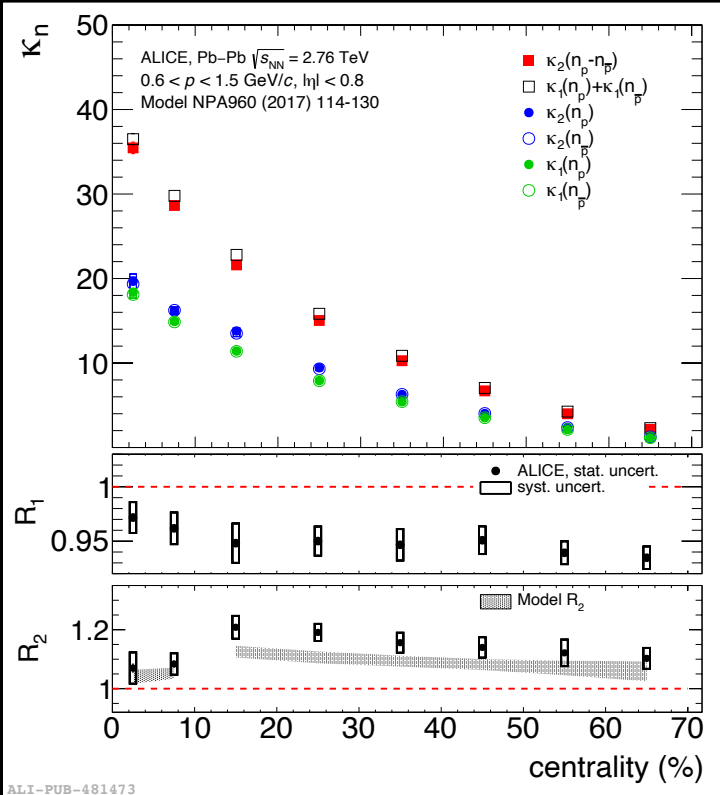
# Net-proton fluctuations in Pb-Pb at 2.76 TeV

$$\kappa_1(\Delta n_B) = \sum_{\Delta n_B=-\infty}^{\infty} \Delta n_B P(\Delta n_B) = \langle \Delta n_B \rangle,$$

$$\kappa_2(\Delta n_B) = \sum_{\Delta n_B=-\infty}^{\infty} (\Delta n_B - \langle \Delta n_B \rangle)^2 P(\Delta n_B) = \langle (\Delta n_B - \langle \Delta n_B \rangle)^2 \rangle$$

$$R_1 = \kappa_2(n_p - n_{\bar{p}}) / \langle n_p + n_{\bar{p}} \rangle, \quad R_2 = \kappa_2(n_p) / \langle n_p \rangle$$

Phys. Lett. B 807 (2020) 135564



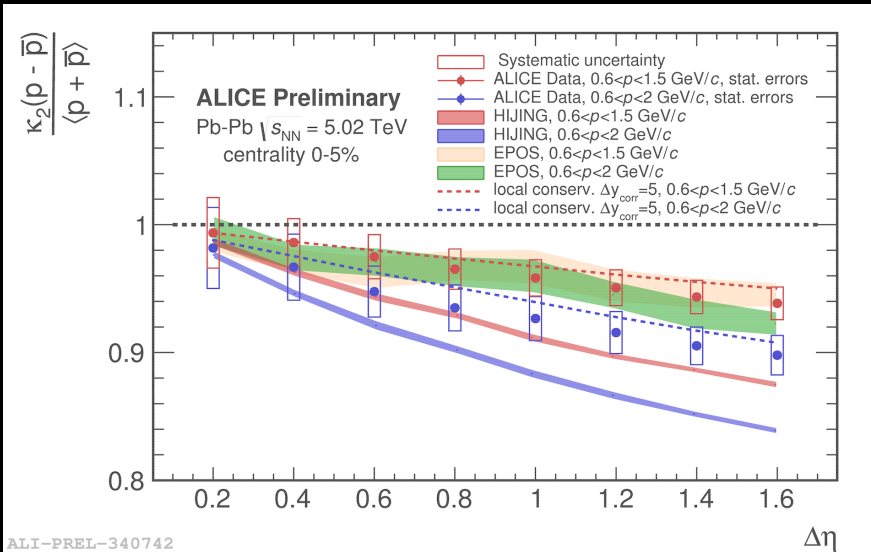
- Event-by-event baryon number conservation leads to subtle long-range correlations arising from very early interactions in the collisions
- Experimental tests of lattice QCD predictions on second and higher order cumulants of net-baryon distributions → critical behavior near QCD phase boundaries

# Net-proton fluctuations in Pb-Pb at 5.02 TeV

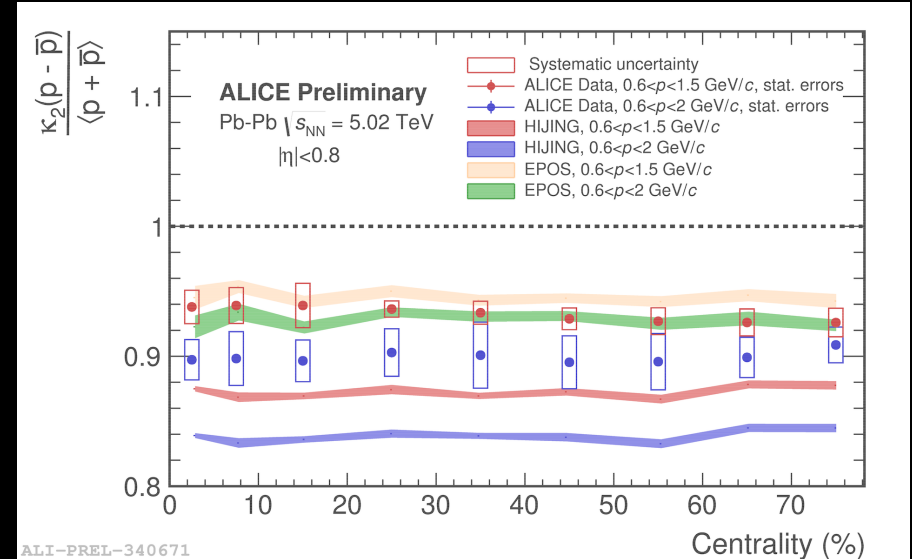
$$\kappa_1(\Delta n_B) = \sum_{\Delta n_B=-\infty}^{\infty} \Delta n_B P(\Delta n_B) = \langle \Delta n_B \rangle,$$

$$\kappa_2(\Delta n_B) = \sum_{\Delta n_B=-\infty}^{\infty} (\Delta n_B - \langle \Delta n_B \rangle)^2 P(\Delta n_B) = \langle (\Delta n_B - \langle \Delta n_B \rangle)^2 \rangle$$

$$R_1 = \kappa_2(n_p - n_{\bar{p}}) / \langle n_p + n_{\bar{p}} \rangle, \quad R_2 = \kappa_2(n_p) / \langle n_p \rangle$$



ALI-PREL-340742



ALI-PREL-340671

- Event-by-event baryon number conservation leads to subtle long-range correlations arising from very early interactions in the collisions
- Experimental tests of lattice QCD predictions on second and higher order cumulants of net-baryon distributions  $\rightarrow$  critical behavior near QCD phase boundaries

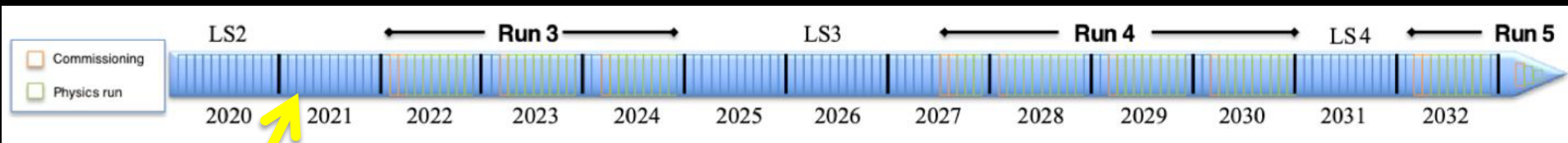
# Summary



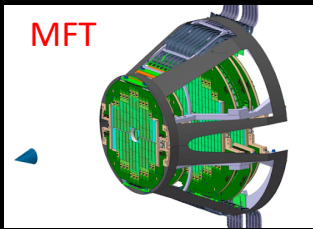
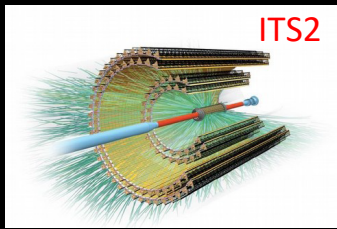
- Relative to pion particle production is driven by characteristics of final state
- Several mechanisms of hadron and nuclei production: (re)scattering, (re)combination, coalescence etc.
- Collective effects are observed in small and large systems
- The strong interaction between protons and hyperons ( $p$ - $\Xi$  and  $p$ - $\Omega$ ) is attractive (no indication of bound states)
- Strongly intensive quantities allow for particle correlation studies free from volume size and volume fluctuations
- Baryon number conservation introduces a subtle long-range correlation in net-baryon correlation studies



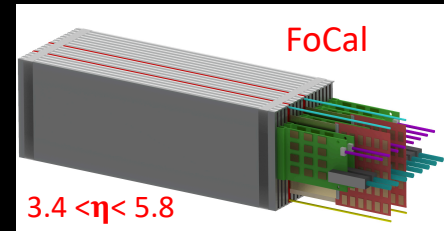
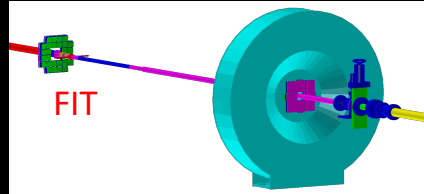
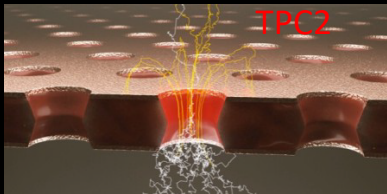
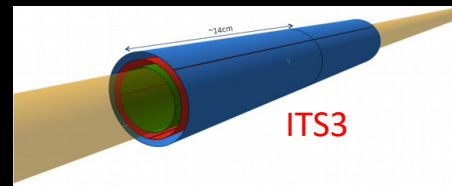
# Outlook



