

# Neutrino signatures from magnetar remnants of BNS mergers

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Fermilab

Kavli Institute for Cosmological Physics (KICP),  
University of Chicago

Zoom seminar on the interface of particle physics and GWs  
January 13, 2026

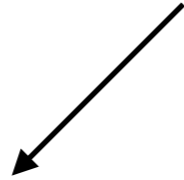
# Prologue

**New physics, understanding the fundamentals,....**

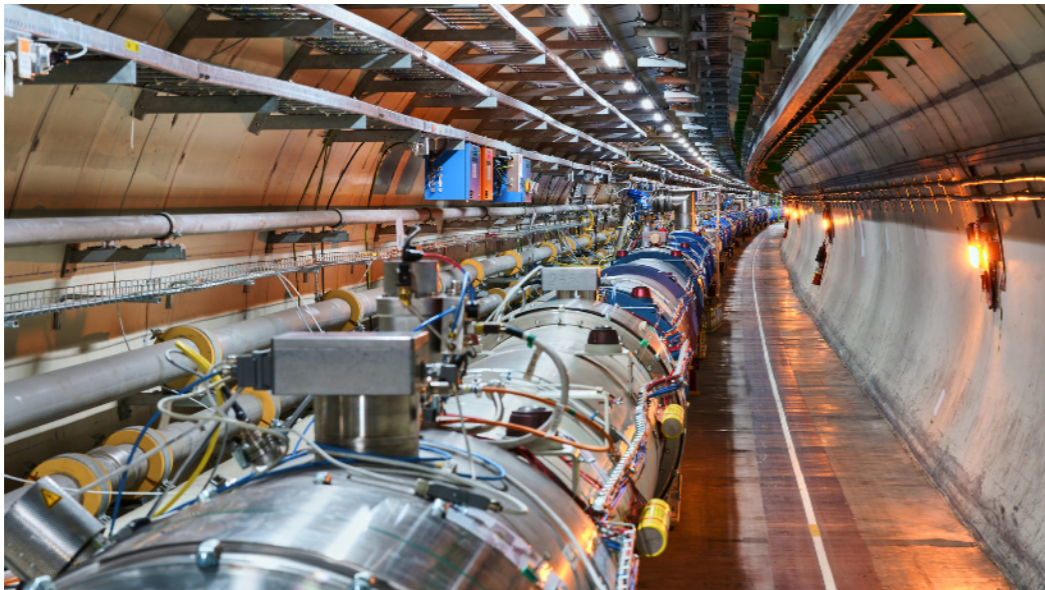
# Prologue

New physics, understanding the fundamentals,....

Man-made Accelerators



LHC



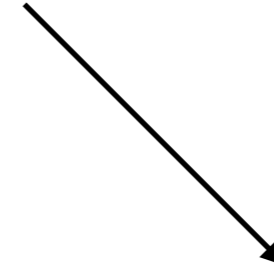
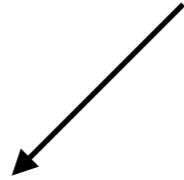
Tevatron



# Prologue

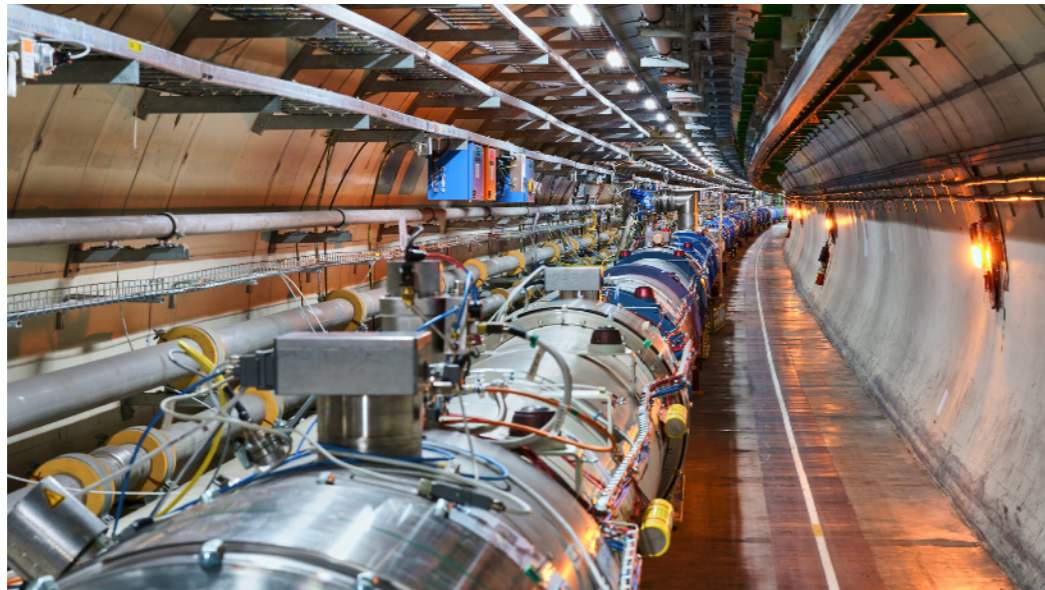
New physics, understanding the fundamentals,....

Man-made Accelerators

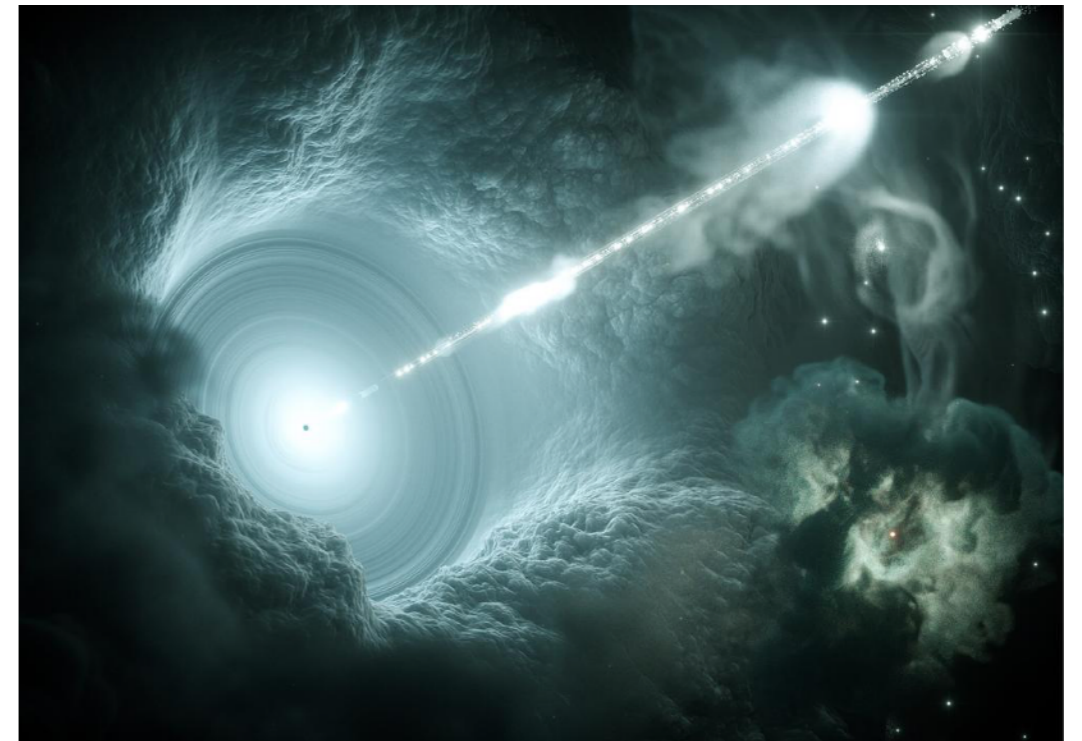


Cosmic Accelerators

LHC



Tevatron

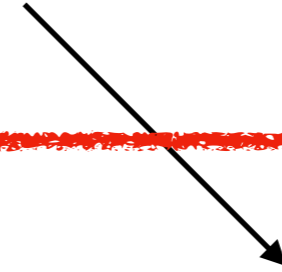
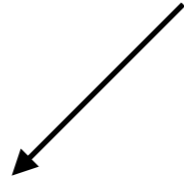


High-energy astrophysical phenomena

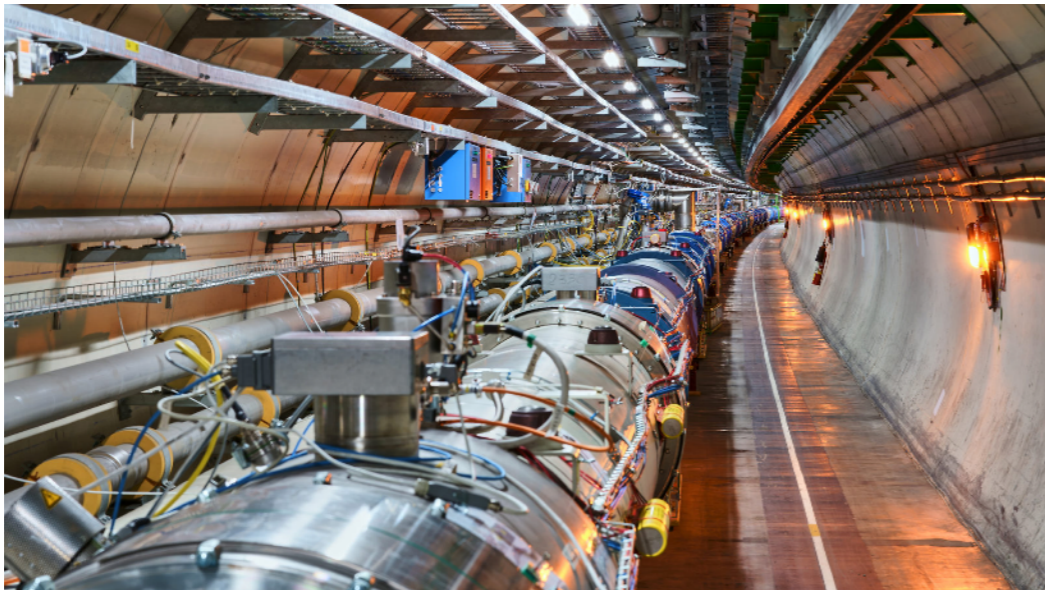
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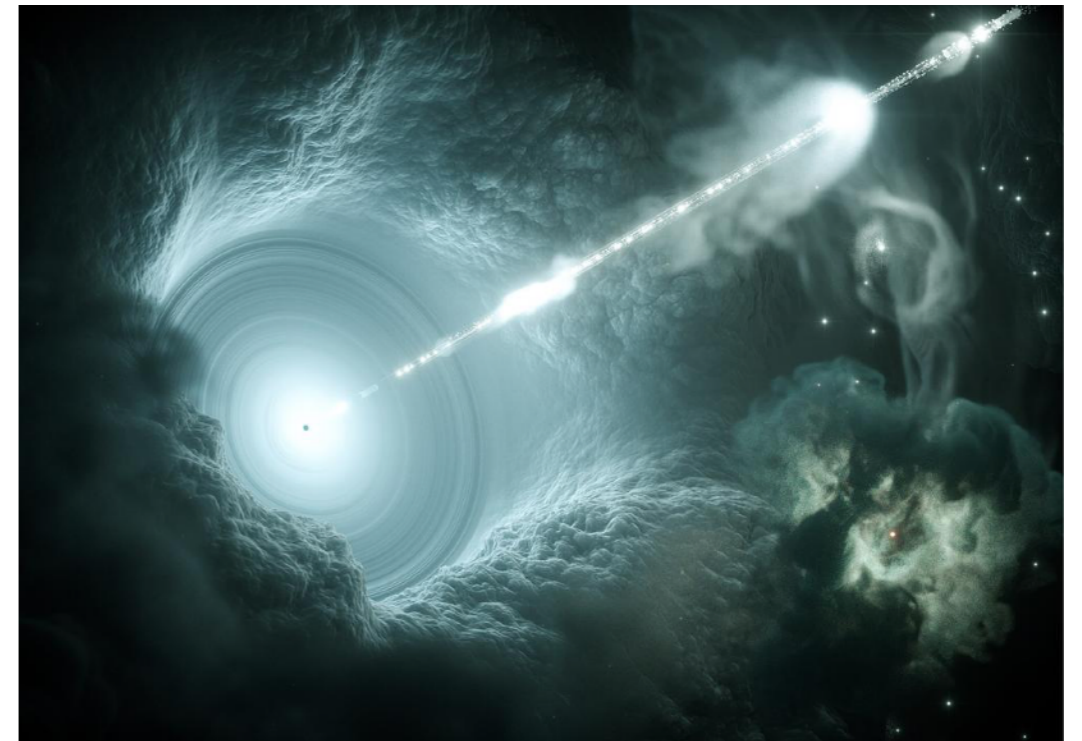
LHC



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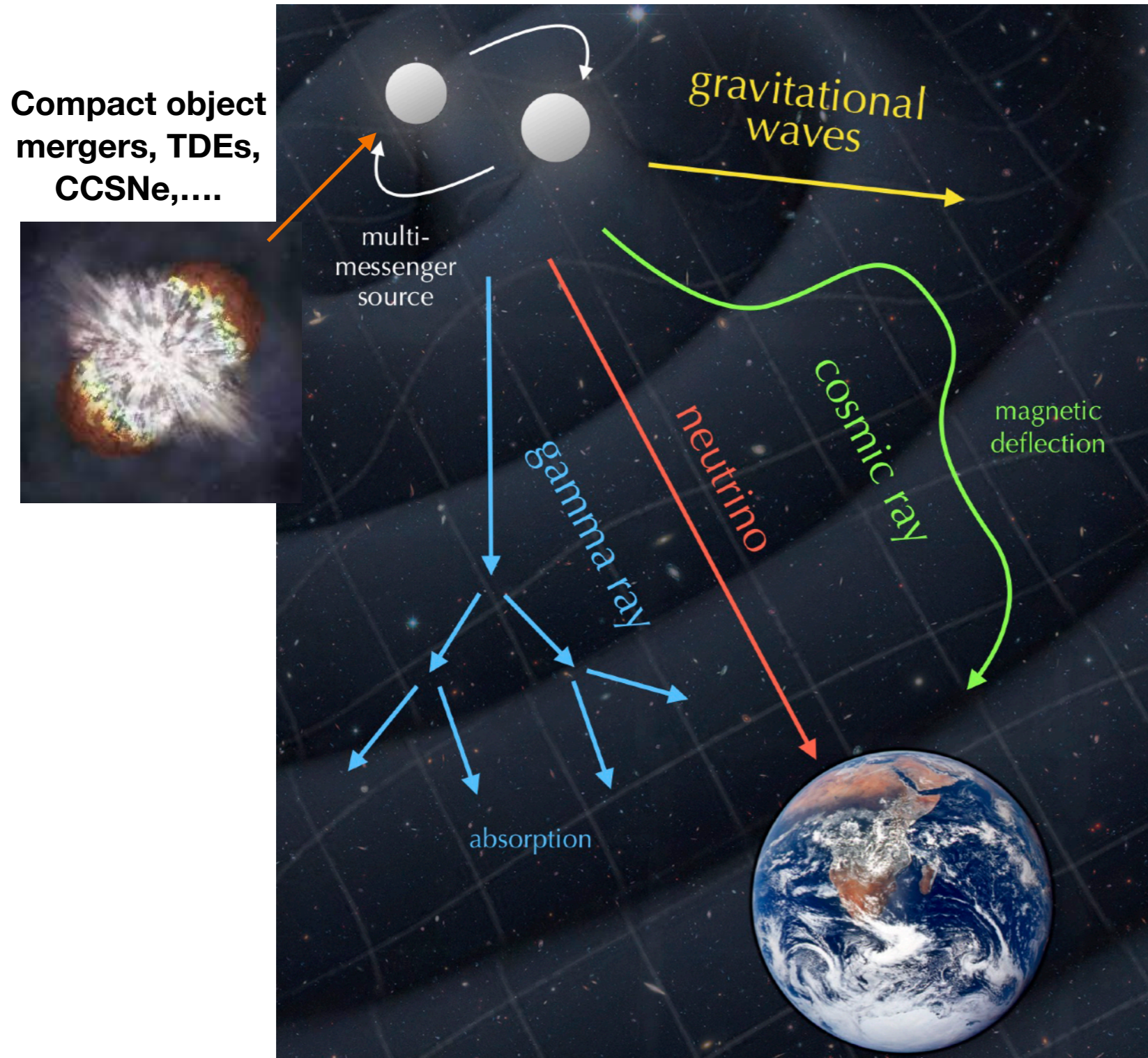