## LS2: upgrades



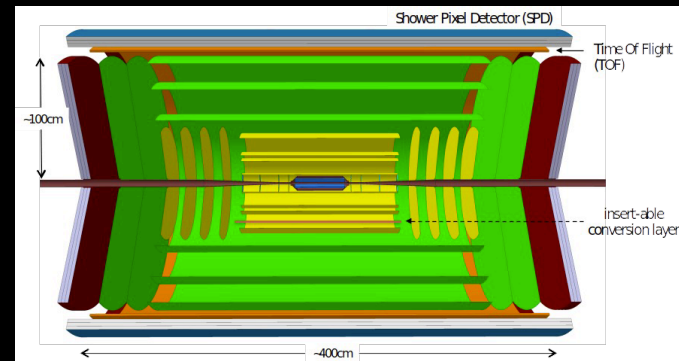
## LS3: upgrades



## LS4: Future heavy-ion detector [arXiv:1902.01211](https://arxiv.org/abs/1902.01211)

### Continuous data taking

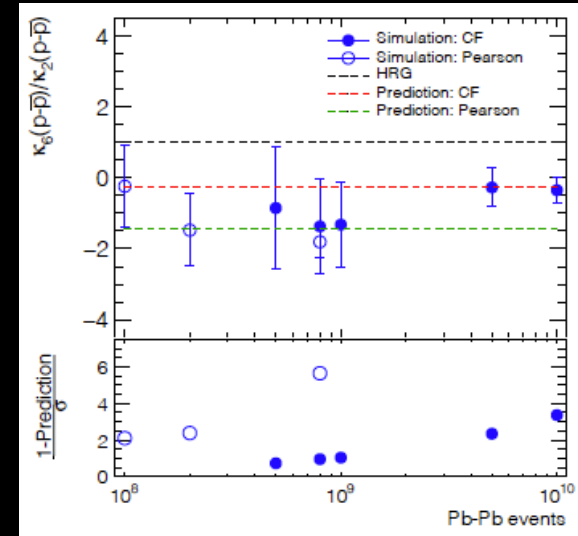
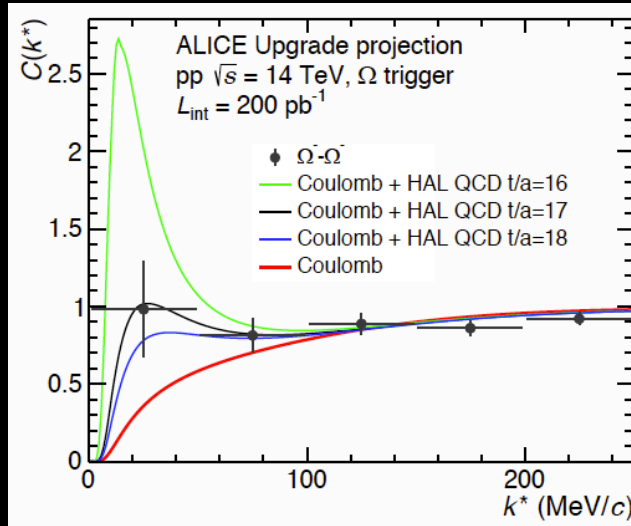
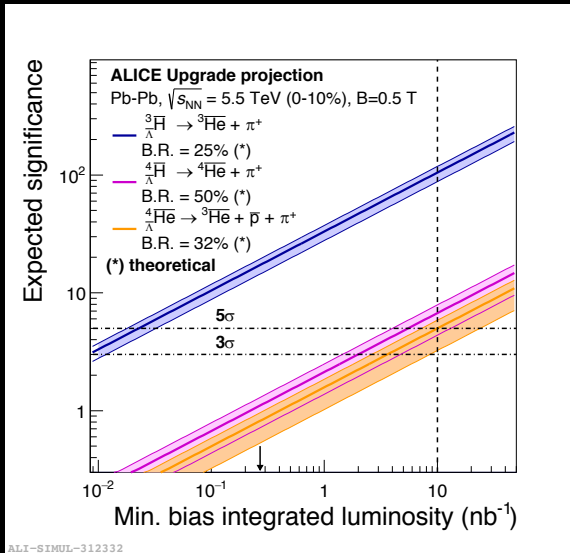
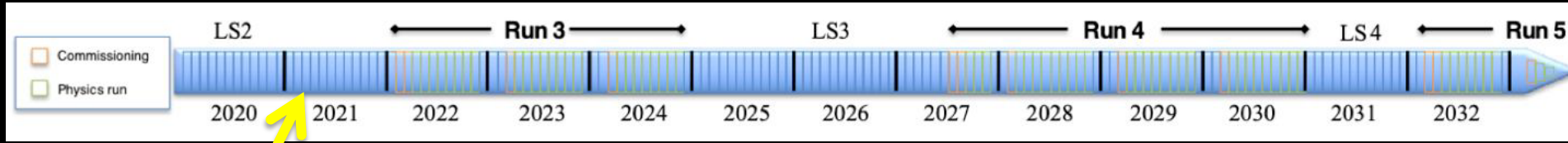
- Detector upgrade
- Online-offline computing system upgrade
- Readout electronics and trigger upgrade



# Outlook



## LHC Run3 projected statistic



More than  $3\sigma$  significance for  $\frac{4}{\Lambda}\bar{\text{He}}$

$\Omega$ - $\Omega$  correlations

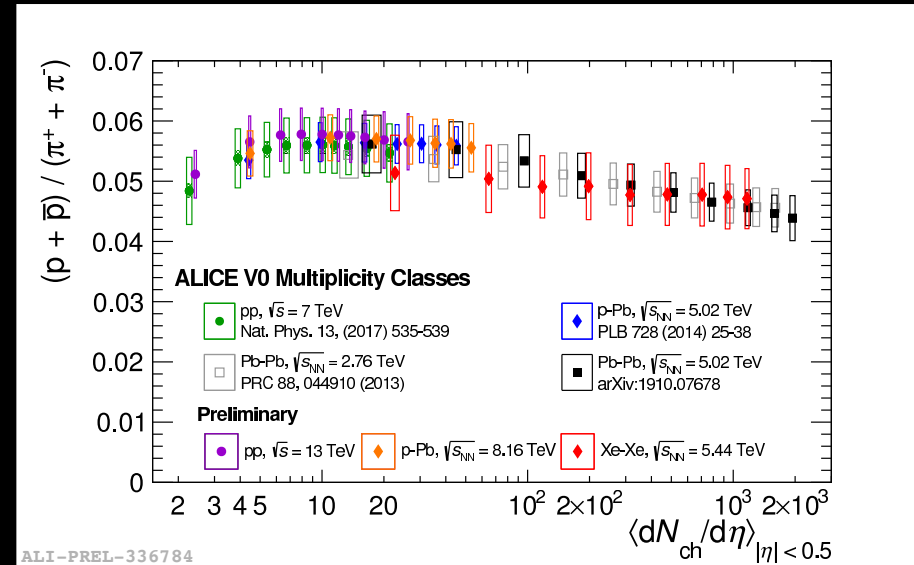
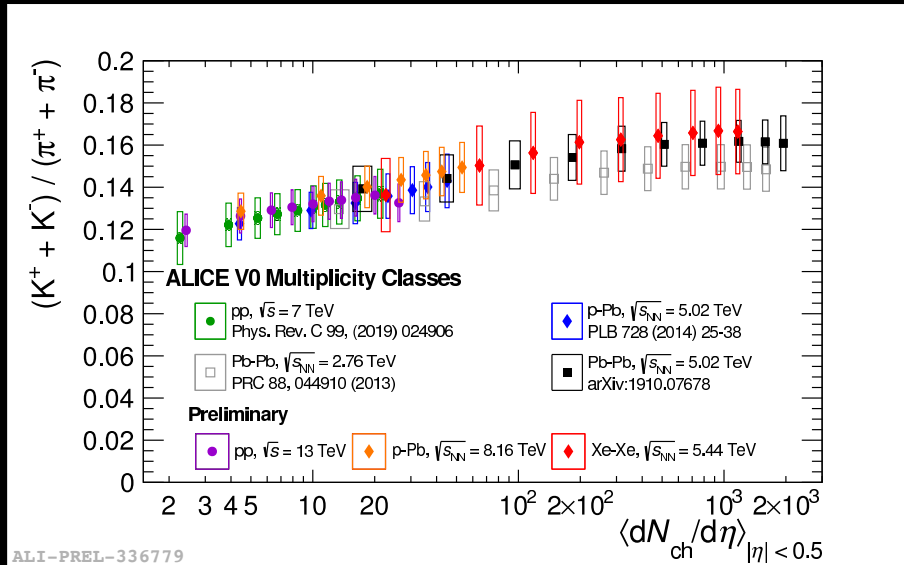
Net-proton 6<sup>th</sup> order cumulants