Cosmic Accelerators

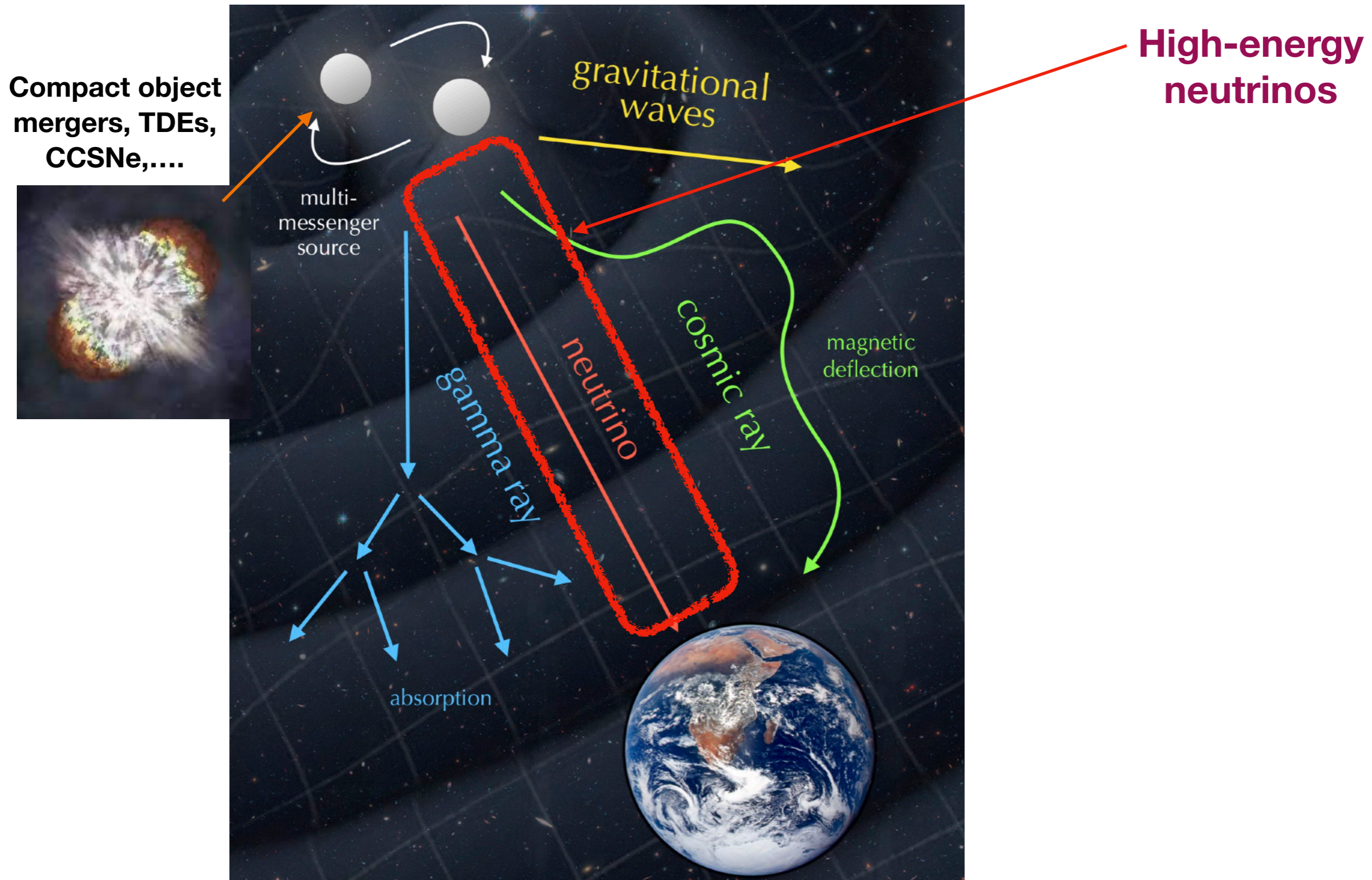


High-energy astrophysical phenomena

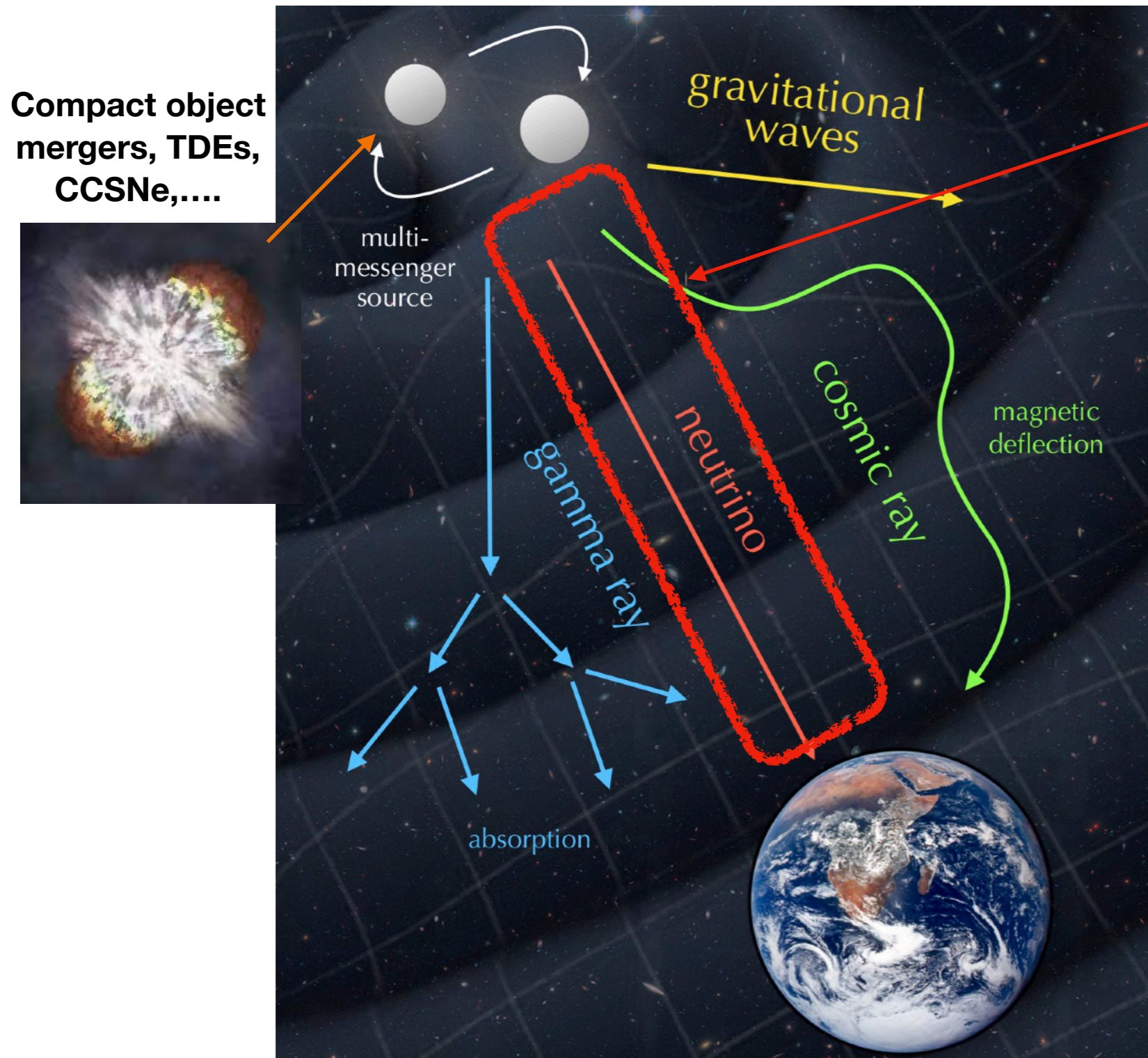
# The multi-messenger paradigm



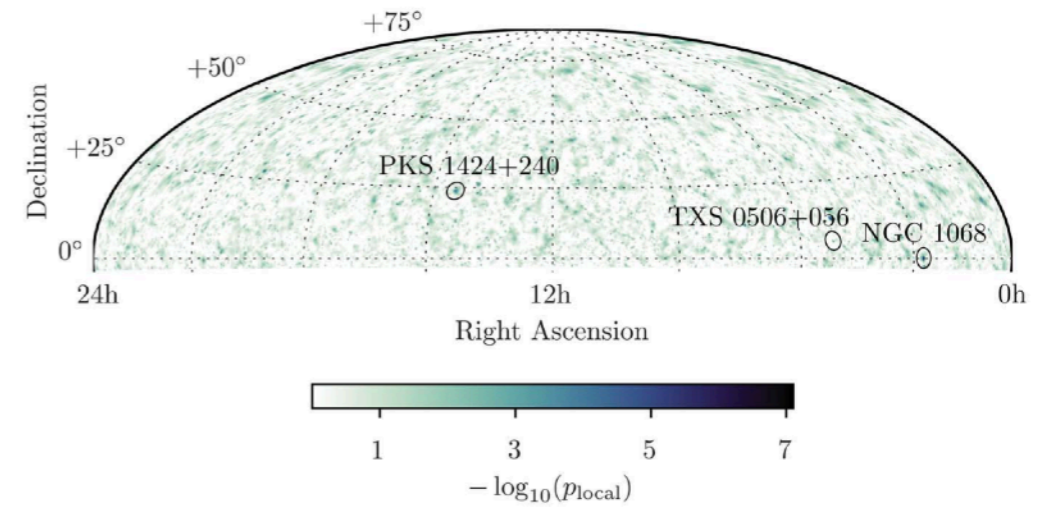
# The multi-messenger paradigm



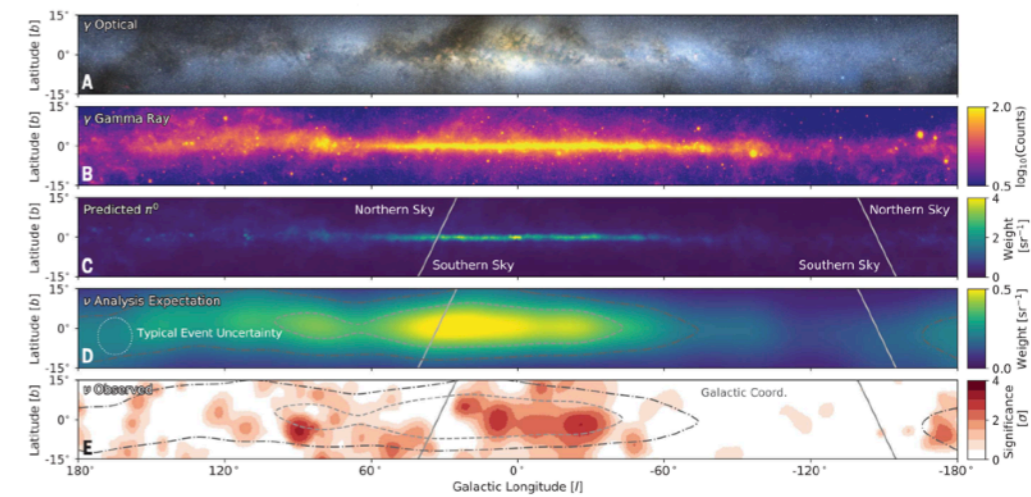
# The multi-messenger paradigm



## High-energy neutrinos



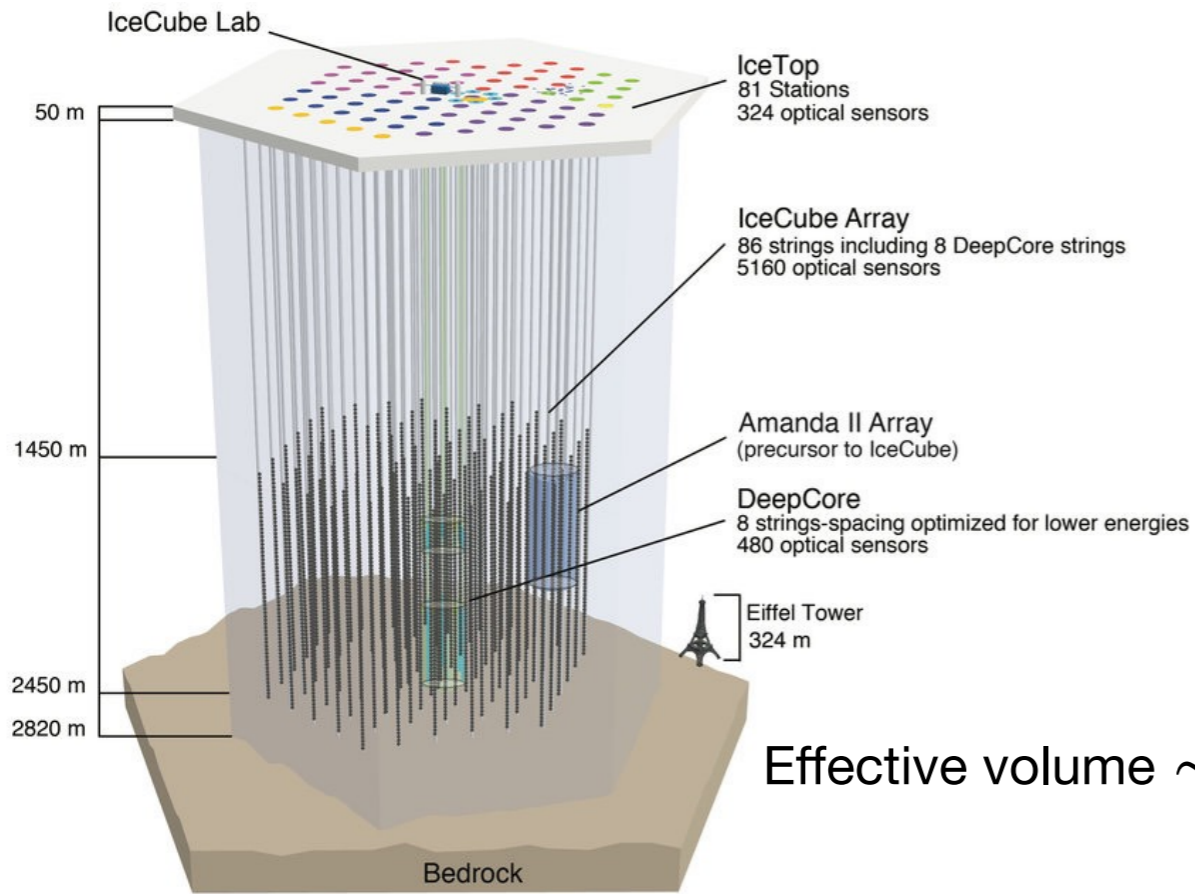
## Point Sources



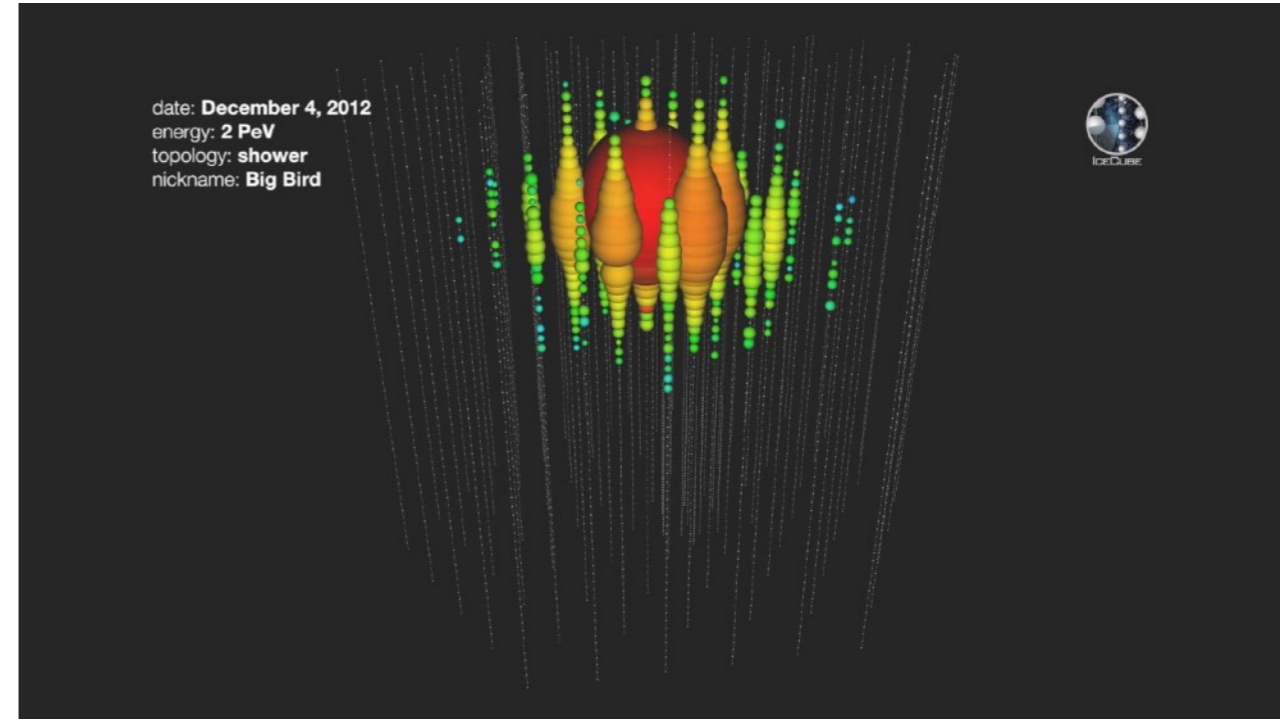
## The Galactic plane

Image credits: NBI  
IceCube Collab.+ Science 2022  
IceCube Collab.+ Science, 380, 2023

# High-energy neutrino detectors



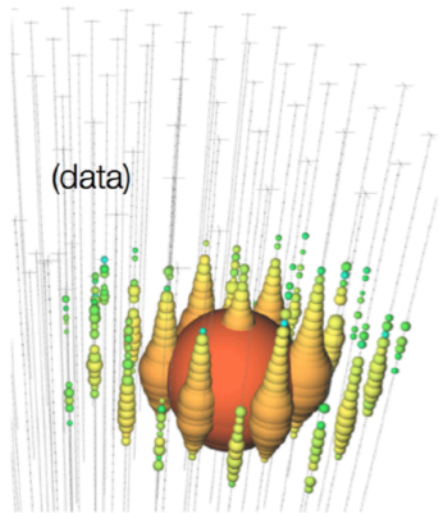
Effective volume  $\sim 1 \text{ km}^3$



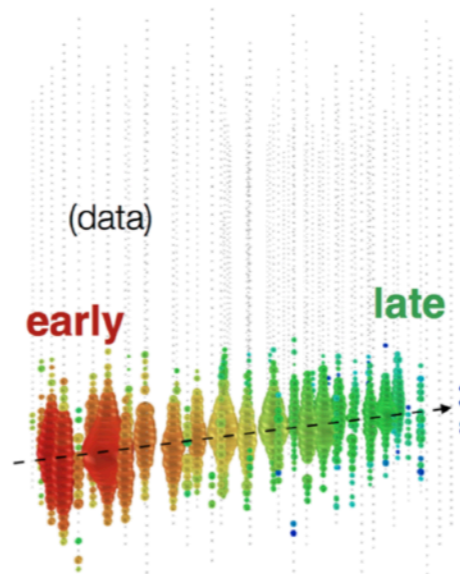
Neutral-current /  $\nu_e$

Charged-current  $\nu_\mu$

Charged-current  $\nu_\tau$



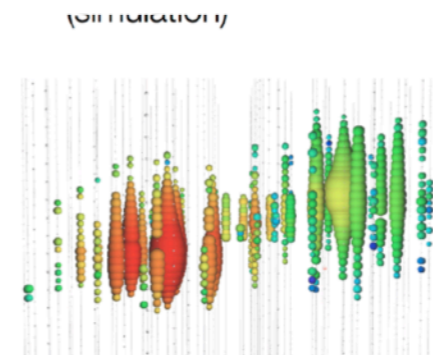
Isolated energy deposition (cascade) with no track



Up-going track

## IceCube observes seven astrophysical tau neutrino candidates

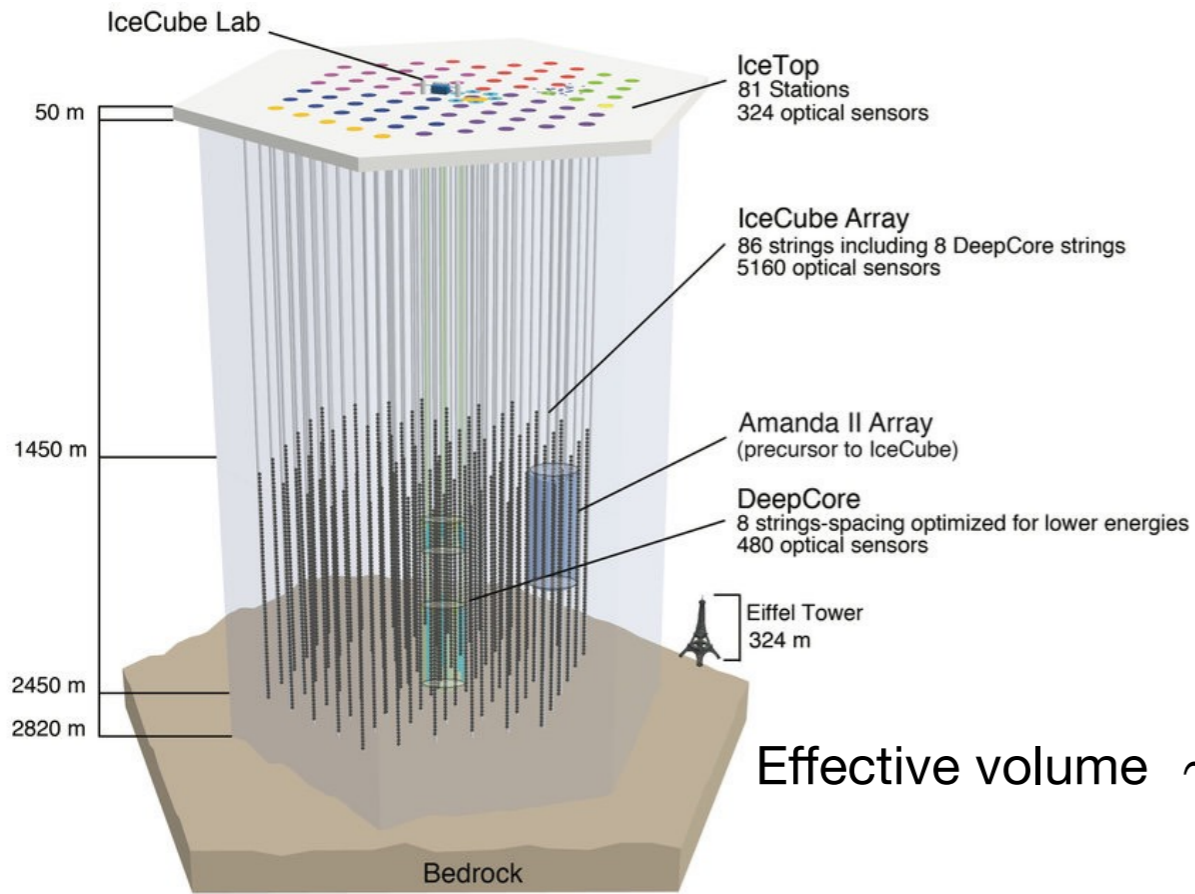
Posted on March 7, 2024 by Alisa King-Klemperer



Double cascade

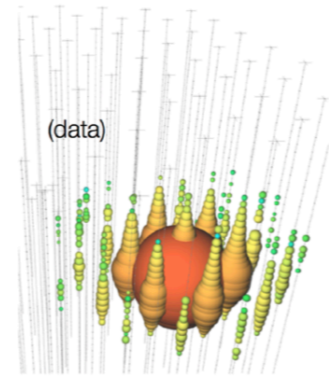
Image credits: [icecube.wisc.edu](http://icecube.wisc.edu)

# High-energy neutrino detectors



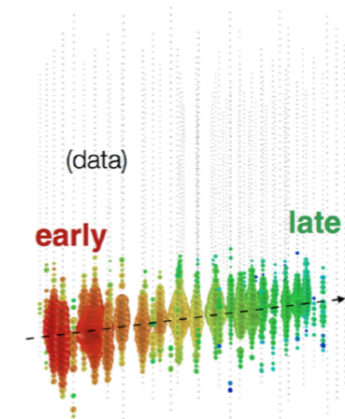
Effective volume  $\sim 1 \text{ km}^3$

Neutral-current /  $\nu_e$



Isolated energy deposition (cascade) with no track

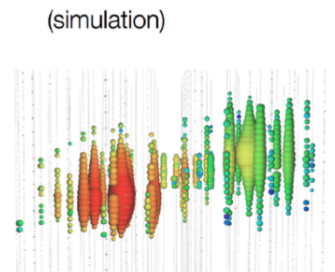
Charged-current  $\nu_\mu$



Up-going track

Charged-current  $\nu_\tau$

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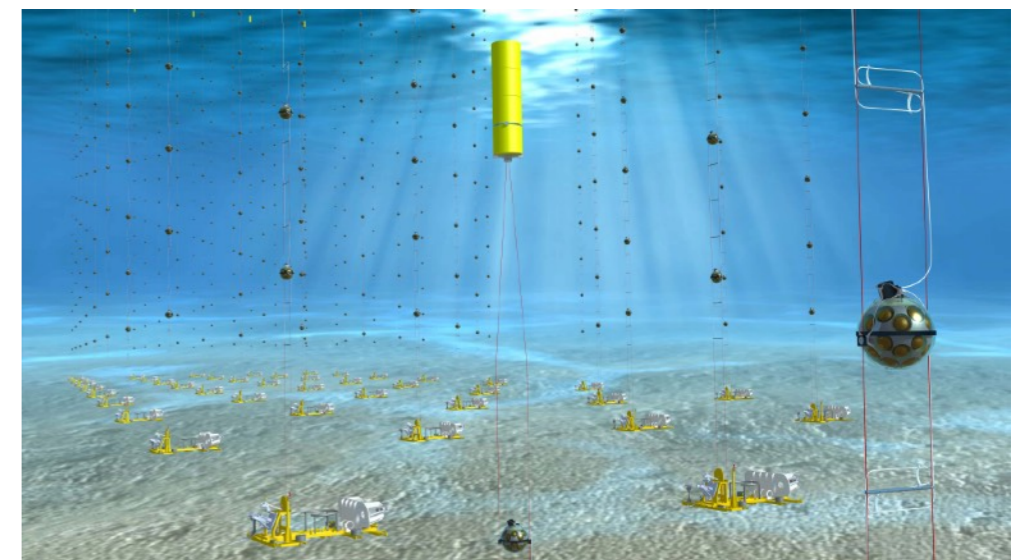


Double cascade

Baikal GVD

ANTARES

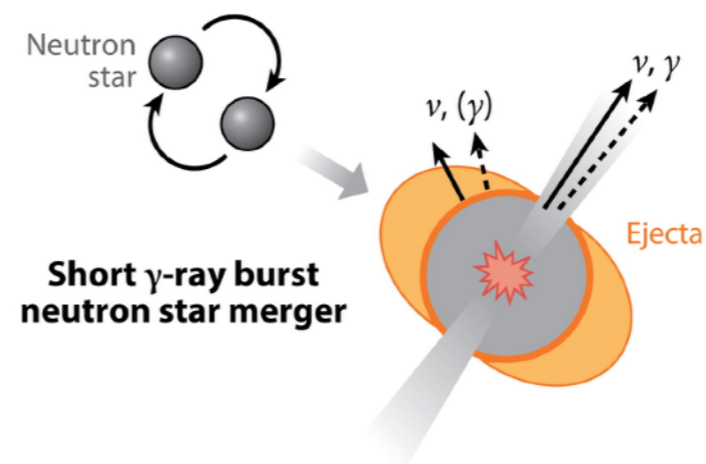
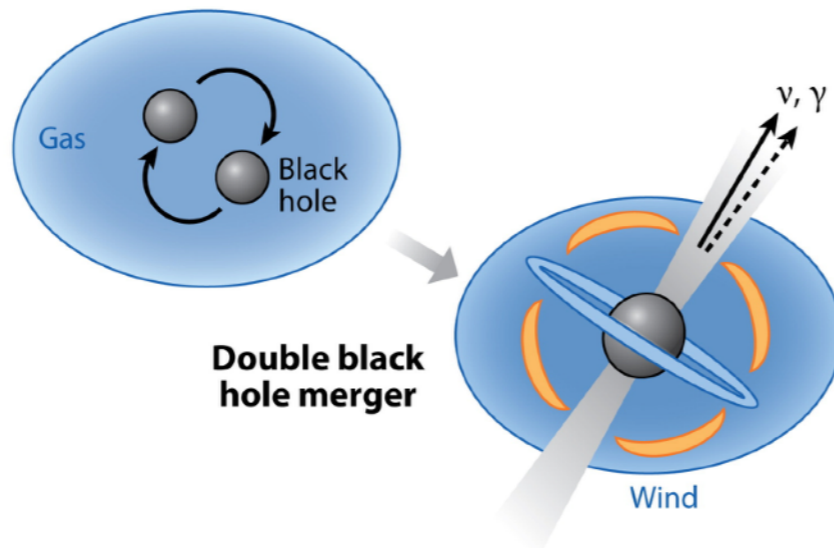
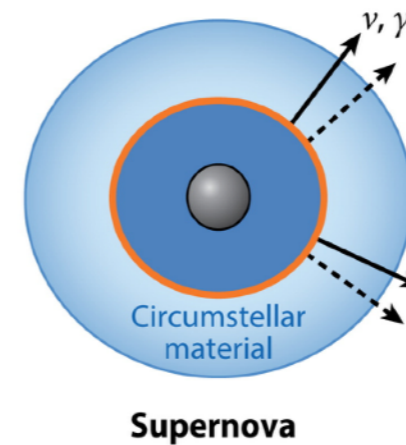
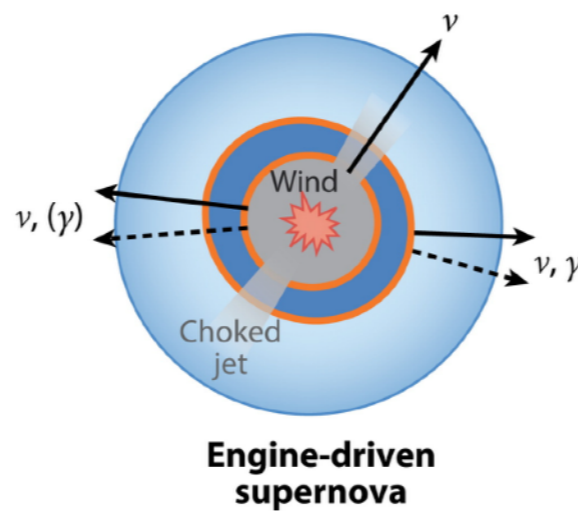
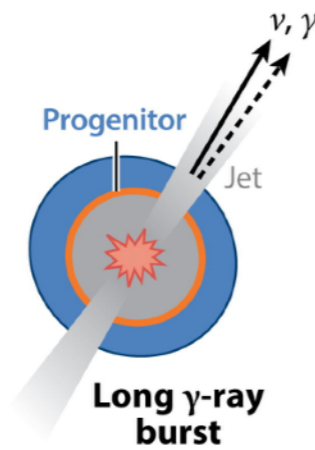
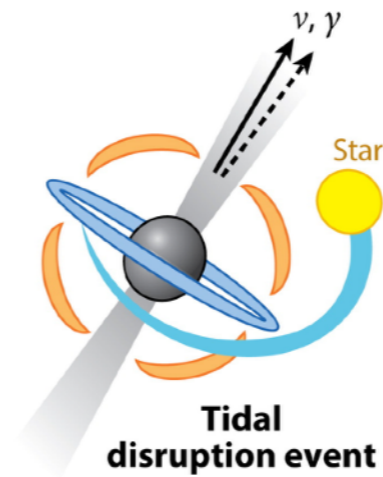
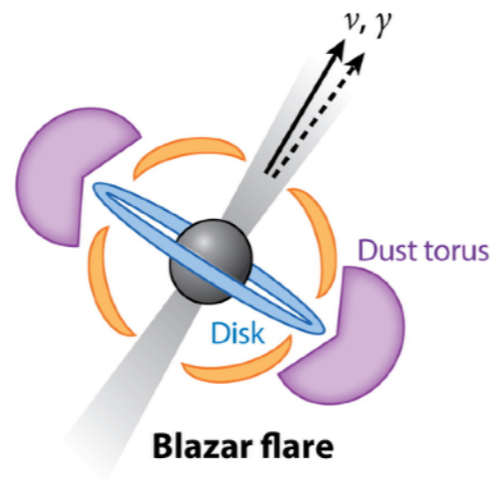
KM3NeT



Future detectors: IceCube-Gen2, RNO-G, GRAND, P-ONE....