**THANK YOU FOR YOUR ATTENTION!**

# Backup



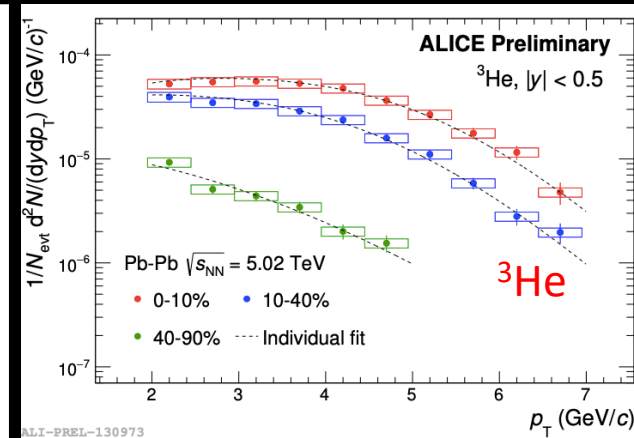
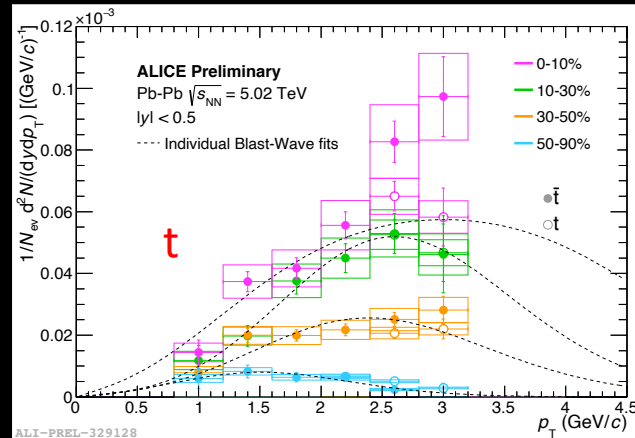
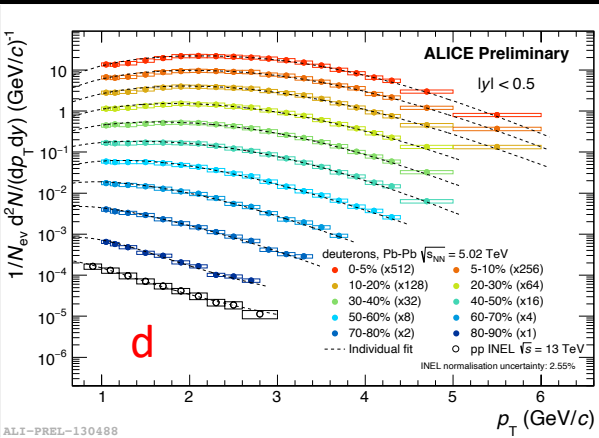
# $K/\pi$ and $p/\pi$ in pp, p-Pb, Pb-Pb and Xe-Xe



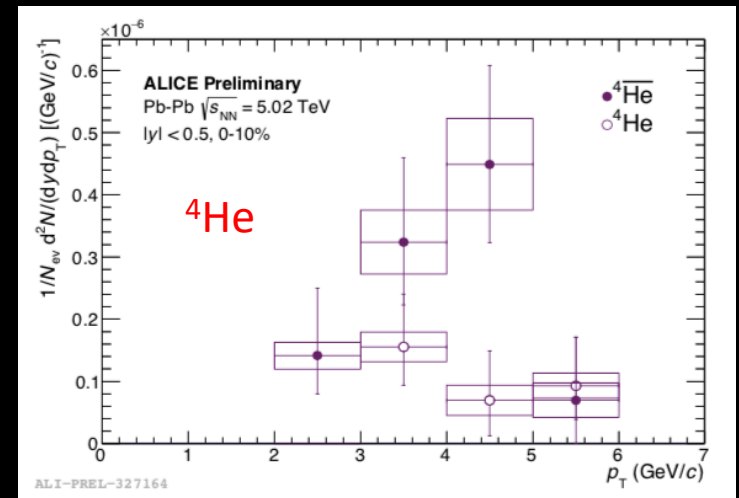
- No significant energy dependence is observed
- $K/\pi$  and  $p/\pi$  are consistent for all collision systems at similar multiplicity

→ Particle production is driven by the characteristics of final state

# (Anti)nuclei $p_T$ spectra in Pb-Pb

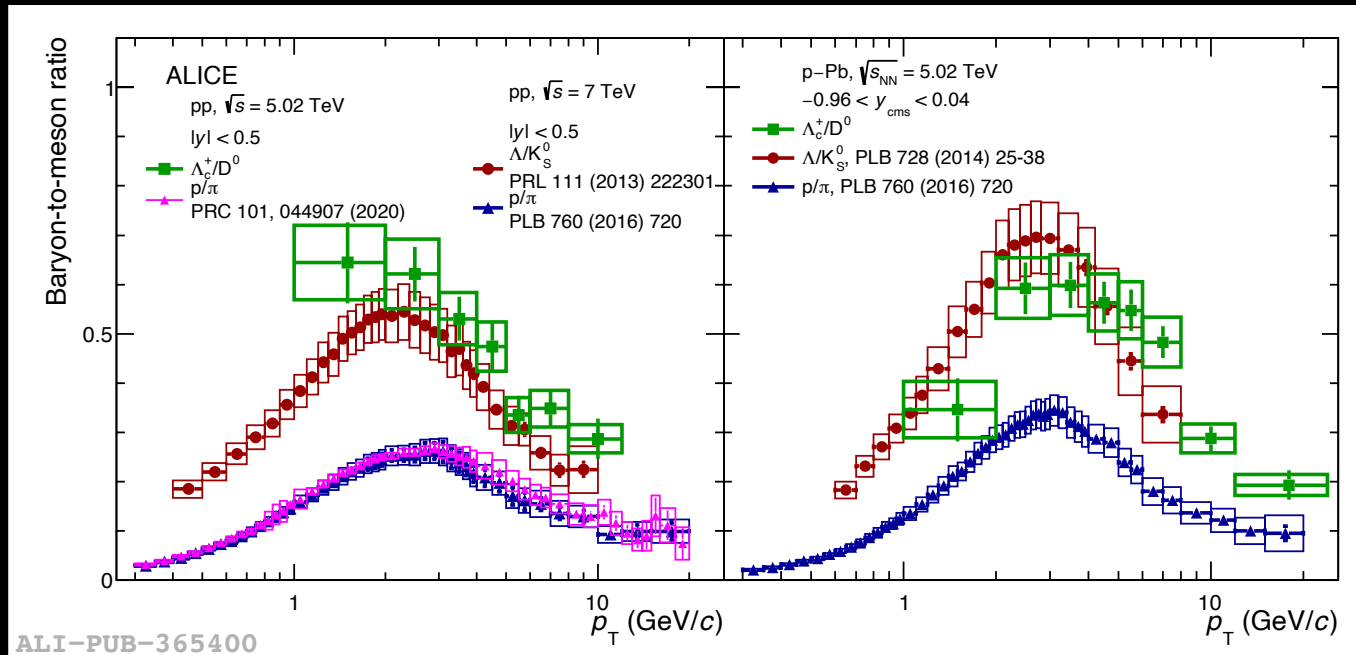


- Light nuclei  $p_T$  spectra are modified by radial flow



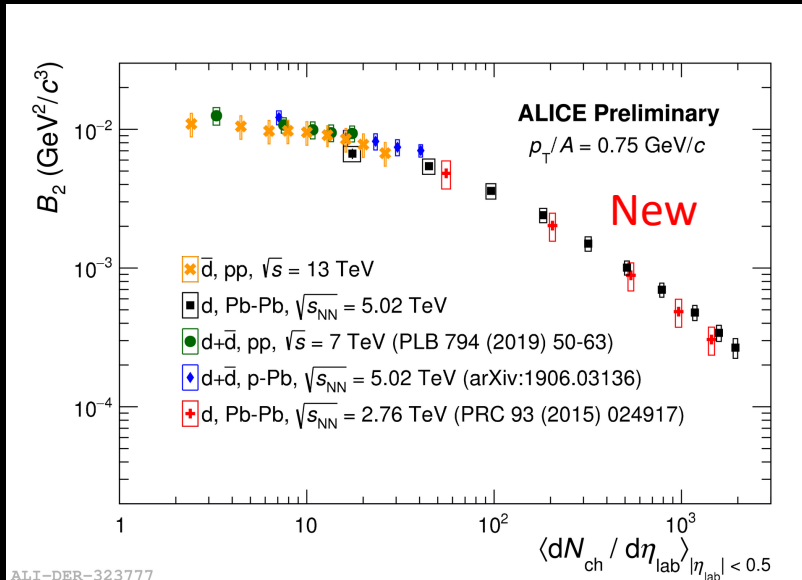
# Baryon-to-meson ratios

arXiv:2011.06078



- Baryon-to-meson ratio for  $\Lambda_c/D_0$  show similarities with those for light-flavor  $p/\pi$  and  $\Lambda/K^0$ s
  - hint for the common production mechanism of light- and heavy-flavor baryons (coalescence vs. fragmentation)

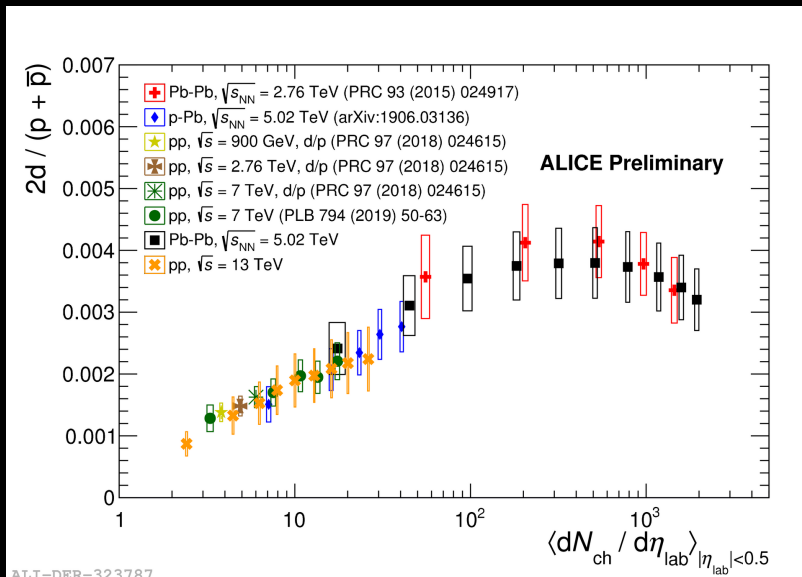
# Formation of light nuclei: (anti) deuterons



- Coalescence of baryons close in phase space (A - mass number)

$$E_A \frac{d^3 N_A}{dp_A^3} = B_A \left( E_p \frac{d^3 N_p}{dp_p^3} \right)^A$$

- $B_2$  shows dependence on multiplicity (no dependence on  $p_T$ )



- d/p vs multiplicity
  - Increase from pp to peripheral Pb-Pb consistent with coalescence model
  - No centrality dependence in high multiplicity Pb-Pb (yields consistent with thermal model)

→ Production mechanisms: thermal vs coalescence?