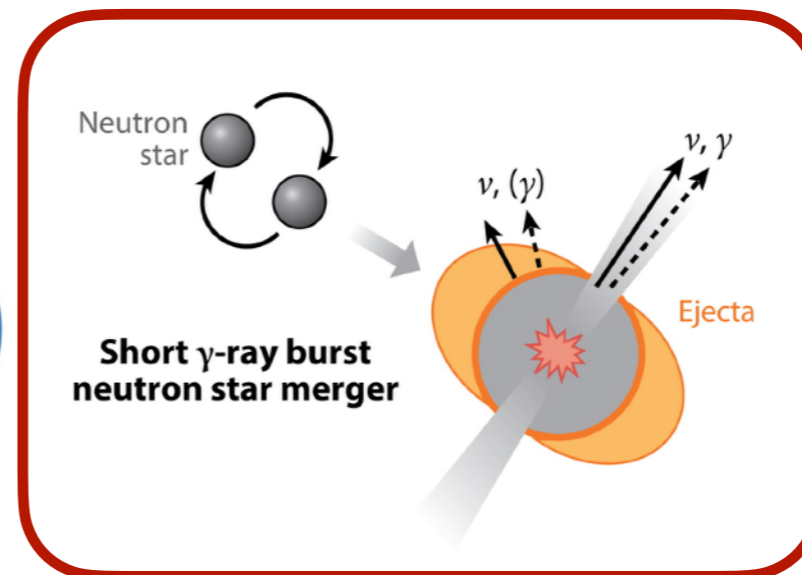
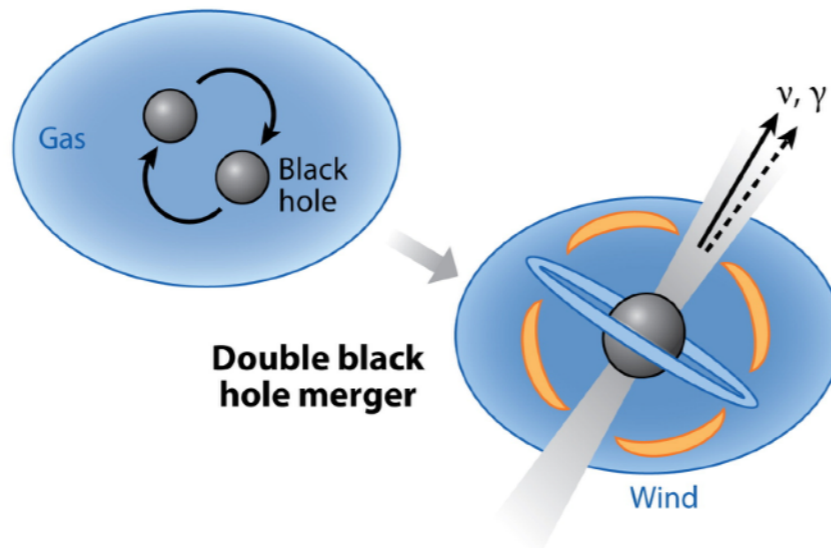
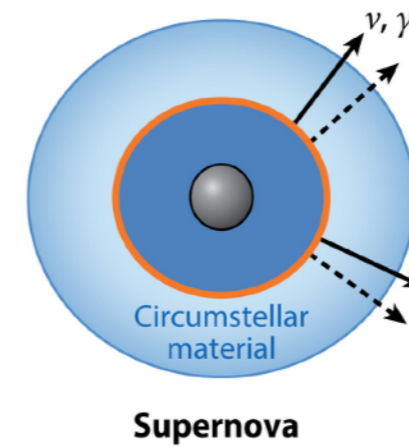
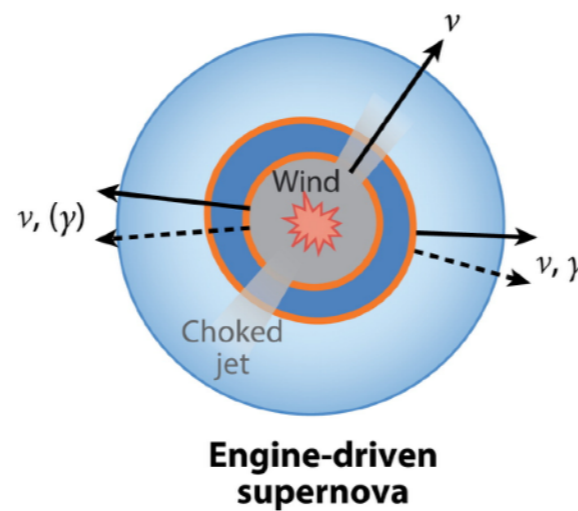
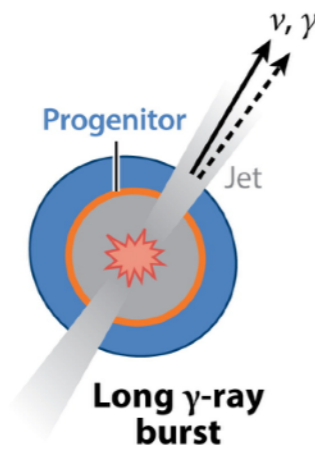
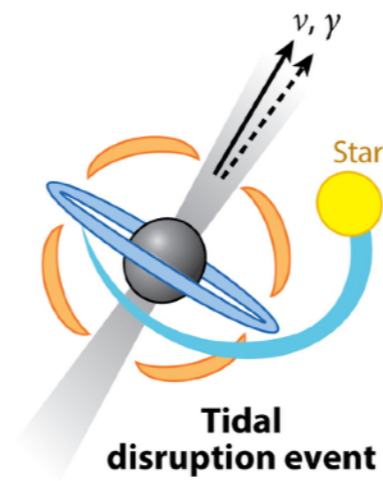
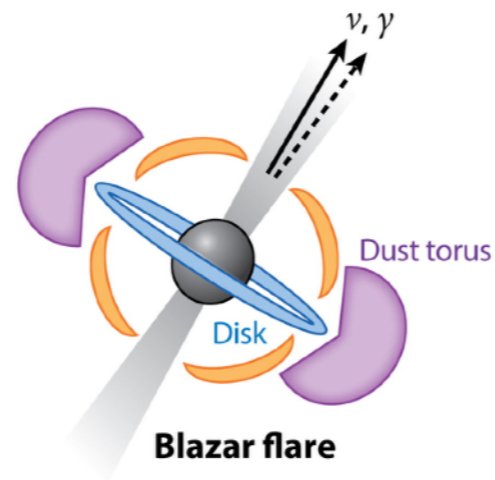
# The high-energy multi-messenger transients

Extreme astrophysical phenomena



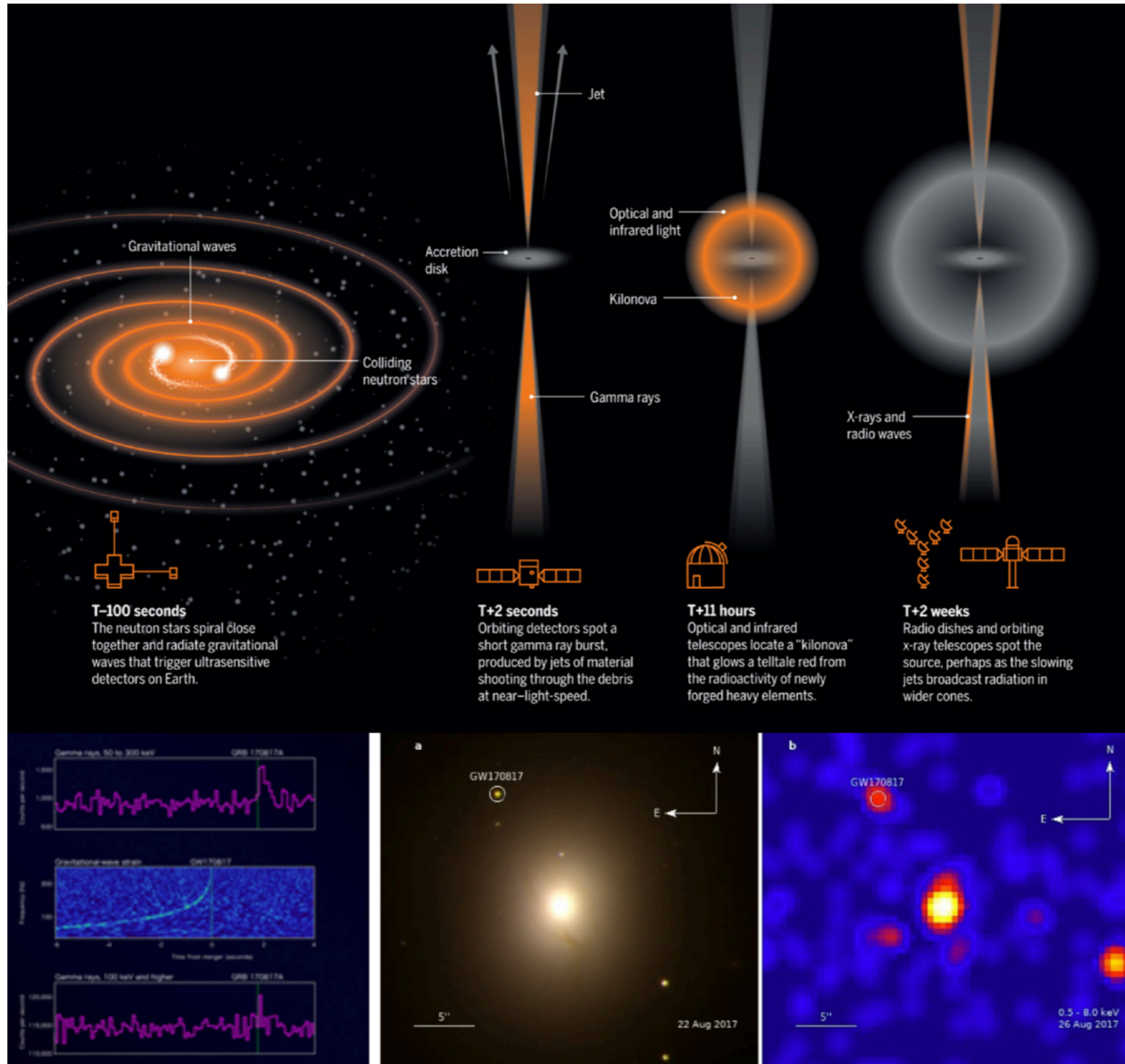
# The high-energy multi-messenger transients

High-energy astrophysical phenomena



# GW170817

~ 40 Mpc (NGC 4993)



No neutrinos :(

**Gamma rays**  
(Fermi+Integral)

**GW**  
(Adv. LIGO+Virgo)

**Optical**  
(HST)

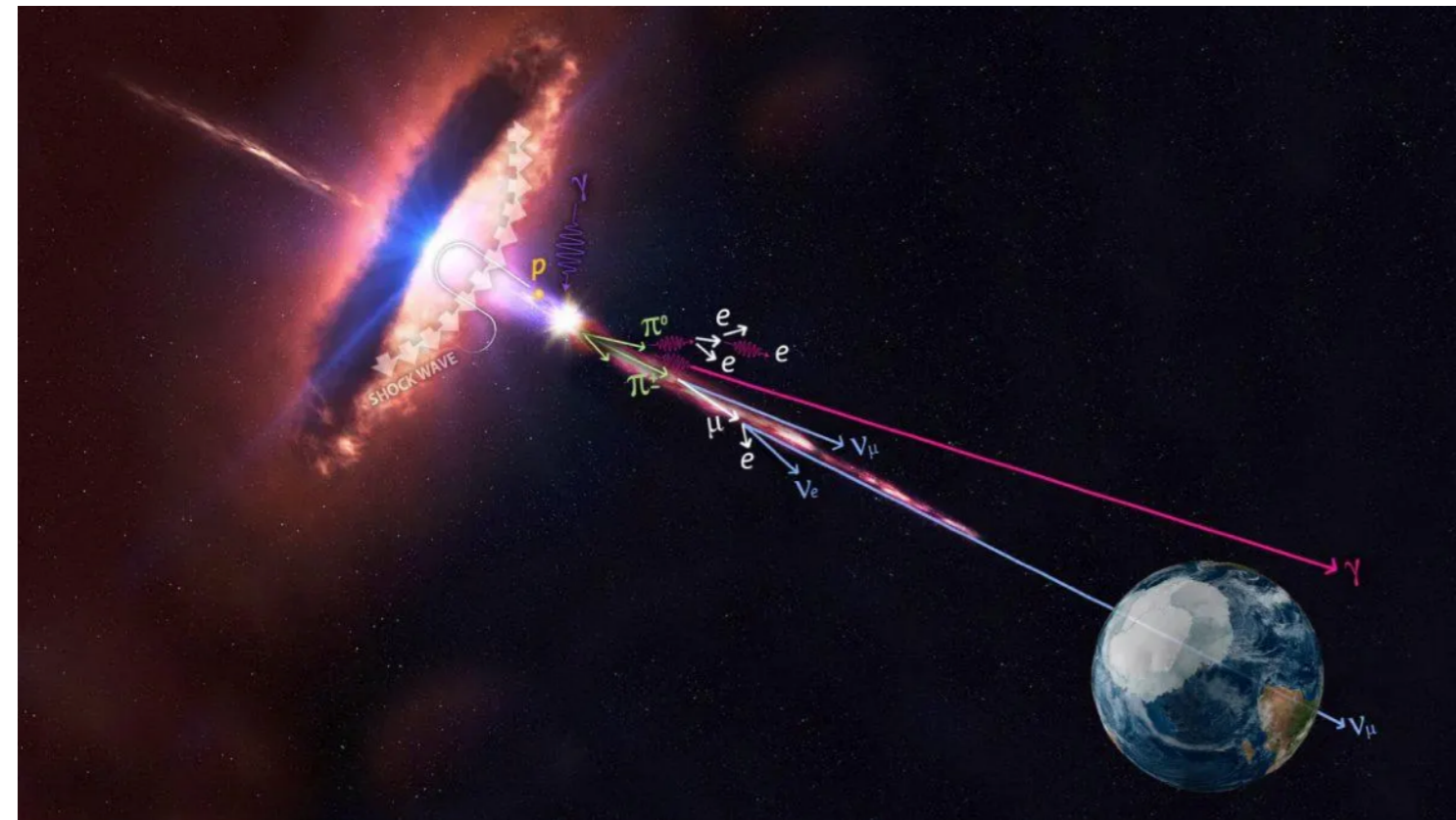
**X-rays**  
(Chandra)

Image credits: <https://ahead.iaps.inaf.it>

Abbott et al. 2017, ApJ 848, L13

Troja, Piro, van Earthen et al., 2017, Nature, 551, 71 13

# High-energy (HE) neutrinos



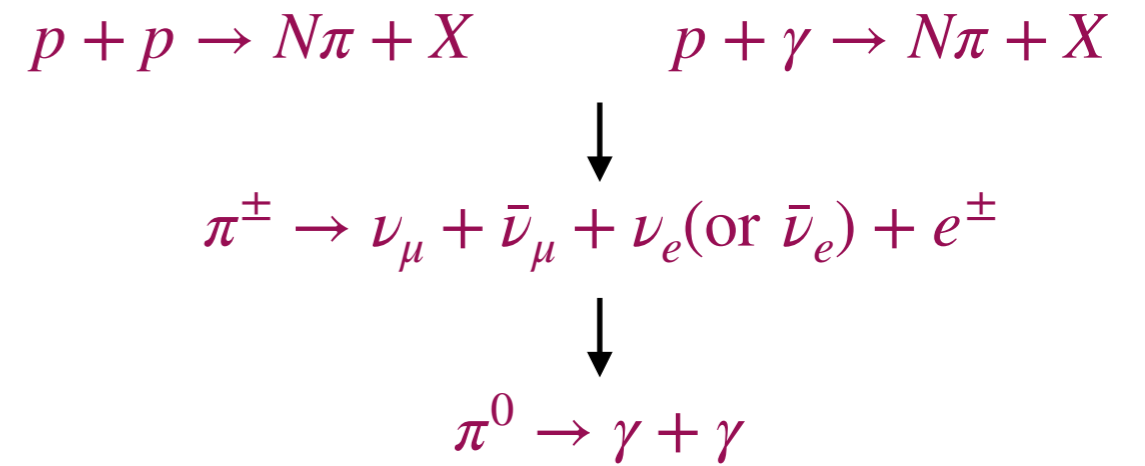
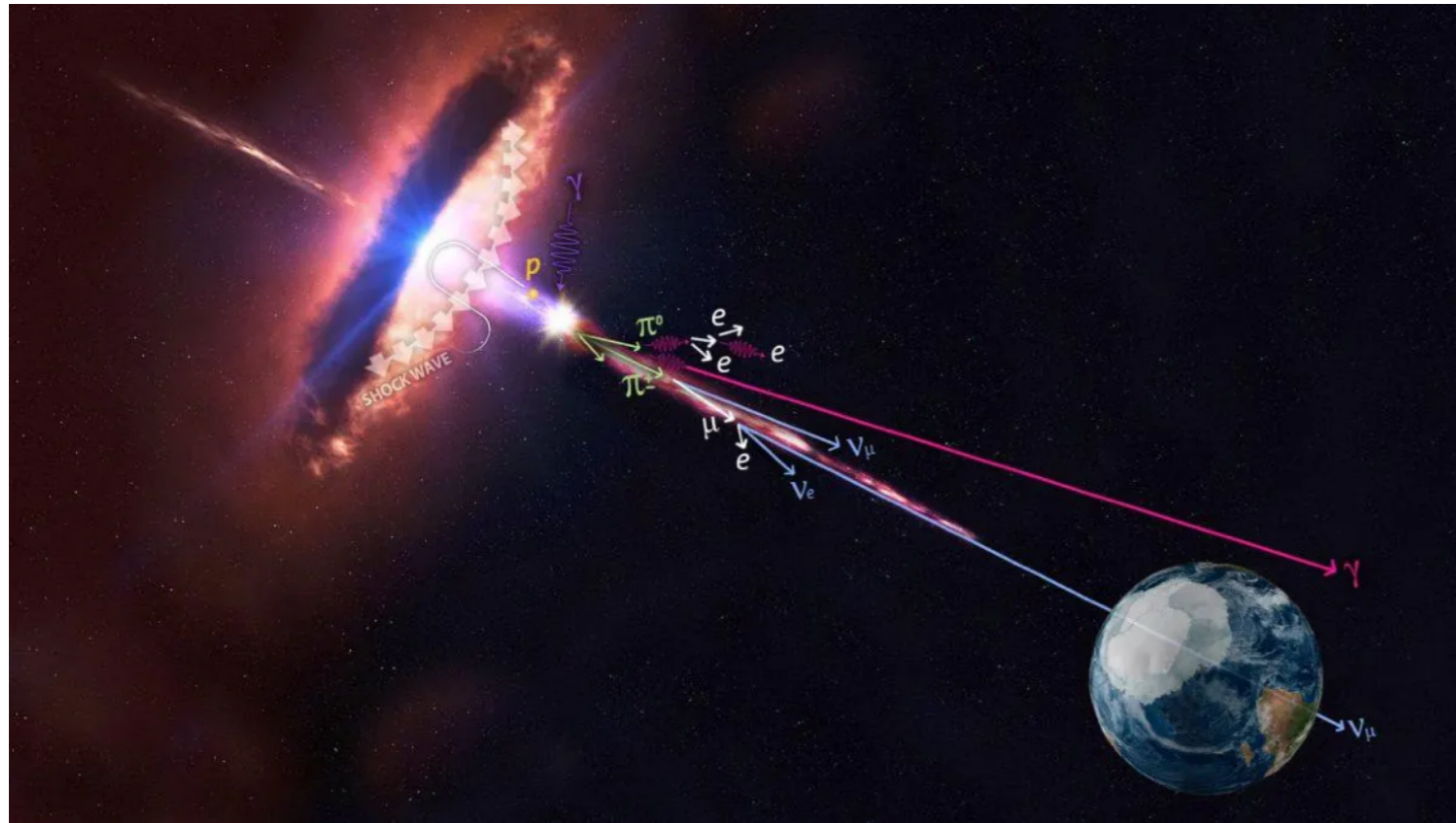
$$p + p \rightarrow N\pi + X$$

$$p + \gamma \rightarrow N\pi + X$$

$$\pi^{\pm} \rightarrow \nu_{\mu} + \bar{\nu}_{\mu} + \nu_e(\text{or } \bar{\nu}_e) + e^{\pm}$$

$$\pi^0 \rightarrow \gamma + \gamma$$

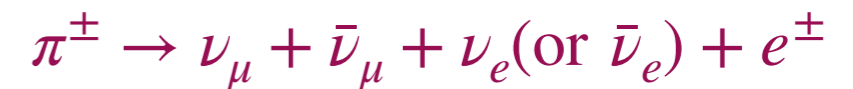
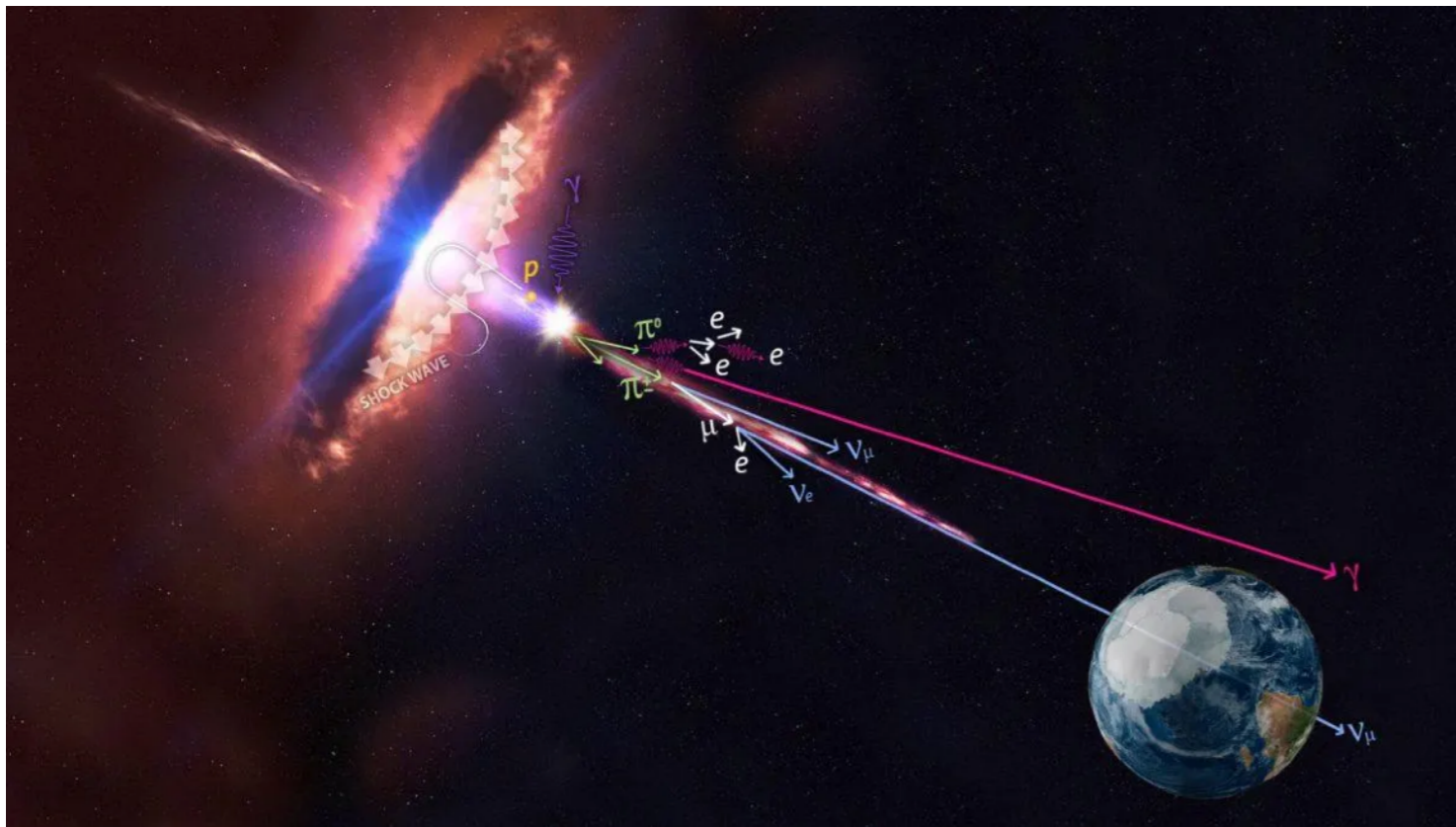
# High-energy (HE) neutrinos



## Conditions for HE- $\nu$ production:

- Acceleration of ions (p and nuclei) to sufficiently high energies - Shocks, magnetic reconnection, stochastic acceleration aided by turbulence
- Rate of acceleration  $>$  Rate of energy loss

# High-energy (HE) neutrinos



Proton energy loss due to p-p interactions

## Conditions for HE- $\nu$ production:

- Acceleration of ions (p and nuclei) to sufficiently high energies - Shocks, magnetic reconnection, stochastic acceleration aided by turbulence
- Rate of acceleration  $>$  Rate of energy loss
- Significant density on target media - matter and radiation
- (a) and (b)  $\rightarrow$  production of charged mesons - pions that decay into neutrinos, charged leptons, and gamma-rays

$$t_{pp}^{-1} = n_N \kappa_{pp} \sigma_{pp} c$$

Nucleon density  $\rightarrow$   $n_N$       Proton inelasticity  $\rightarrow$   $\kappa_{pp}$       p-p cross-section  $\rightarrow$   $\sigma_{pp}$

$$t_{p\gamma}^{-1}(\epsilon_p) = \frac{c}{2\gamma_p^2} \int_{\bar{\epsilon}_{th}}^{\infty} d\bar{\epsilon} \kappa_{p\gamma}(\bar{\epsilon}) \sigma_{p\gamma}(\bar{\epsilon}) \bar{\epsilon} \int_{\bar{\epsilon}/2\gamma_p}^{\infty} d\epsilon \epsilon^{-2} n_\epsilon$$

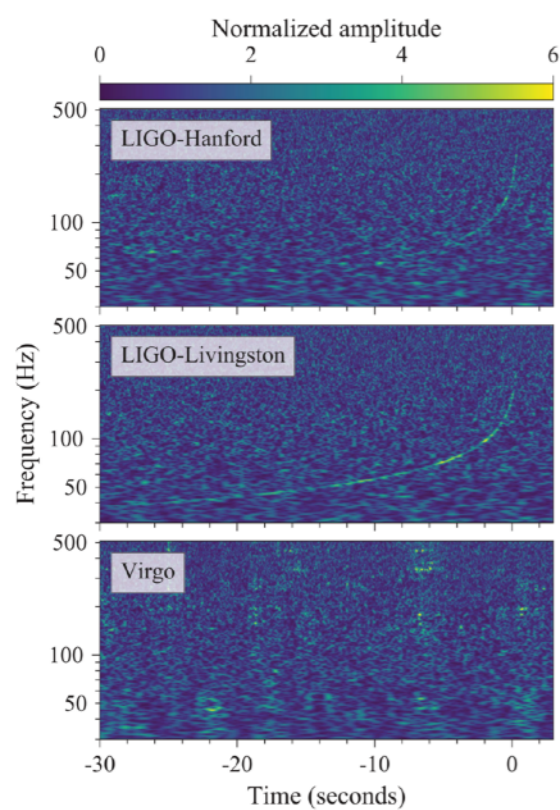
Proton energy  $\rightarrow$   $\epsilon_p$       Photon energy in proton rest frame  $\rightarrow$   $\bar{\epsilon}$       p- $\gamma$  cross-section  $\rightarrow$   $\sigma_{p\gamma}$

Proton energy loss due to p- $\gamma$  interactions

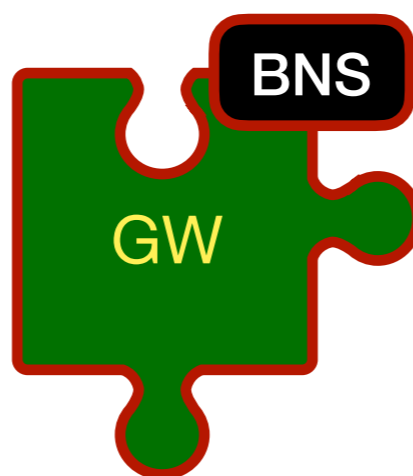
# BNS mergers: particle accelerators and multi-messenger zoo

**BNS**

# BNS mergers: particle accelerators and multi-messenger zoo



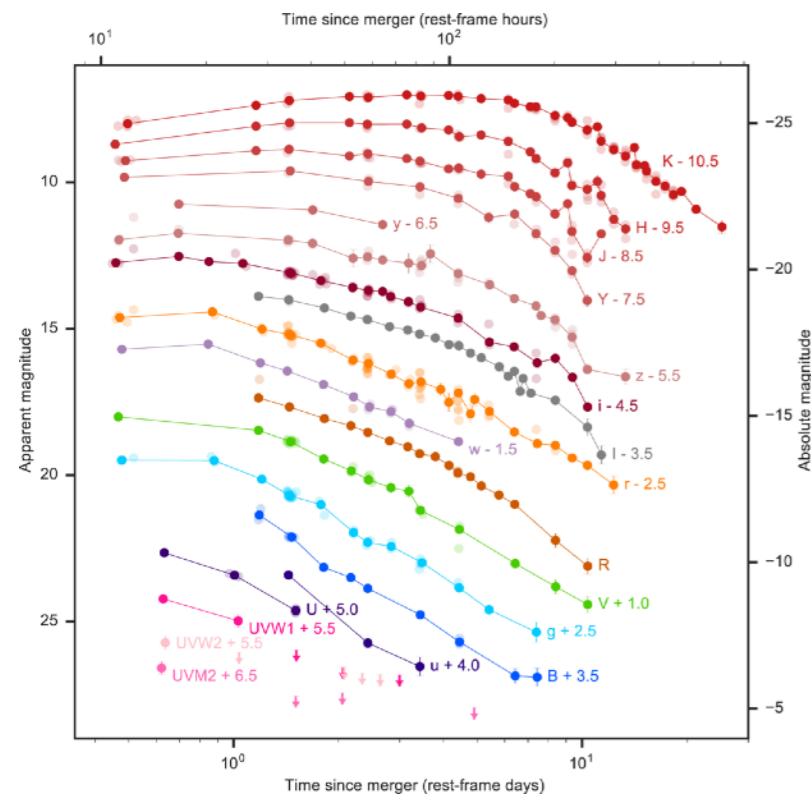
**Observed**



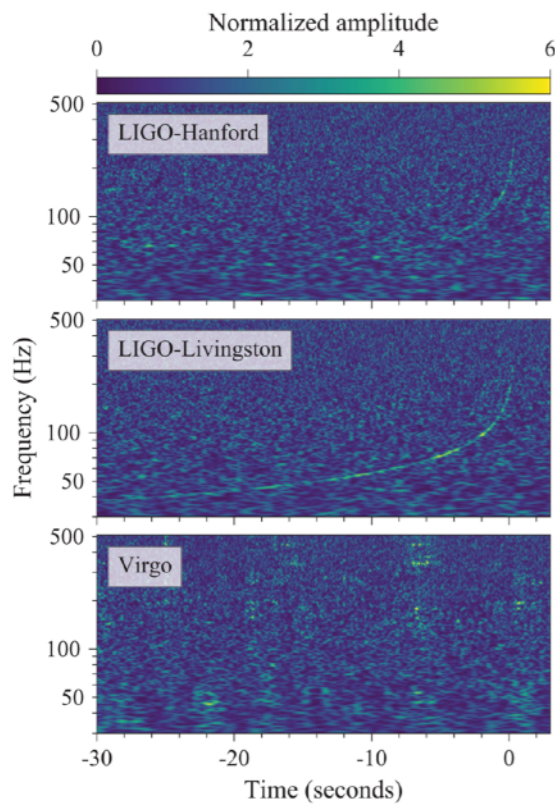
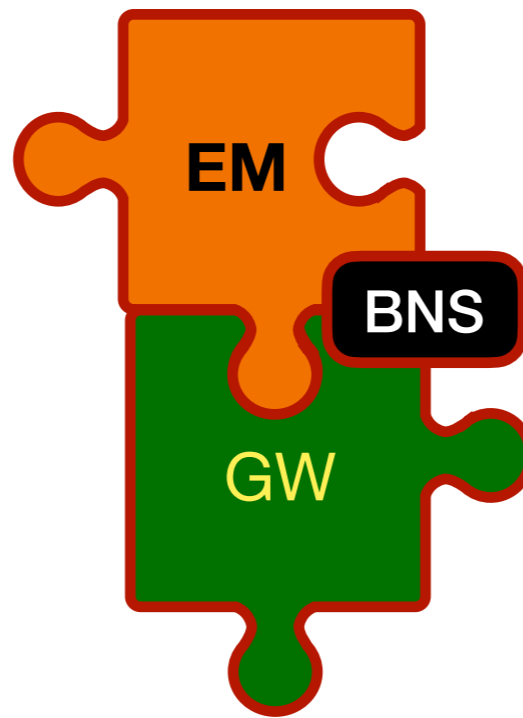
*S. Gezari, Annu. Rev. Astron. Astrophys. 2021. 59:21–58*  
*Kimura+, PRD (2018), Fang & Metzger (2017)*  
*Mukhopadhyay & Kimura (2024)*  
*LIGO Collab (2017)*