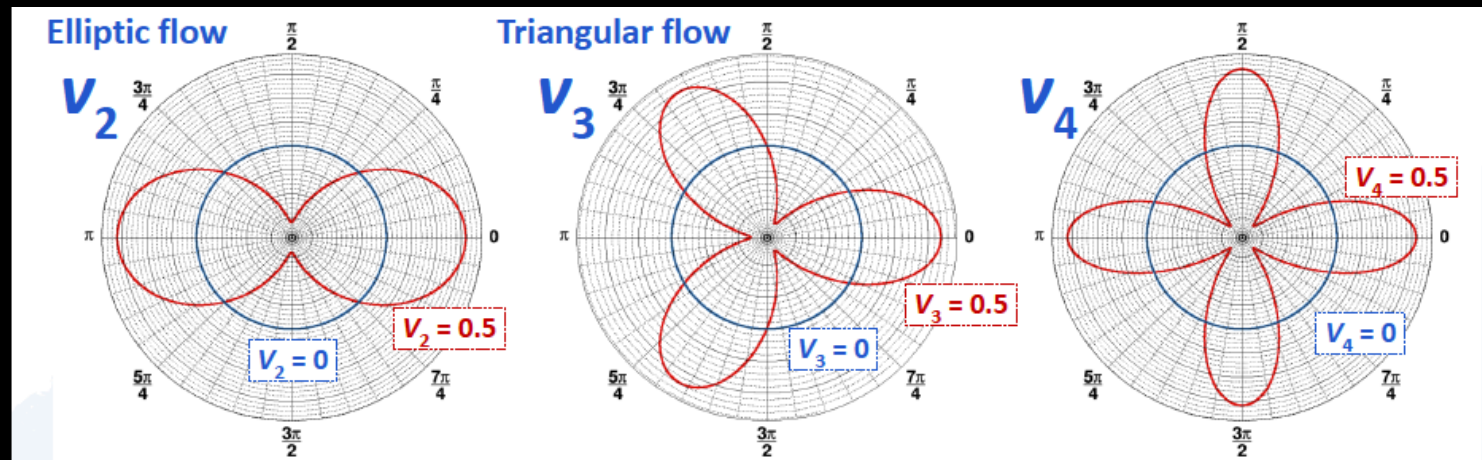
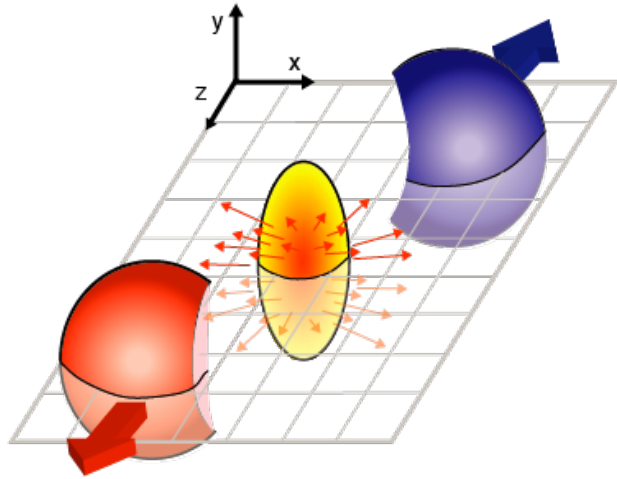
# ANISOTROPIC FLOW

# Anisotropic flow

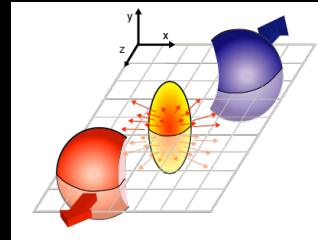
- Strongly interacting system:  
spatial anisotropy  $\rightarrow$  momentum anisotropy
- Quantified in terms of Fourier coefficients  $v_n$

$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi p_T dp_T dy} \left( 1 + 2 \sum_{n=1}^{\infty} v_n \cos[(\varphi - \Psi_n)] \right)$$

$$v_n(p_T, y) = \langle \cos[n(\varphi - \Psi_n)] \rangle$$



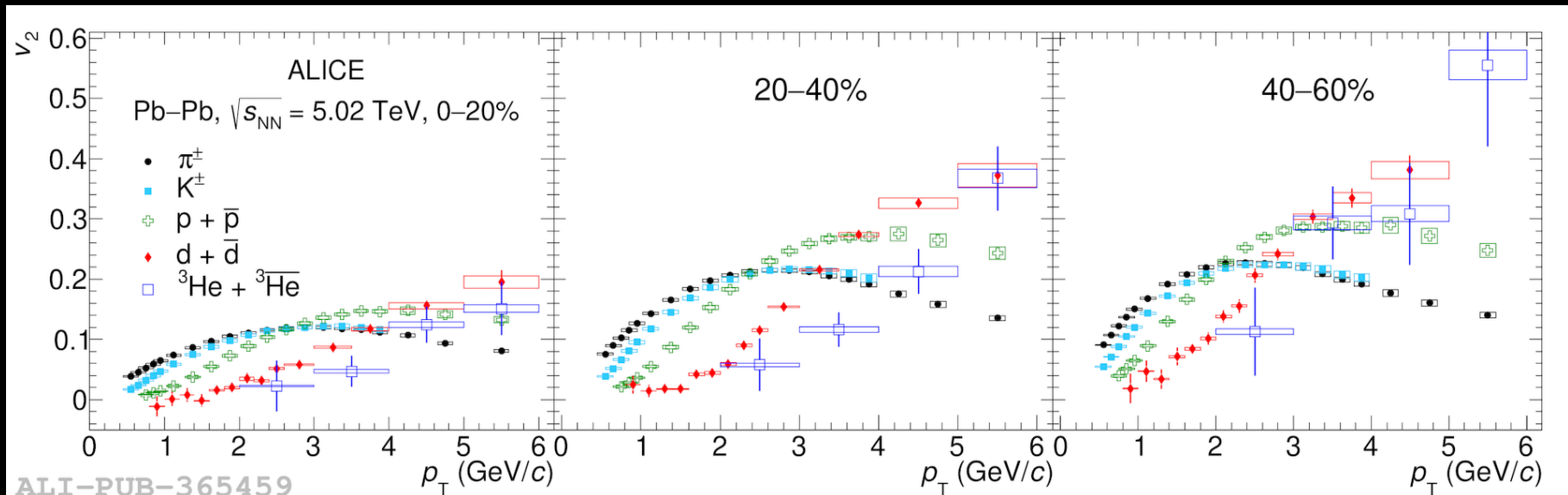
# Hadron and nuclei $v_2$ in Pb-Pb



$$f(\varphi) = \frac{1}{2\pi} \left[ 1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\varphi - \Psi_n)] \right]$$

$$v_n(p_T, y) = \langle \cos[n(\varphi - \Psi_n)] \rangle$$

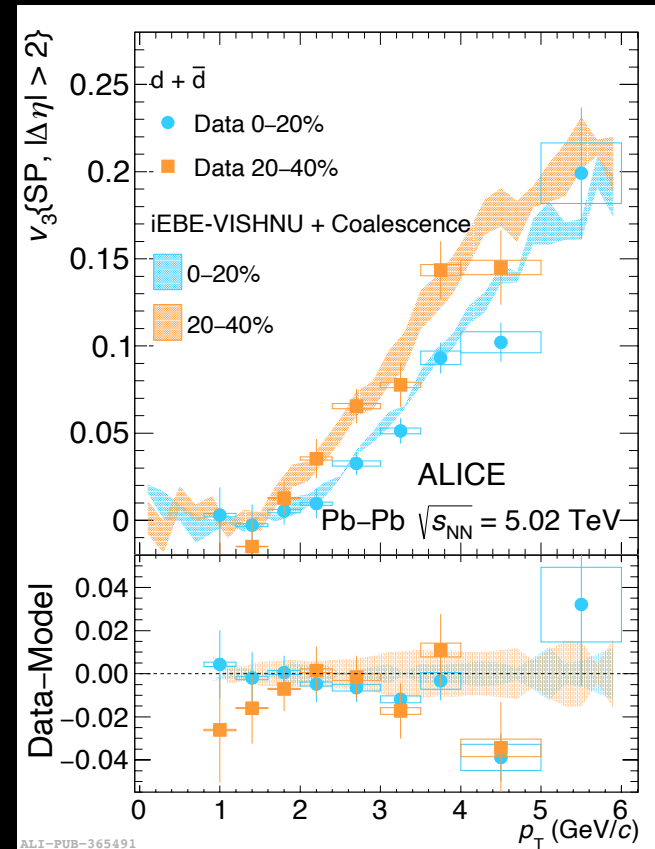
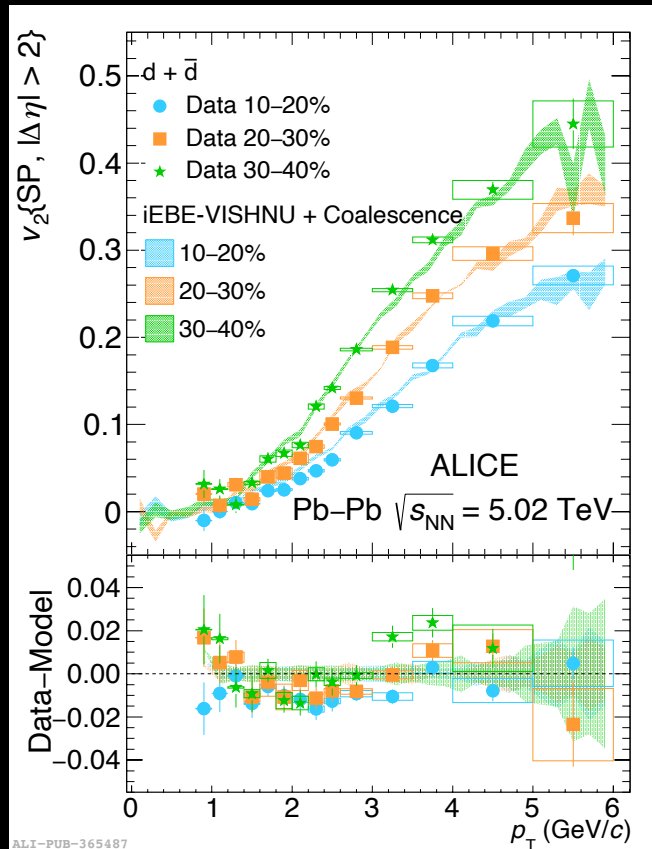
Phys. Rev. C 102 (2020) 055203