# BNS mergers: particle accelerators and multi-messenger zoo

## Observed



Kilonova emission  
Afterglow emission  
Short GRB

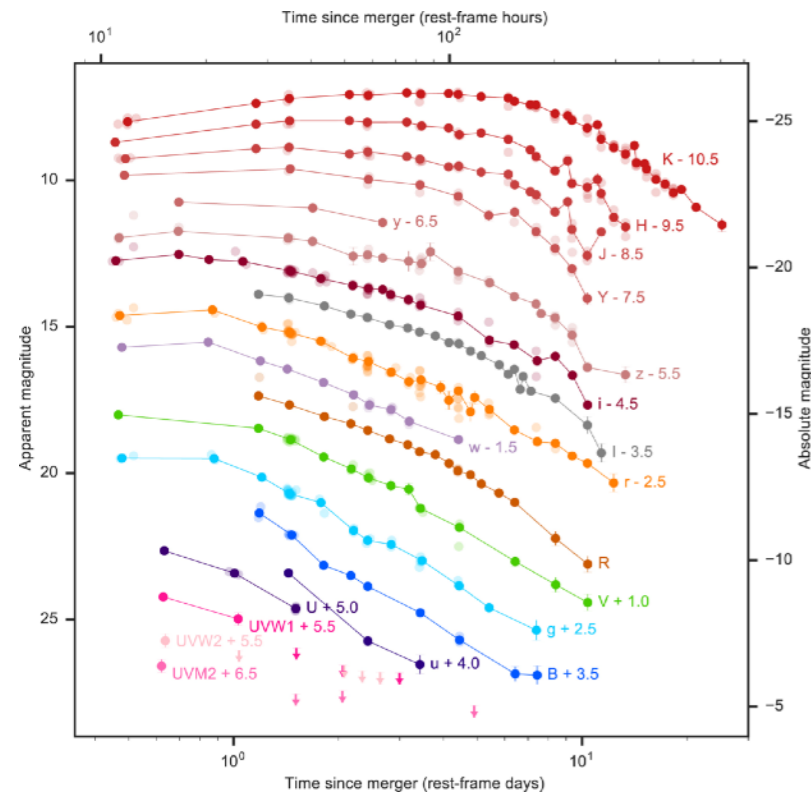


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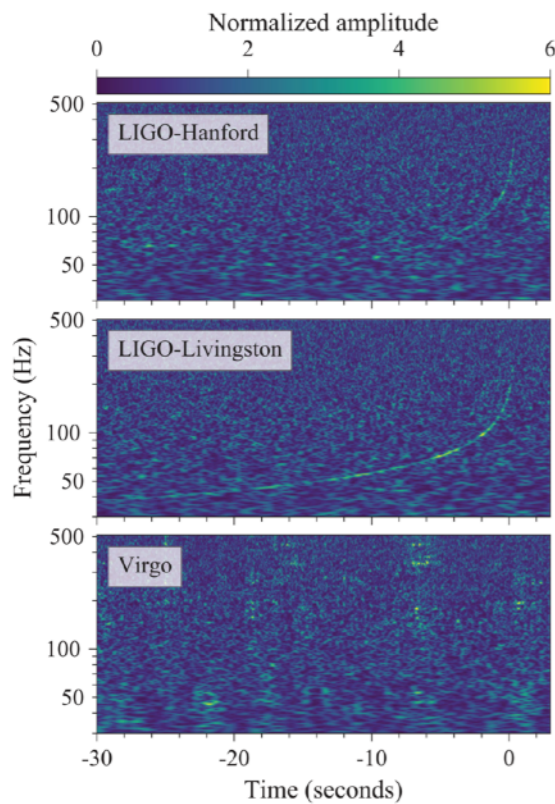
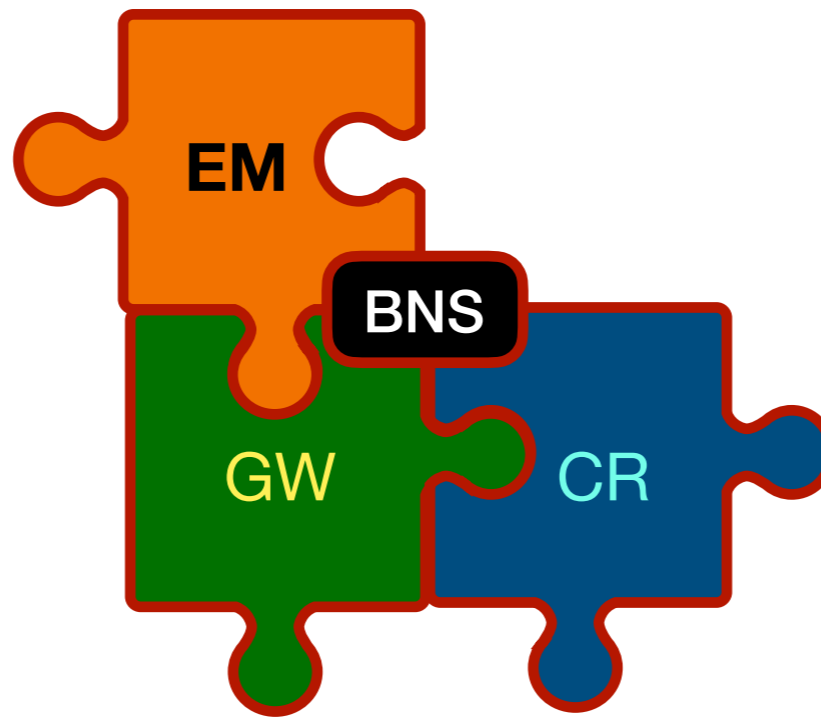
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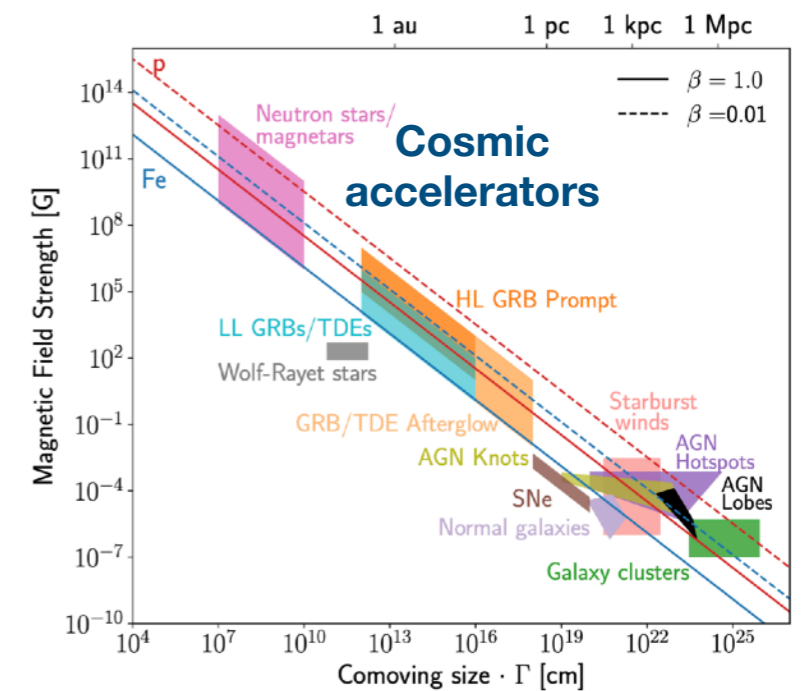
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Kilonova emission  
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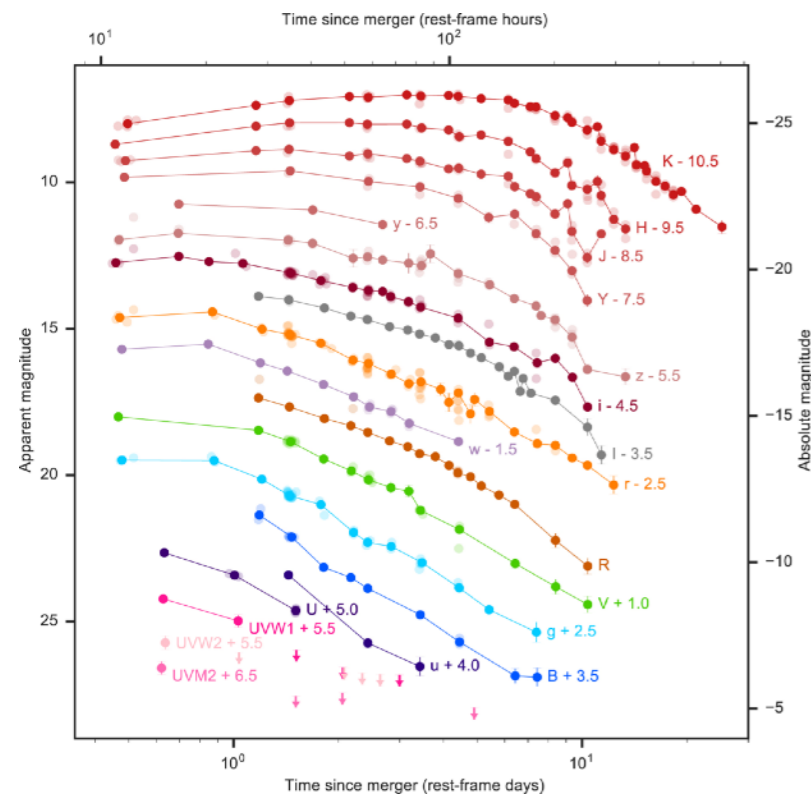
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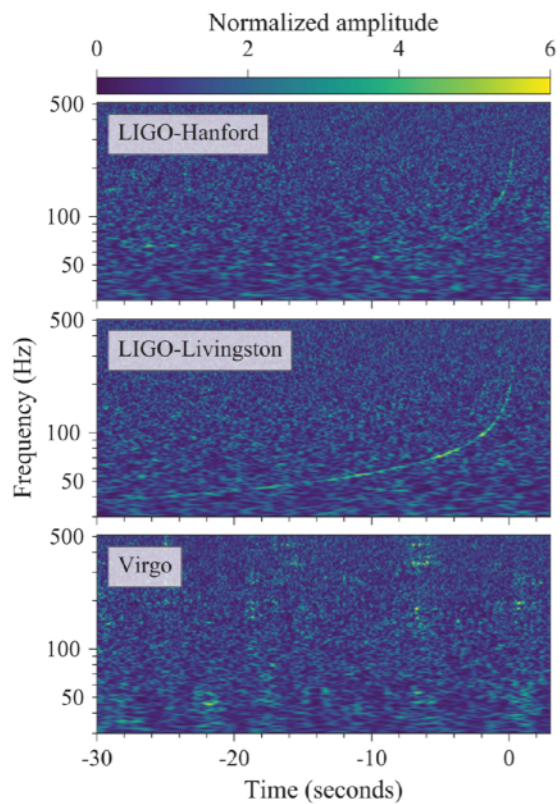
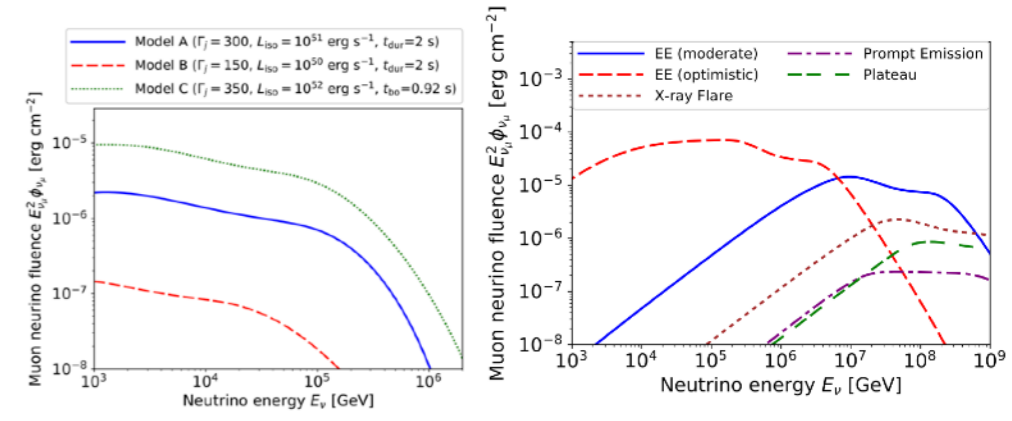
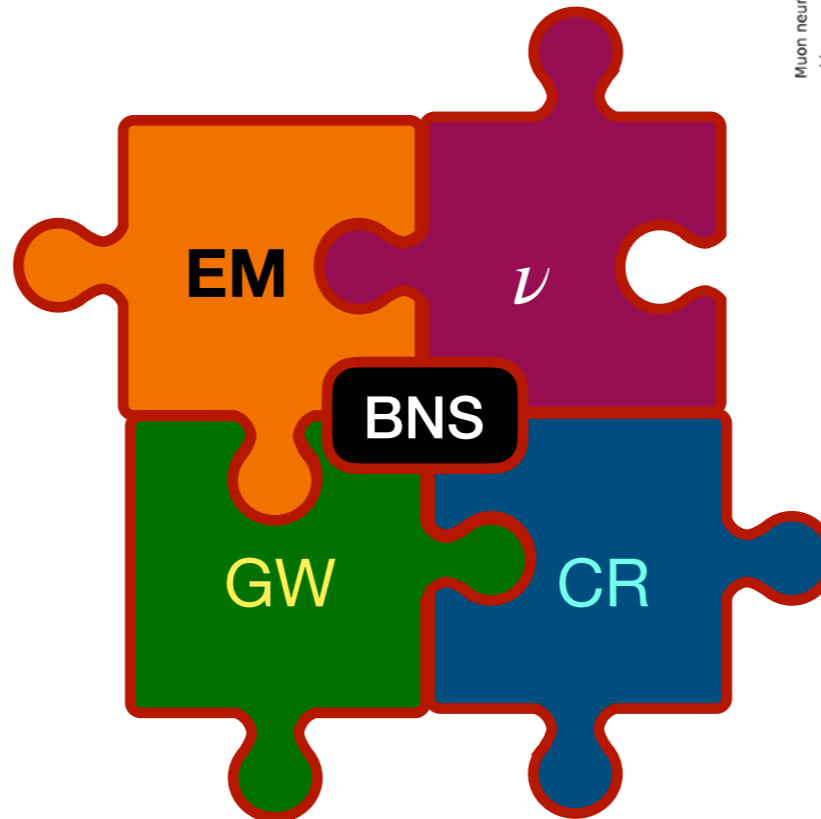
S. Gezari, *Annu. Rev. Astron. Astrophys.* 2021. 59:21–58  
 Kimura+, *PRD* (2018), Fang & Metzger (2017)  
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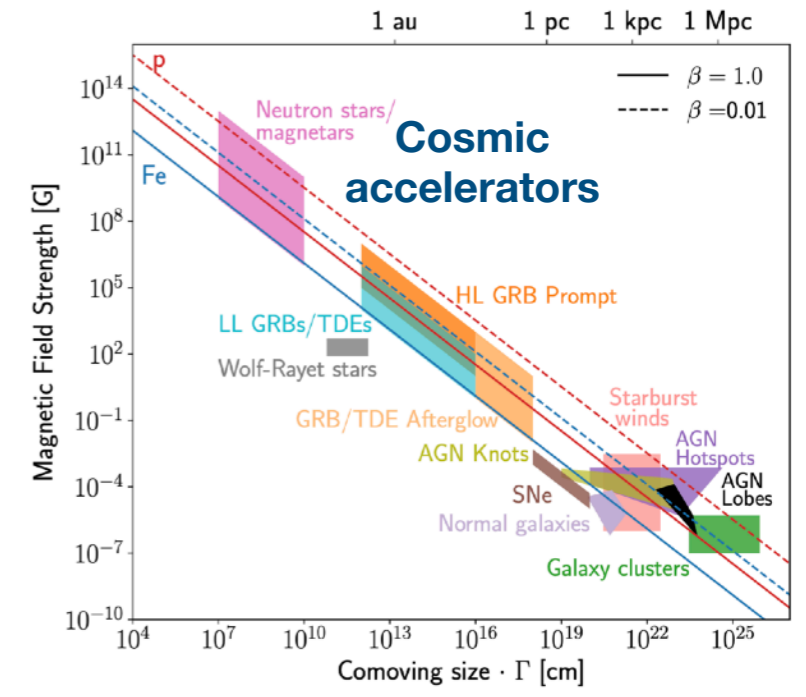
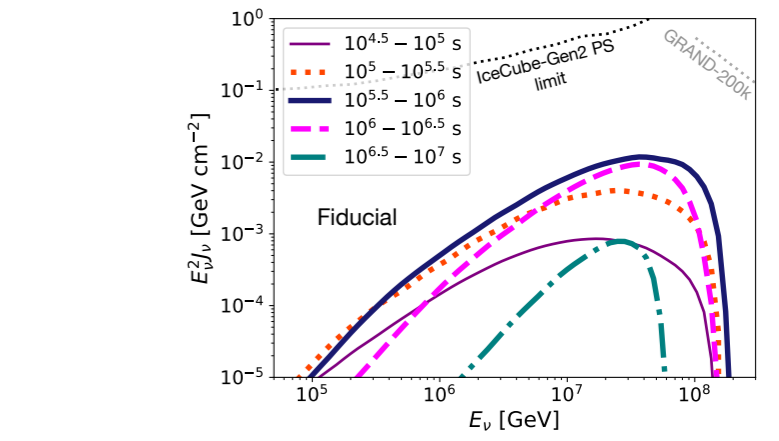
## Observed



Kilonova emission  
Afterglow emission  
Short GRB



## Observed



Batista et al., *Front. Astron. Space Sci.* 6 (2019), 23  
 Kimura+, *PRD* (2018), Fang & Metzger (2017)  
 Mukhopadhyay et al. (2024)  
 LIGO Collab (2017)

## Based on

### High-energy neutrino and EM emissions from magnetars

Based on: **High-energy neutrino signatures from pulsar remnants of binary neutron-star mergers: coincident detection prospects with gravitational waves**

MM, S.S. Kimura, B.D. Metzger

[Astrophys. J. 987 \(2025\) 2 \(arXiv: 2407.04767\)](#)

Electromagnetic signatures from pulsar remnants of binary neutron-star mergers

MM, S.S. Kimura

[Astrophys. J. Lett. 989 \(2025\) 2, L41 \(arXiv: 2506.09157\)](#)

Hunting for high-energy and ultrahigh energy neutrinos from BNS mergers at next-generation GW and neutrino detectors

Based on: **Gravitational wave triggered high energy neutrino searches from BNS mergers: prospects for next generation detectors**

MM, S. S. Kimura, K. Murase

[Phys. Rev. D 109, 4, 043053 \(2024\) \(arXiv: 2310.16875\)](#)

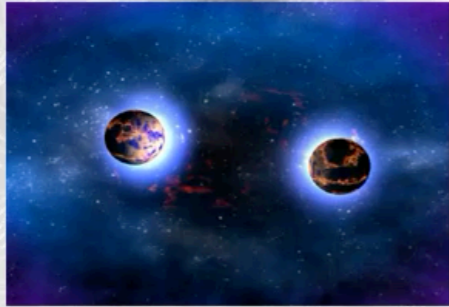
Ultrahigh energy neutrino searches using next-generation gravitational wave detectors at radio neutrino detectors: GRAND, IceCube-Gen2 Radio, and RNO-G

MM, K. Kotera, S. Wissel, K. Murase, S.S. Kimura

[Phys. Rev. D 110, 6, 063004 \(arXiv: 2406.19440\)](#)