- $p_T < 2$  GeV/c: mass ordering
- $p_T \sim 2.5$  GeV/c: crossing between  $v_2$  of baryons and mesons
- $p_T > 2.5$  GeV/c: baryons  $v_2 >$  mesons  $v_2$  (flow driven by quark content)
- Light nuclei participating in flow

# (Anti)deuteron $v_2$ and $v_3$ in Pb-Pb

Phys. Rev. C 102 (2020) 055203



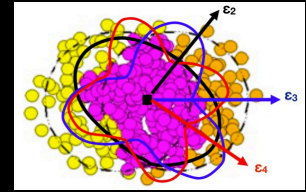
- Coalescence model with phase-space distributions of protons and neutrons from iEBE-VISHNU in agreement with data

# Correlations between flow amplitudes

Observables sensitive to initial state and dynamical evolution of QGP

Symmetric cumulants

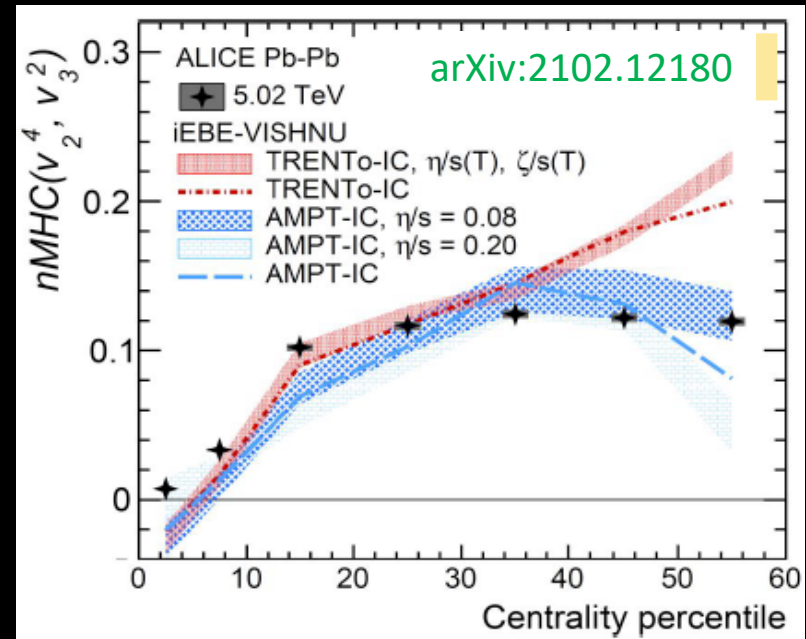
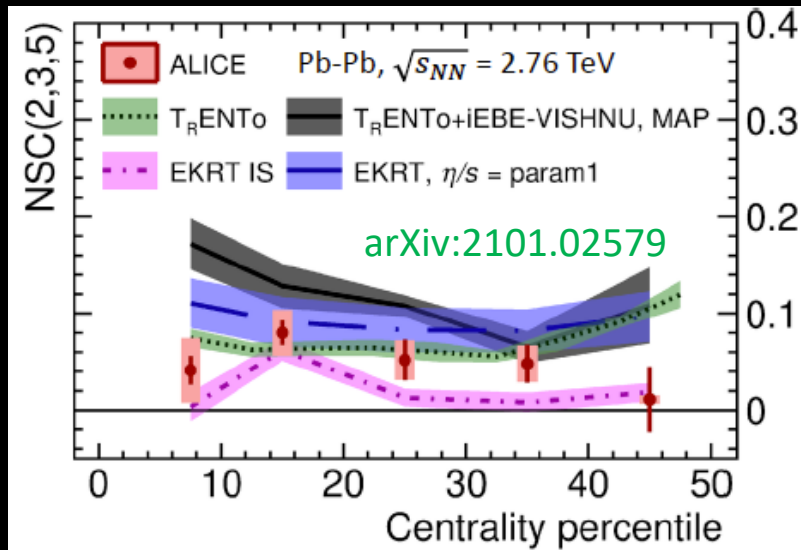
$$SC(k, l, m) \equiv \langle v_k^2 v_l^2 v_m^2 \rangle - \langle v_k^2 v_l^2 \rangle \langle v_m^2 \rangle - \langle v_k^2 v_m^2 \rangle \langle v_l^2 \rangle - \langle v_l^2 v_m^2 \rangle \langle v_k^2 \rangle + 2 \langle v_k^2 \rangle \langle v_l^2 \rangle \langle v_m^2 \rangle$$



$$NSC(k, l, m) \equiv \frac{SC(k, l, m)}{\langle v_k^2 \rangle \langle v_l^2 \rangle \langle v_m^2 \rangle}$$

Mixed harmonic cumulants

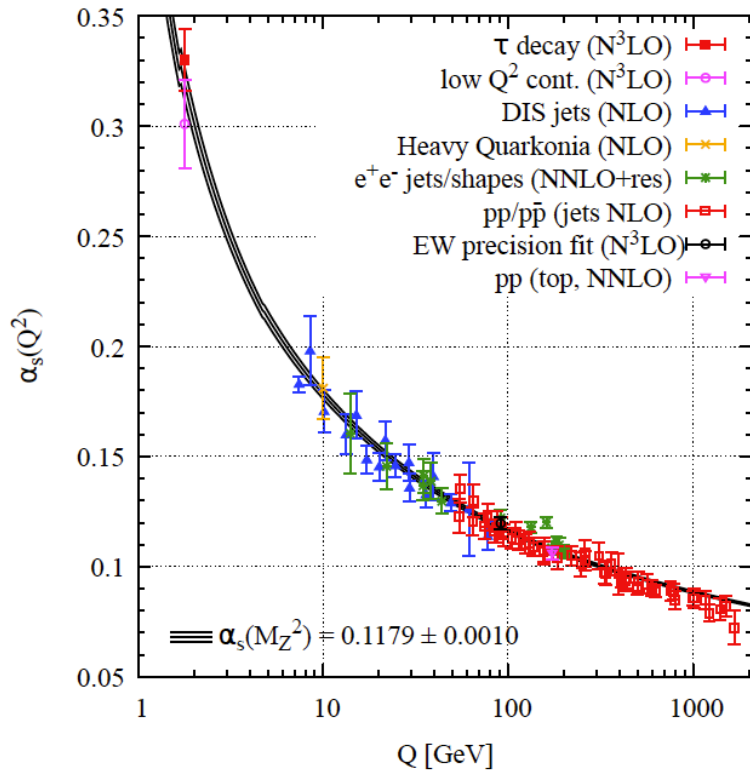
$$MHC(v_2^4, v_3^2) = \langle v_2^4 v_3^2 \rangle - 4 \langle v_2^2 v_3^2 \rangle \langle v_2^2 \rangle - \langle v_2^4 \rangle \langle v_3^2 \rangle + 4 \langle v_2^2 \rangle^2 \langle v_3^2 \rangle$$



- Symmetric cumulants and mixed harmonic cumulants are insensitive to non-flow contributions

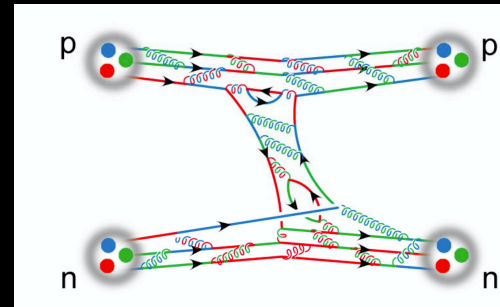
# The strong interaction

PDG 2020



Running coupling constant  $\alpha_s$ :

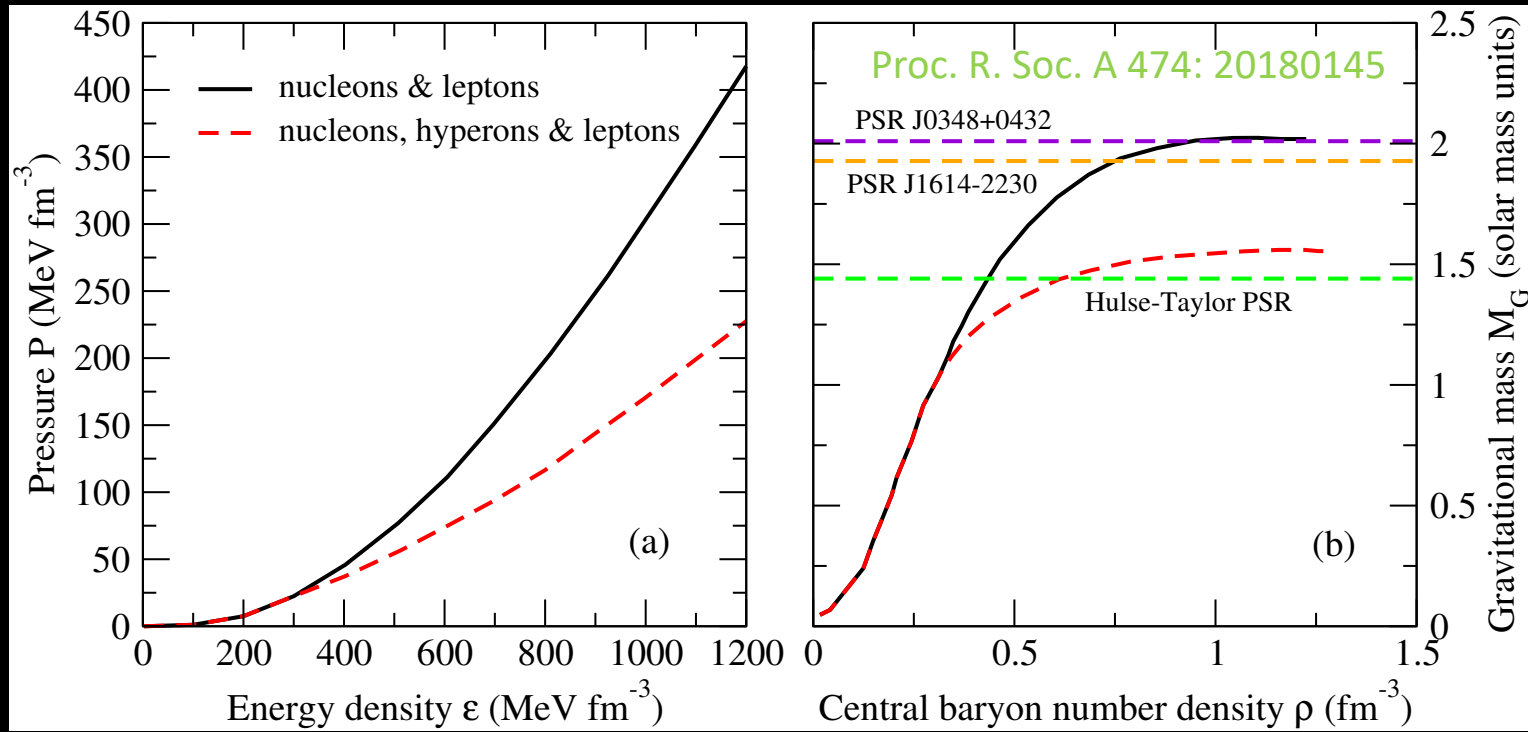
- Asymptotic freedom at small distance ( $Q^2 \gg \Lambda_{\text{QCD}}^2$ )
- Confinement at large distance ( $Q \lesssim 1 \text{ GeV}$ )
  - Perturbative methods not applicable
  - Effective theories with hadrons as degrees of freedom
  - **Lattice QCD calculations**



$$\alpha_s(Q^2) \equiv \frac{g_s^2(Q^2)}{4\pi} = \frac{12\pi}{(33 - 2N_f)\ln(Q^2/\Lambda_{\text{QCD}}^2)}$$

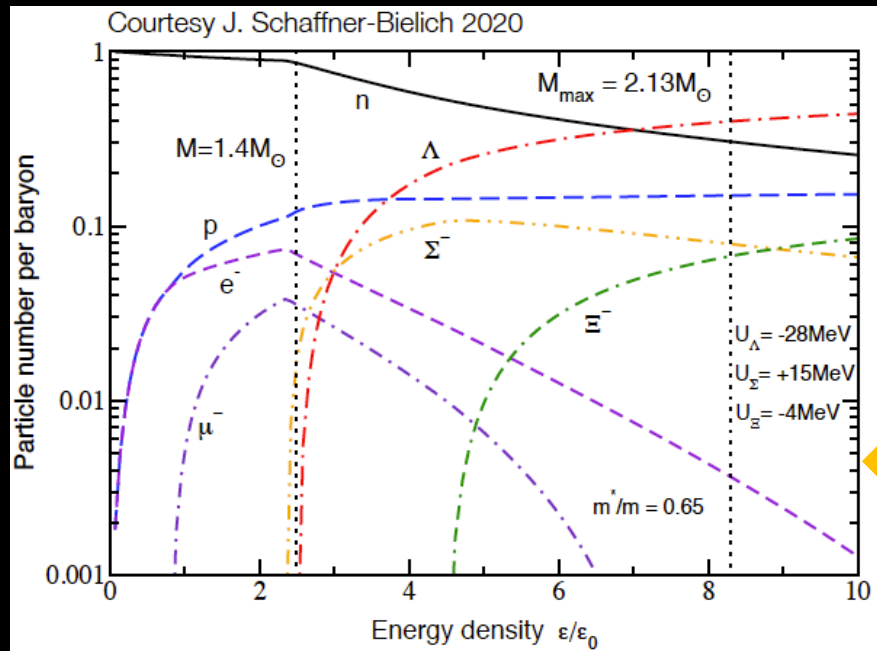
What is the strong interaction between different hadron species?

# Hyperon puzzle

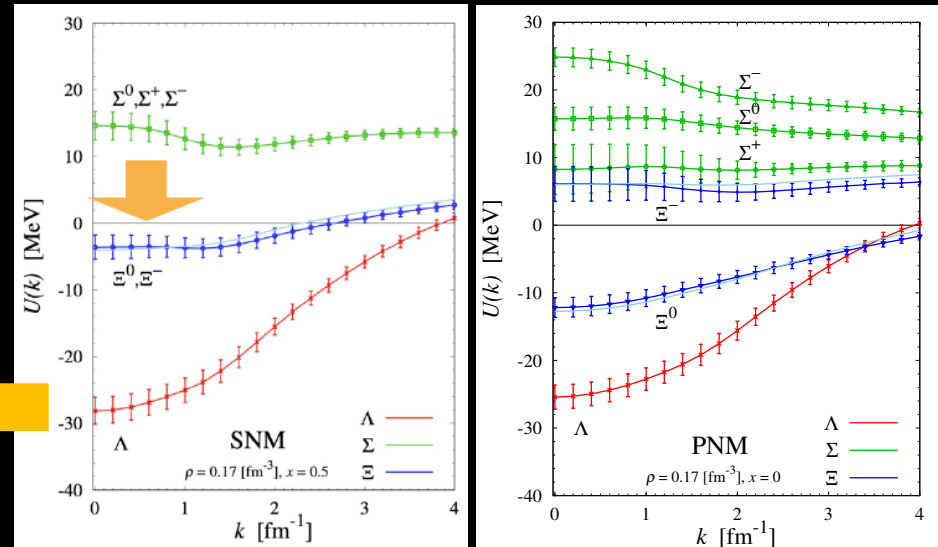


- Hyperons may appear in the inner core of neutron stars at densities of about  $2-3\rho_0$
- Their presence in the neutron star interior leads to a softening of the EoS and consequently to a reduction of the maximum mass (current predictions  $1.4-1.8 M_\odot$ )
- **Astrophysical observations of pulsars rule out almost all currently proposed EoS with hyperons...?**
- **Additional repulsion: Y-Y repulsive potential, hyperonic three-body forces (e.g. NNY, NYY, YYY), quark-gluon plasma below the hyperon threshold (hybrid neutron stars)...?**

# Hyperon puzzle and lattice QCD



HAL QCD, T. Ionue et al. AIP Conf. Proc. 2130 (2019) 1, 020002



Hyperon single-particle potential

- $\Xi^-$  attractive single-particle potential in symmetric nuclear matter (SNM) and repulsive in pure neutron matter (PNM)
- $\Xi^-$  appears at larger densities in neutron stars

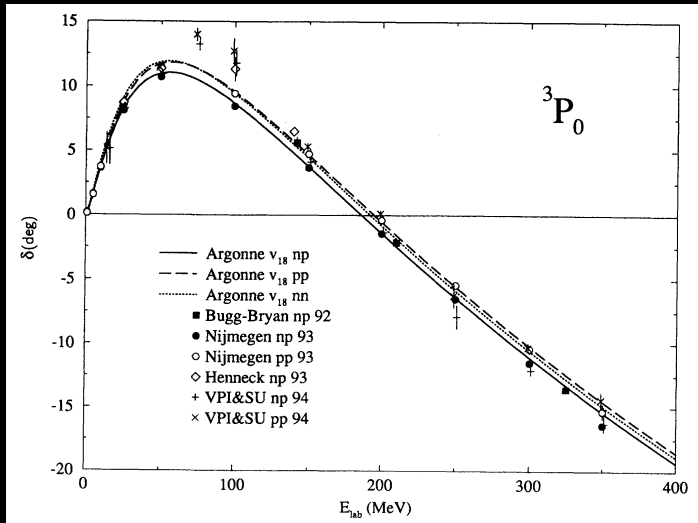
→ Resulting EoS of neutron stars is stiffer and matches astrophysical observation...