# Fate of NS-NS mergers

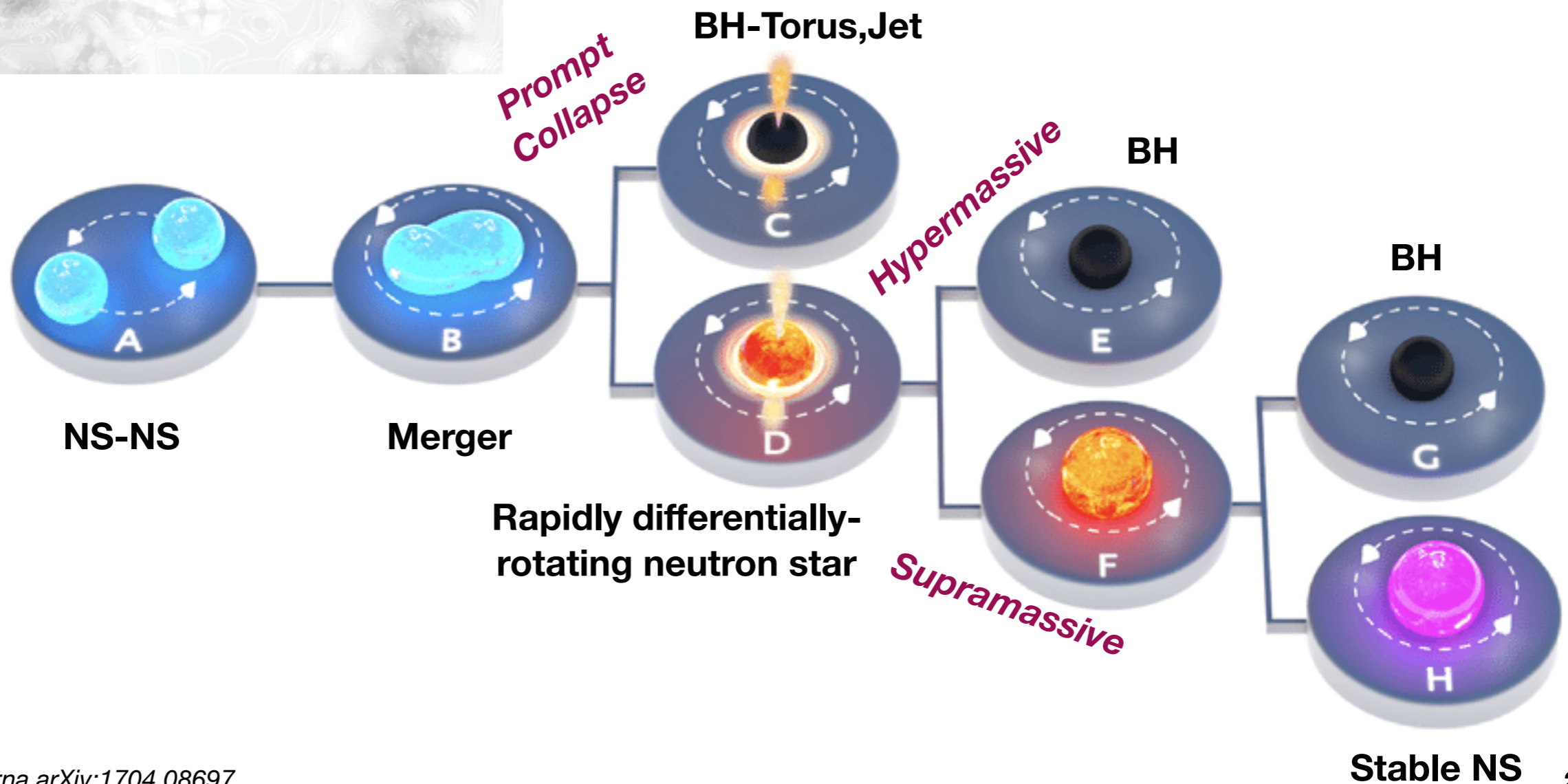
SWIFT NEUTRON STAR  
COLLISION V. 2



ANIMATION: DANA BERRY  
310-441-1735

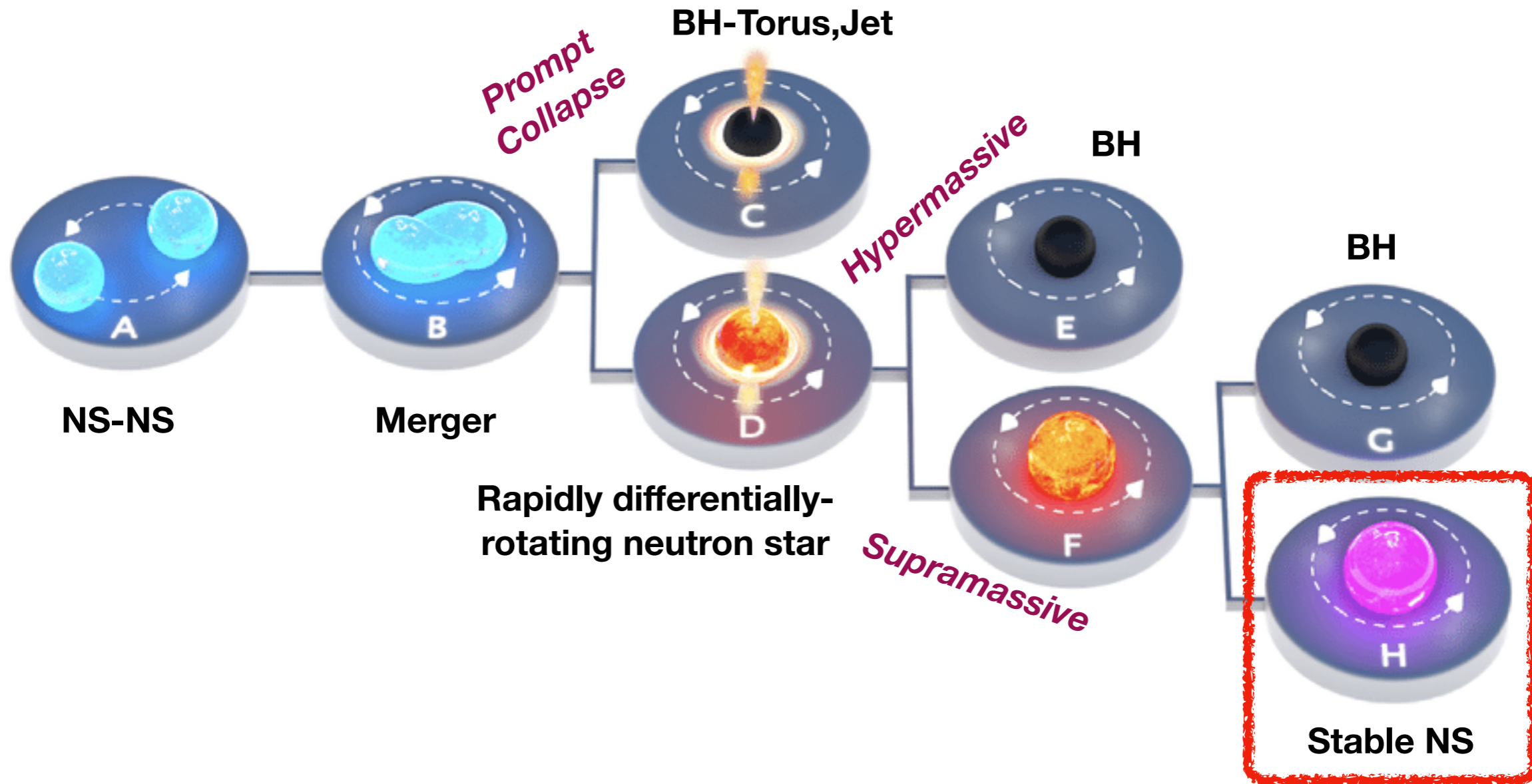
PRODUCED BY ERICA DREZEK

Fate decided by EOS, Mass, Spin, ....



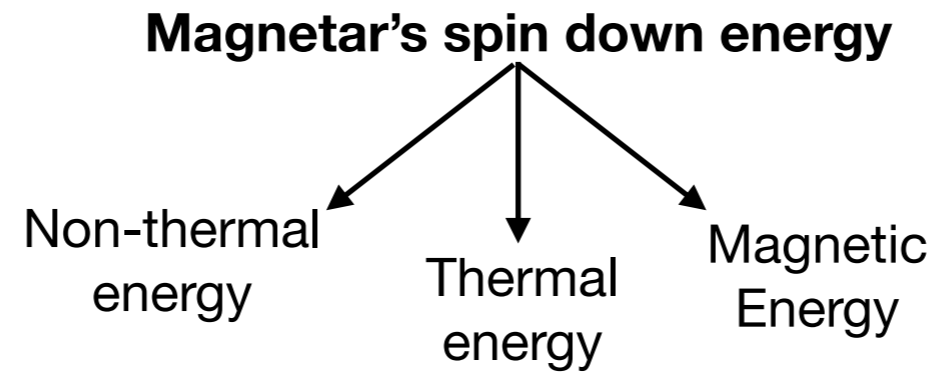
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Fate decided by EOS, Mass, Spin, ....

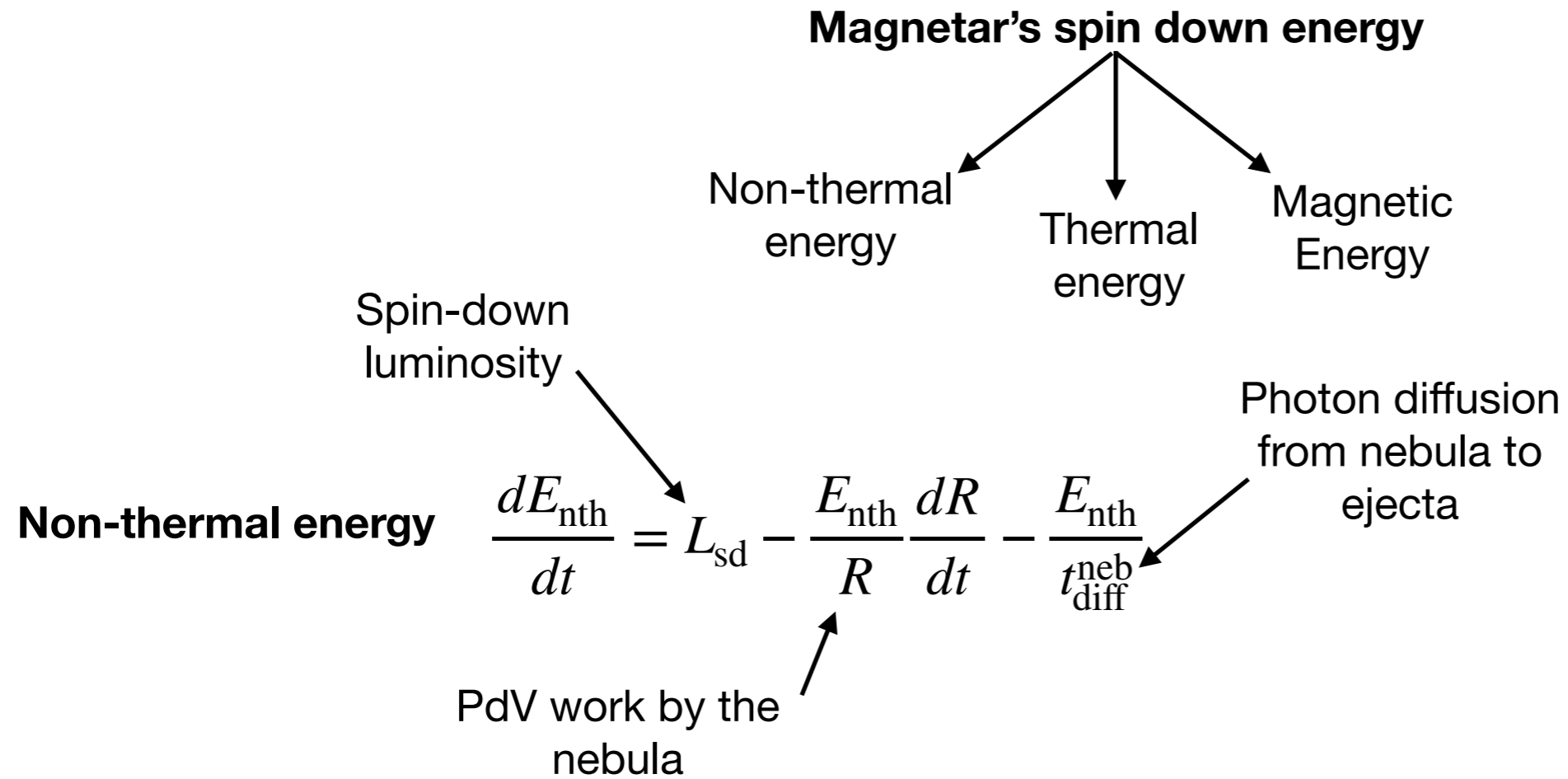




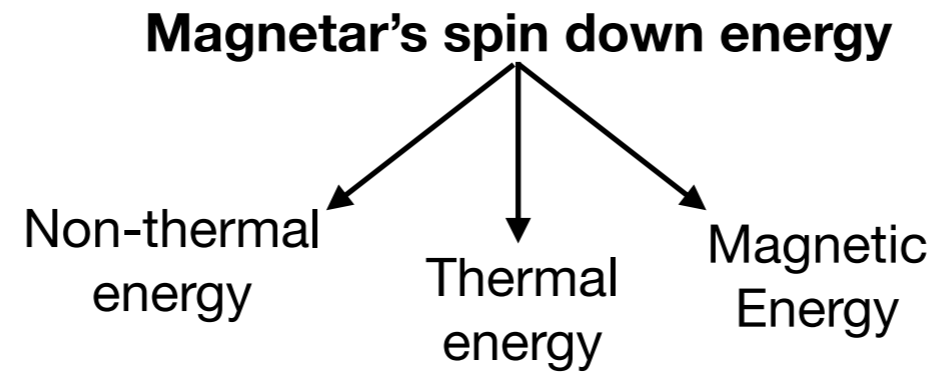
# Model: Evolution of thermal, non-thermal, and magnetic energies



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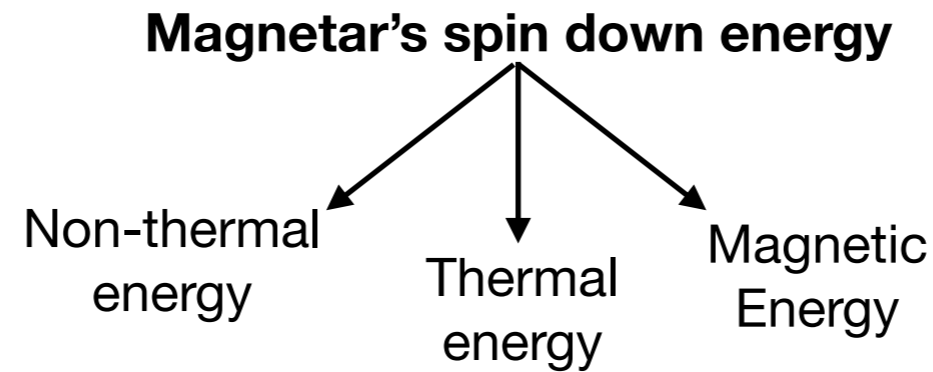


**Non-thermal energy**  $\frac{dE_{\text{nth}}}{dt} = L_{\text{sd}} - \frac{E_{\text{nth}}}{R} \frac{dR}{dt} - \frac{E_{\text{nth}}}{t_{\text{diff}}^{\text{neb}}}$

**Thermal energy**  $\frac{dE_{\text{th}}}{dt} = \left(1 - \mathcal{A}\right) \frac{E_{\text{nth}}}{t_{\text{diff}}^{\text{neb}}} - \frac{E_{\text{th}}}{R} \frac{dR}{dt} - \frac{E_{\text{th}}}{t_{\text{diff}}^{\text{ej}}} + Q_{\text{rp}}^{\text{heat}}$

Fraction of non-thermal photons that escape  $\rightarrow$   $\mathcal{A}$   
 Heating rate due to decay of r-process elements in the ejecta  $\rightarrow$   $Q_{\text{rp}}^{\text{heat}}$   
 Photon diffusion through ejecta  $\rightarrow$   $t_{\text{diff}}^{\text{ej}}$

# Model: Evolution of thermal, non-thermal, and magnetic energies



**Non-thermal energy**  $\frac{dE_{\text{nth}}}{dt} = L_{\text{sd}} - \frac{E_{\text{nth}}}{R} \frac{dR}{dt} - \frac{E_{\text{nth}}}{t_{\text{diff}}^{\text{neb}}}$

**Thermal energy**

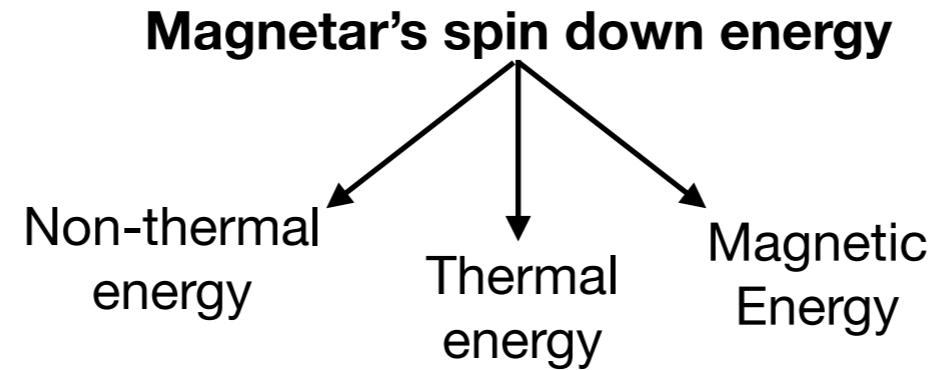
$$\frac{dE_{\text{th}}}{dt} = (1 - \mathcal{A}) \frac{E_{\text{nth}}}{t_{\text{diff}}^{\text{neb}}} - \frac{E_{\text{th}}}{R} \frac{dR}{dt} - \frac{E_{\text{th}}}{t_{\text{diff}}^{\text{ej}}} + Q_{\text{rp}}^{\text{heat}}$$

Magnetic field strength  
amplification parameter

$$\epsilon_B \sim 10^{-4}$$

**Magnetic energy**  $\frac{dE_B}{dt} = \epsilon_B L_{\text{sd}} - \frac{E_B}{R} \frac{dR}{dt}$

# Model: Evolution of thermal, non-thermal, and magnetic energies



**Non-thermal energy**  $\frac{dE_{\text{nth}}}{dt} = L_{\text{sd}} - \frac{E_{\text{nth}}}{R} \frac{dR}{dt} - \frac{E_{\text{nth}}}{t_{\text{diff}}^{\text{neb}}}$