# Experimental data to constrain the strong interaction

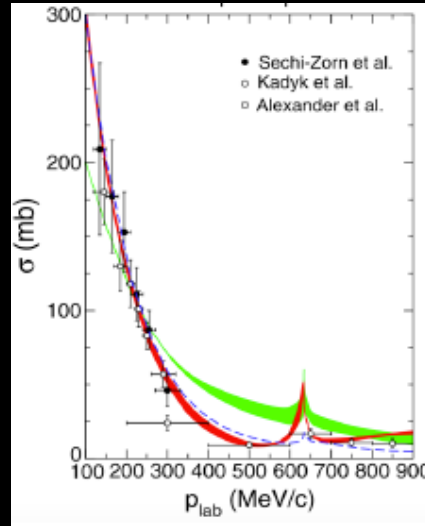


$N+N \rightarrow N+N$



R. B. Wiringa et al., PRC 51 (1995) 38

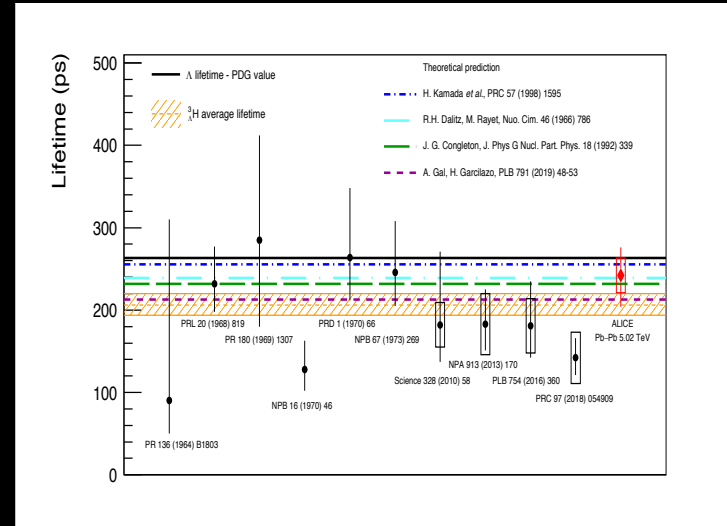
$N + \Lambda \rightarrow N + \Lambda$



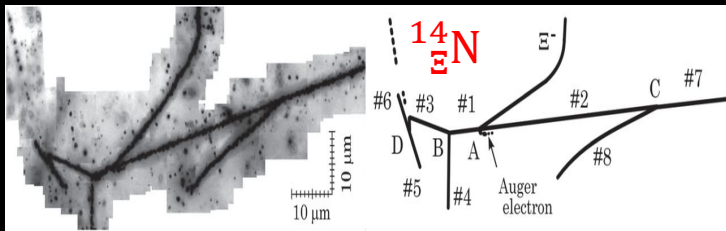
LO: H. Poinder et al. NPA 779 (2006) 244

NLO: J. Haidenbauer et al. NPA 915 (2013) 24

Hypertriton lifetime



ALICE PLB797 (2019) 134905

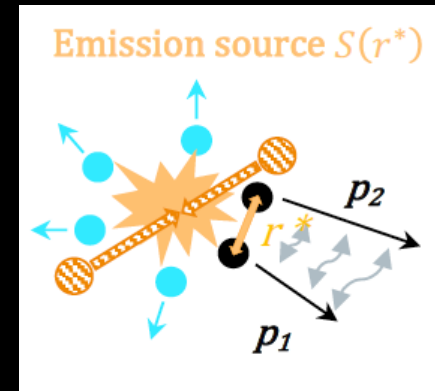
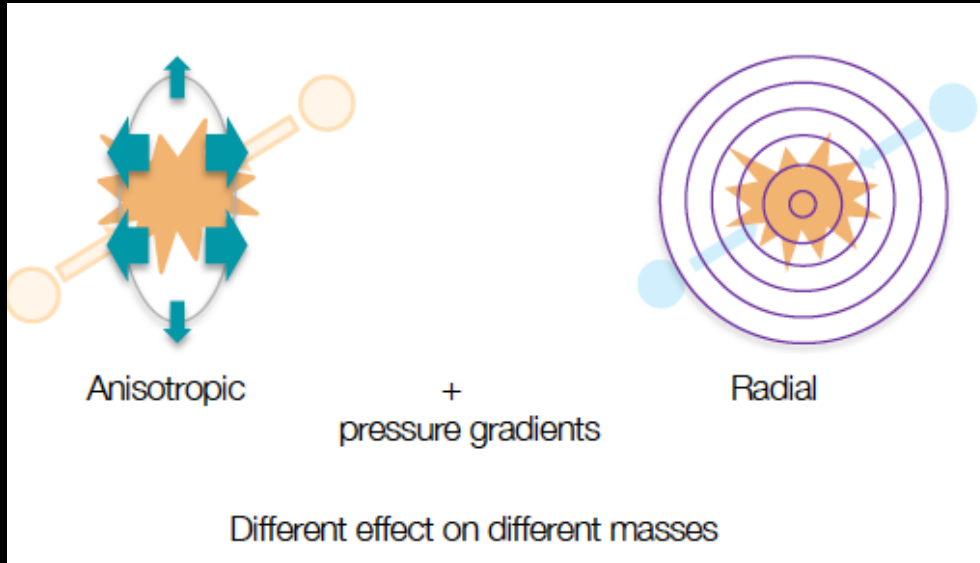


K. Nakazawa, PTEP 2015 (2015) 033D02

- Good constraints for NN interaction
- Small statistics of scattering data for hyperons
- $\sim 1000$   $\Lambda$ -hypernuclei and **one  $^{14}_{\Lambda}\text{N}$  discovered by now**

# Common baryon source

Collective flow and feed-down from short lived resonances modify source size



Resonances with  $c\tau \sim r_0 \sim 1\text{fm}$  ( $\Delta^{++}, N^*, \Sigma^*$ )

Particle	Primordial fraction	Resonances $\langle c\tau \rangle$
Proton	33 %	1.6 fm
Lambda	34 %	4.7 fm

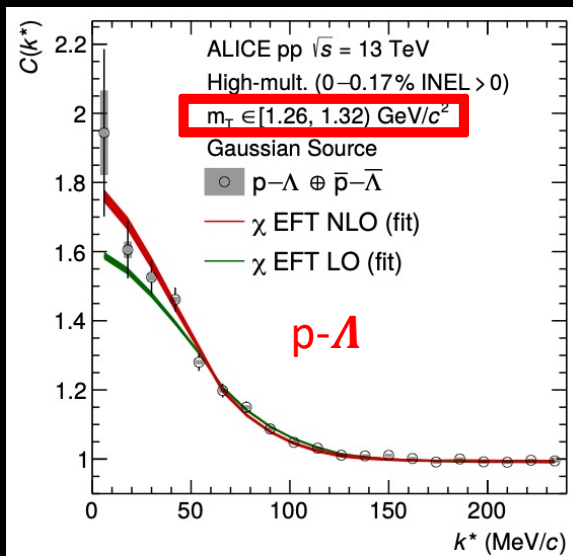
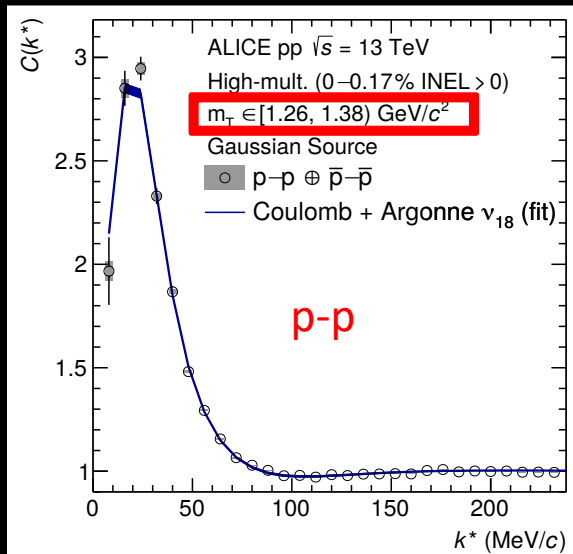
U. Wiedemann U. Heinz (PRC56 R610, 1997)

$$m_T = \sqrt{k_T^2 + m^2}, k_T = \frac{1}{2} |p_{T1} + p_{T2}|$$

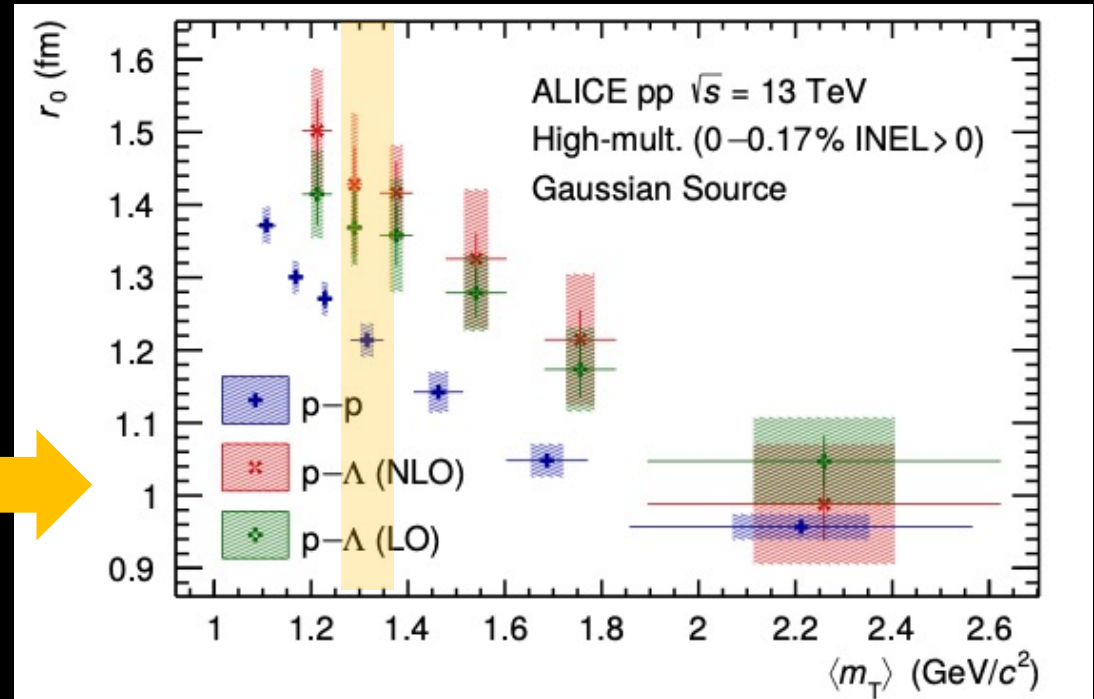
$$S(r^*) = G(r^*, r_{core}(m_T)) = \frac{1}{(4\pi r_{core}^2)^{3/2}} \exp\left(-\frac{r^{*2}}{4r_{core}^2}\right) \otimes E(r^*, M_{res}, \tau_{res}, p_{res})$$

$$E(r^*, M_{res}, \tau_{res}, p_{res}) = \frac{1}{s} \exp\left(-\frac{1}{s}\right), S = \beta\gamma\tau_{res} = \frac{p_{res}}{M_{res}} \tau_{res}$$

# Common baryon source



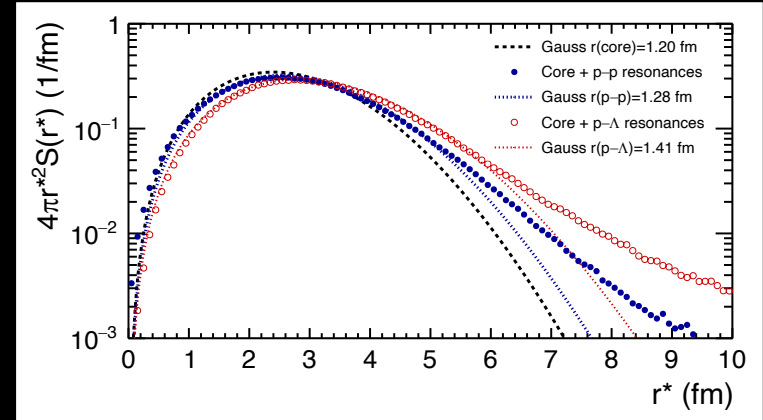
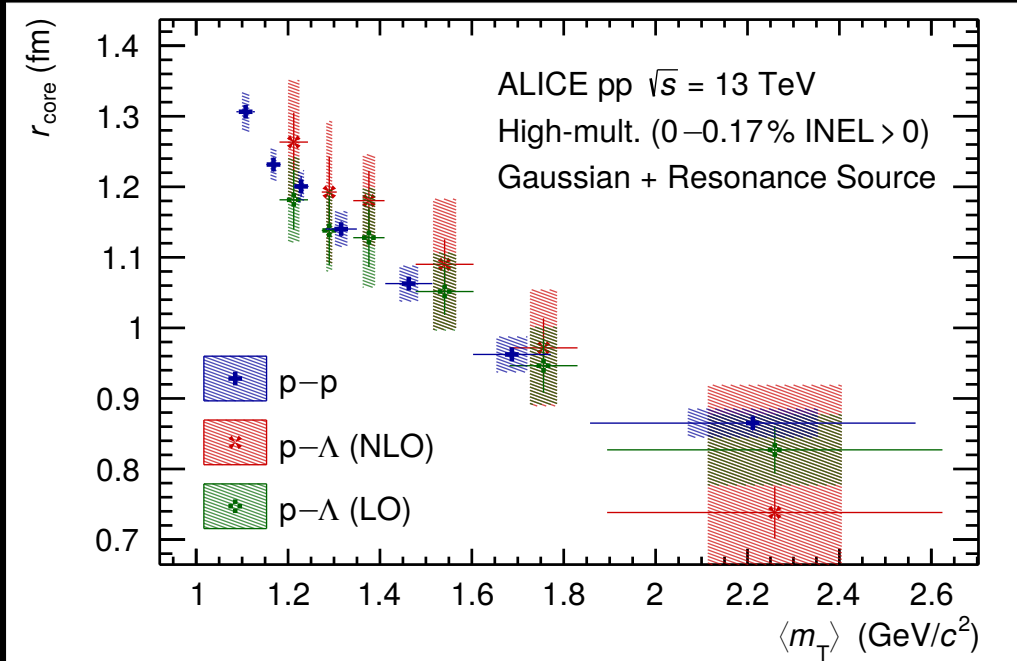
ALICE PLB 811 (2020) 135849



- Different source size depending on the baryon pair
- Gaussian source scales with  $m_T$

# Common baryon source

ALICE PLB 811 (2020) 135849



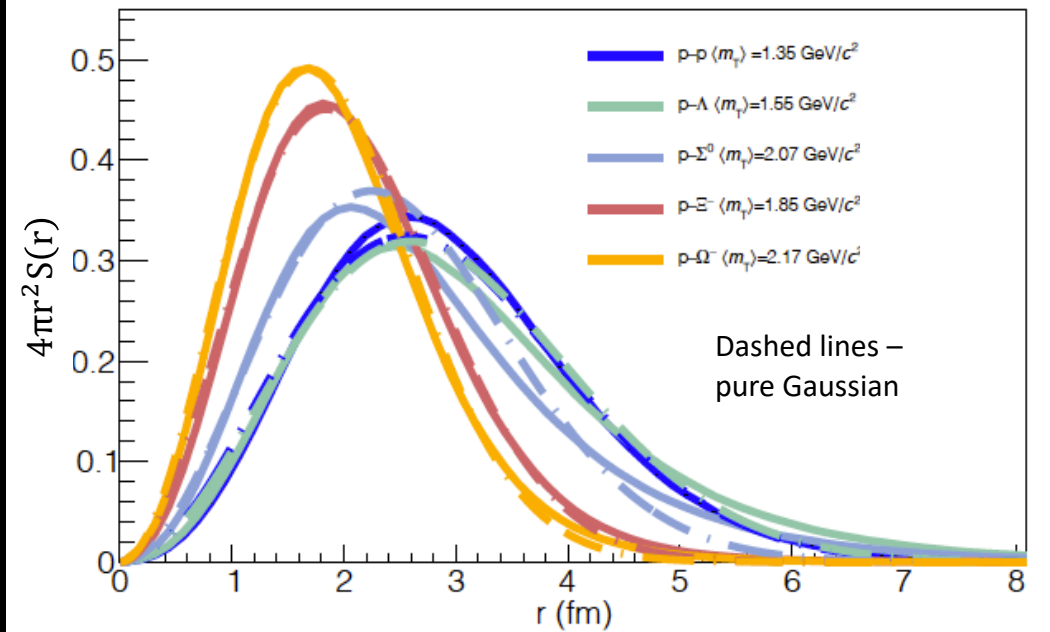
- All baryon pairs with included resonances show common  $m_T$  scaling  
→ Indication of common baryon source  
→ One can use p–p correlation to fix source size for other baryon pairs

\* F. Becattini et al. J. Phys. G: Nucl. Part. Phys. **38** 025002

\*\* T. Pierog et al. PRC 92 (2015) 034906

# Common baryon source

Source using a Gaussian core plus resonances



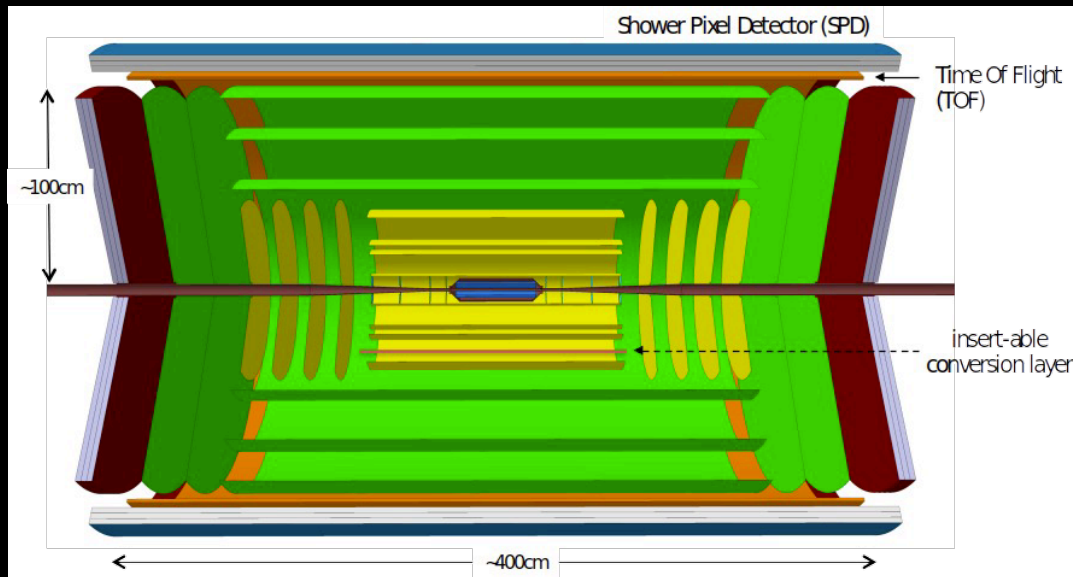
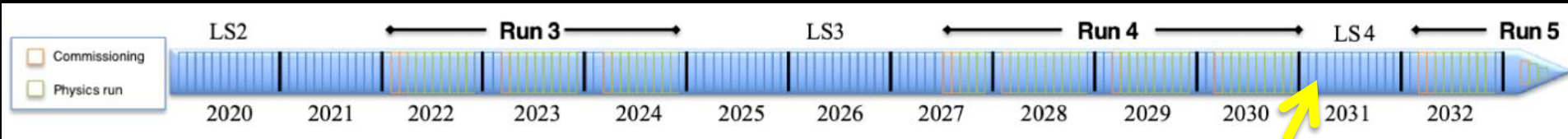
Pair	$r_{\text{Core}}$ [fm]	$r_{\text{Eff}}$ [fm]
p-p	1.1	1.2
p- $\Lambda$	1.0	1.3
p- $\Sigma^0$	0.87	1.02
p- $\Xi^-$	0.93	1.02
p- $\Omega^-$	0.86	0.95

- Emission source for heavier pairs using p-p correlation function plus resonances
- Gaussian source with  $r_{\text{eff}} = 1.02 \pm 0.05$  ( $0.95 \pm 0.06$ ) fm used for the p- $\Xi^-$  (p- $\Omega^-$ ) emission

# Future Heavy Ion Detector

Possibility to extend heavy ion measurements at the LHC beyond 2030

arXiv:1902.01211



## Physics:

- Heavy flavor and quarkonia
- Low mass dielectrons
- Soft photons and hadrons
- BSM

## Design guidelines:

- All silicon detector
- High rate capability:  $\sim 10^{34}/\text{cm}^2/\text{s}$  ( $\sim 50\times$  Run3-4)
- Vertex spatial resolution:  $\sim 1 \mu\text{m}$
- Tracking over wide kinematic range
  - $30 \text{ MeV}/c < p_T < 10 \text{ GeV}/c$
  - $|\eta| < 4.0$