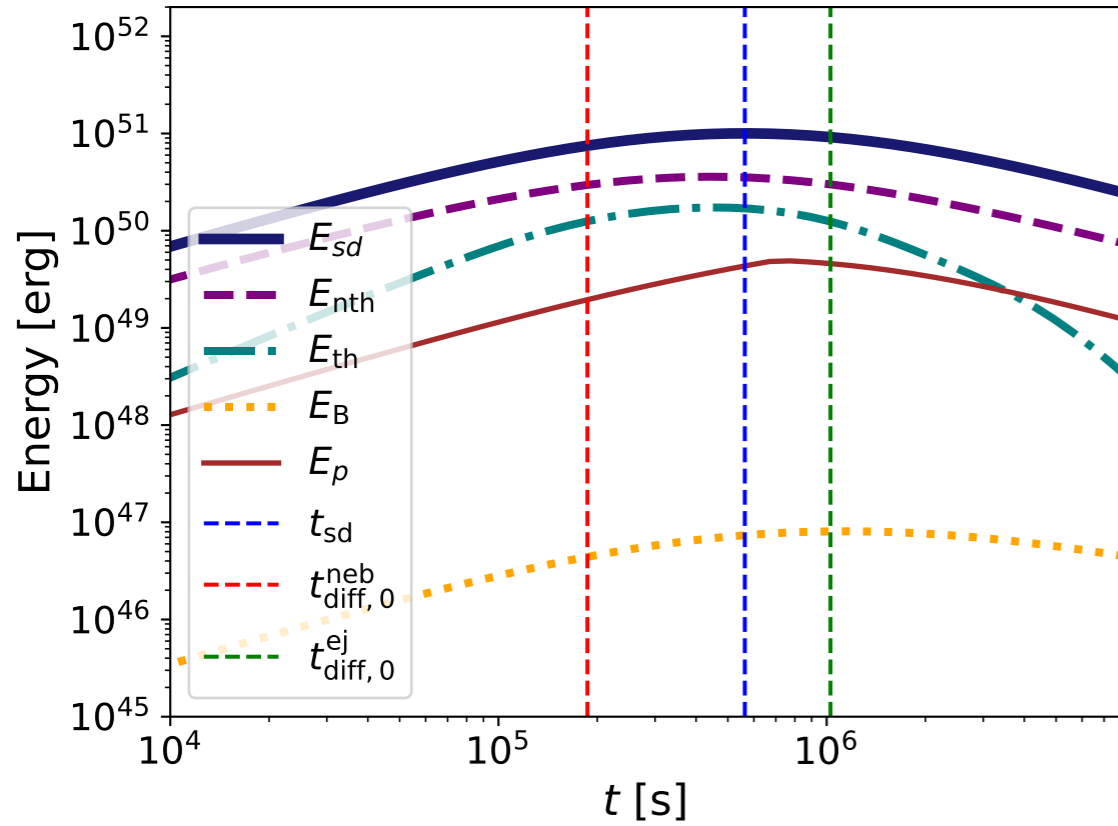
**Thermal energy**

$$\frac{dE_{\text{th}}}{dt} = (1 - \mathcal{A}) \frac{E_{\text{nth}}}{t_{\text{diff}}^{\text{neb}}} - \frac{E_{\text{th}}}{R} \frac{dR}{dt} - \frac{E_{\text{th}}}{t_{\text{diff}}^{\text{ej}}} + Q_{\text{rp}}^{\text{heat}}$$

**Magnetic energy**  $\frac{dE_B}{dt} = \epsilon_B L_{\text{sd}} - \frac{E_B}{R} \frac{dR}{dt}$

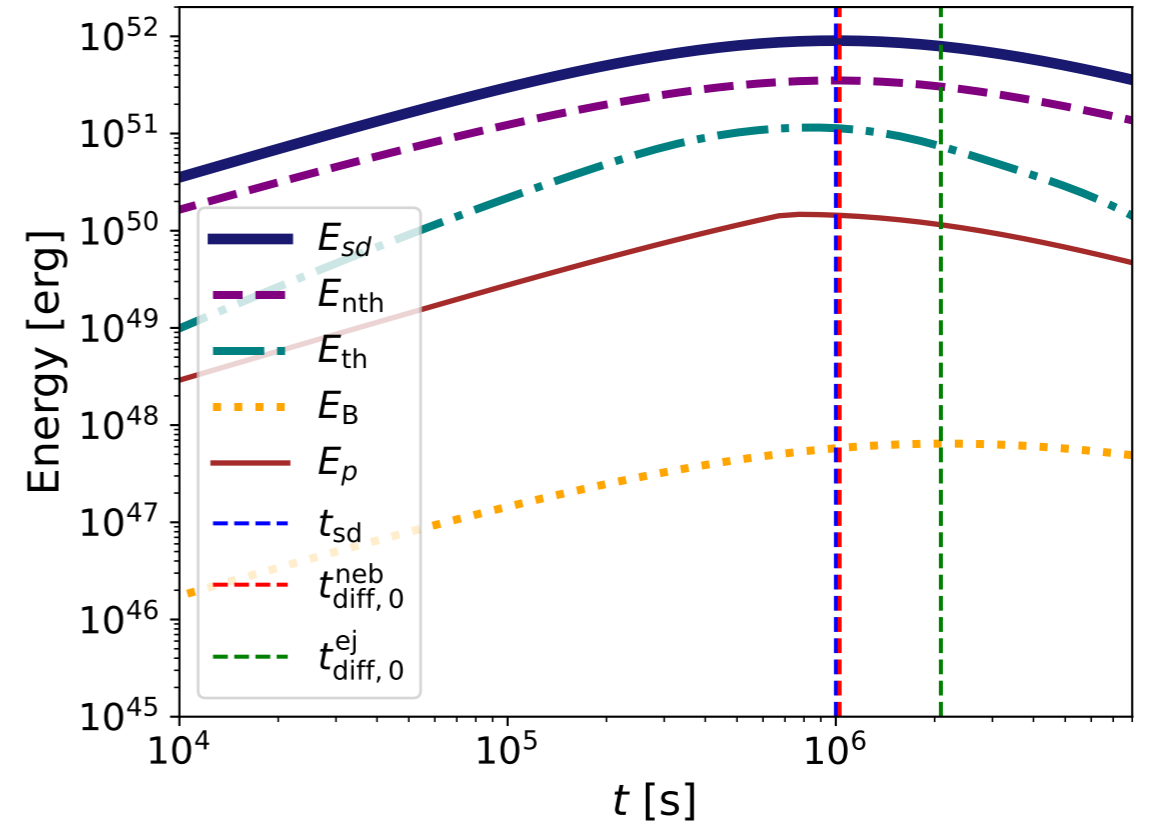
**Work done on the ejecta**  $\frac{d}{dt} E_{\text{kin}} = \frac{d}{dt} \left[ M_{\text{ej}} c^2 (\Gamma_{\text{ej}} - 1) \right] = \frac{v}{R} (E_{\text{nth}} + E_{\text{th}} + E_B)$

# Model: Evolution of thermal, non-thermal, and magnetic energies



Fiducial:

$$B_d = 10^{14} \text{ G}, P_i = 0.003 \text{ s}, M_{\text{ej}} = 0.03 M_{\odot}, \beta_{\text{ej}} = 0.03$$



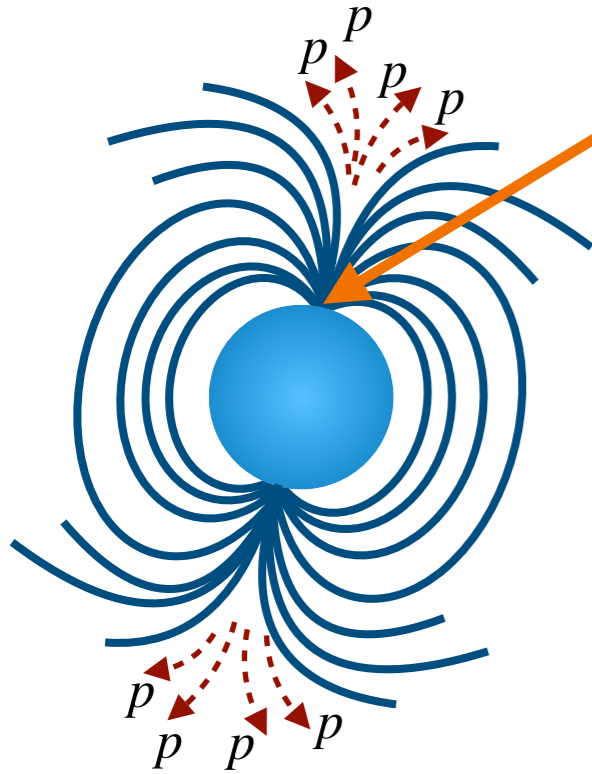
Optimistic:

$$B_d = 2.5 \times 10^{13} \text{ G}, P_i = 0.001 \text{ s}, M_{\text{ej}} = 0.1 M_{\odot}, \beta_{\text{ej}} = 0.1$$

$$L_{\text{sd}} = \alpha \frac{\mu^2 \Omega^4}{c^3} = 7.13 \times 10^{45} \text{ erg s}^{-1} \left( \frac{B_d}{10^{14} \text{ G}} \right)^2 \left( \frac{P_i}{0.003 \text{ s}} \right)^{-4} \left( 1 + \frac{t}{t_{\text{sd}}} \right)^{-2}$$

$$t_{\text{sd}} = 5.63 \times 10^5 \text{ s} \left( \frac{B_d}{10^{14} \text{ G}} \right)^{-2} \left( \frac{P_i}{0.003 \text{ s}} \right)^2$$

# Cosmic ray (CR) proton acceleration: injection spectra



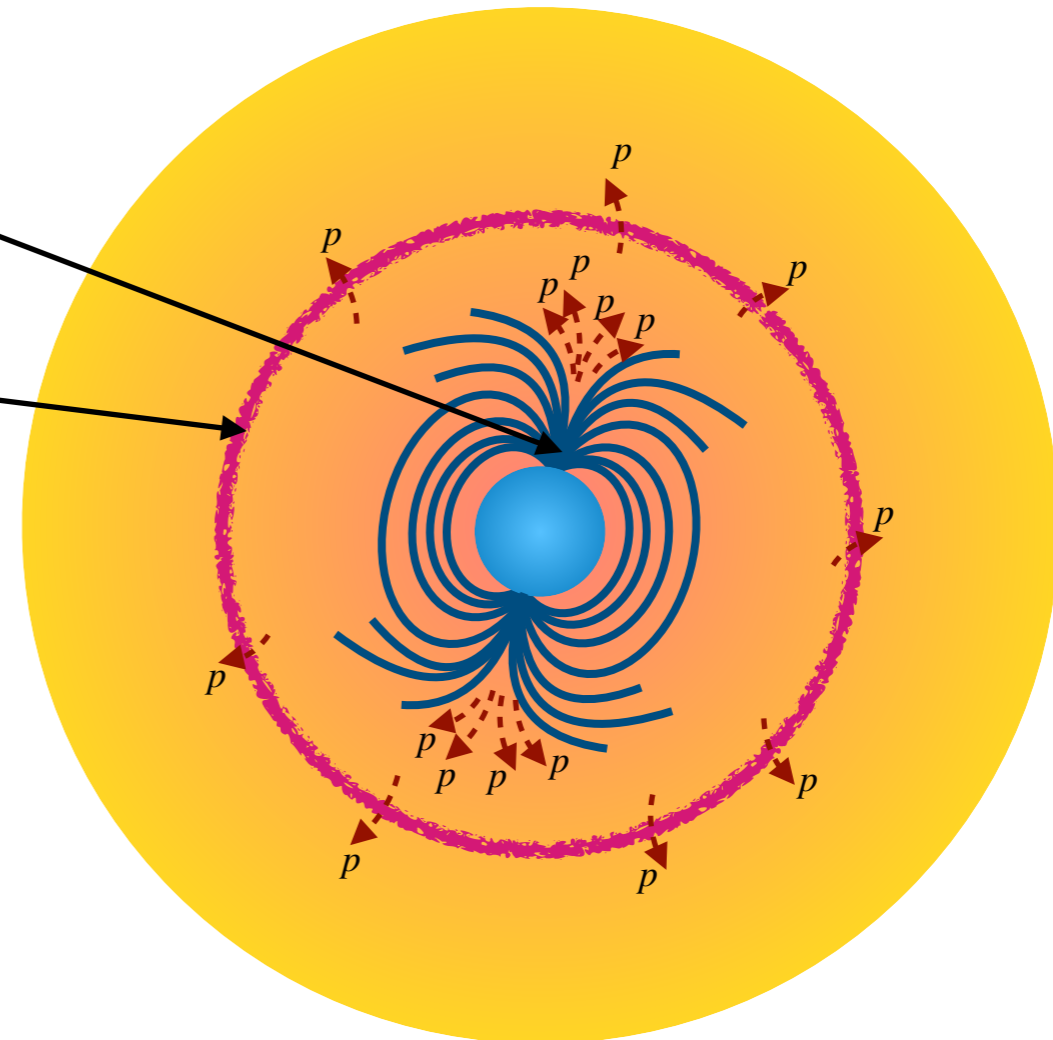
CR protons extracted from the magnetar surface: Goldreich-Julian (GJ) number density of charges

$$n_{\text{GJ}} = -\frac{\boldsymbol{\Omega} \cdot \mathbf{B}}{2\pi Z e c}$$

$$\dot{N}_p = n_{\text{GJ}} 2A_{\text{pc}} c = \frac{4\pi^2 R_*^3 B_0}{Z e c P^2}$$

Acceleration sites:

Polar cap  
+  
Termination shock (TS)

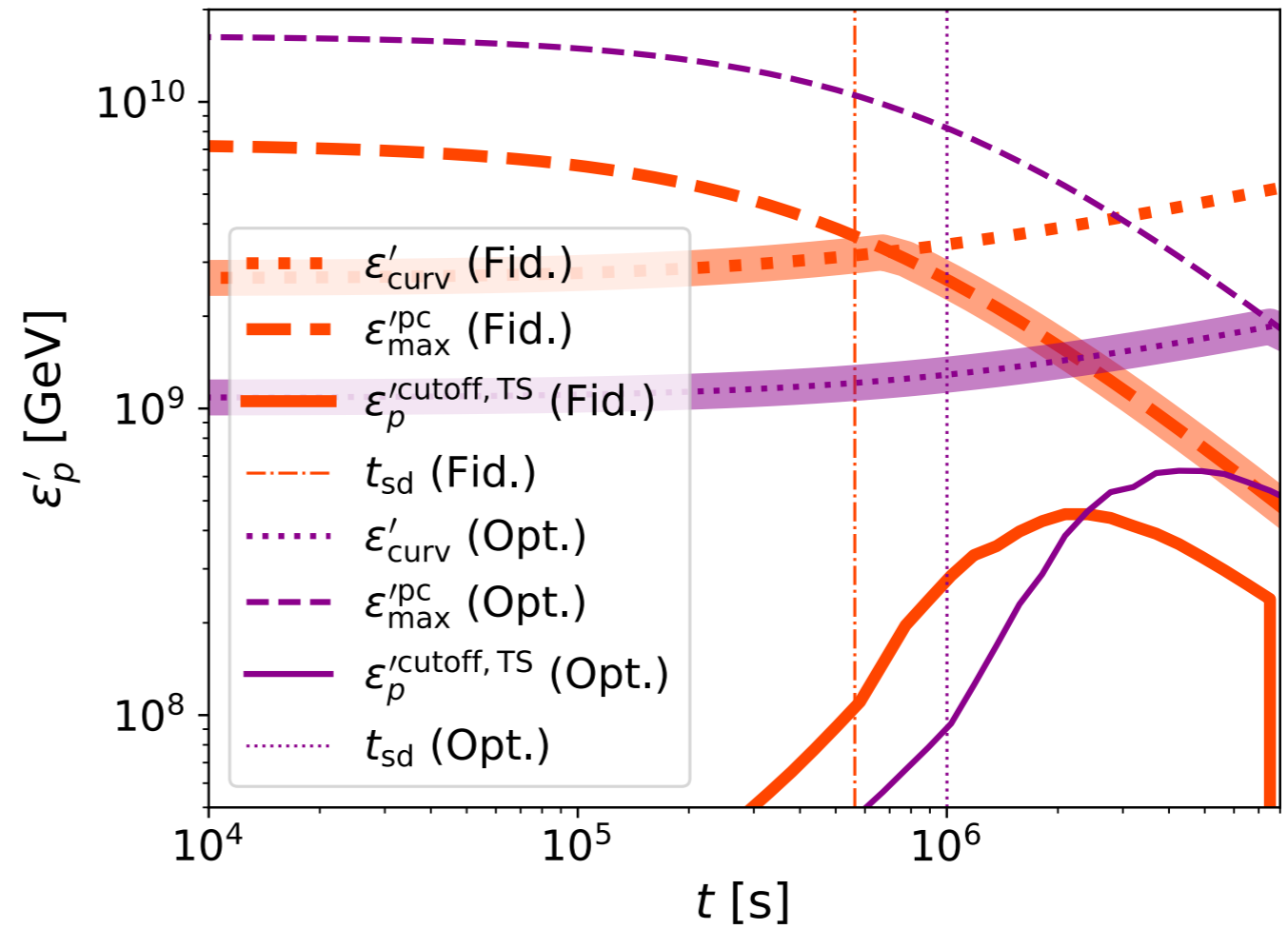
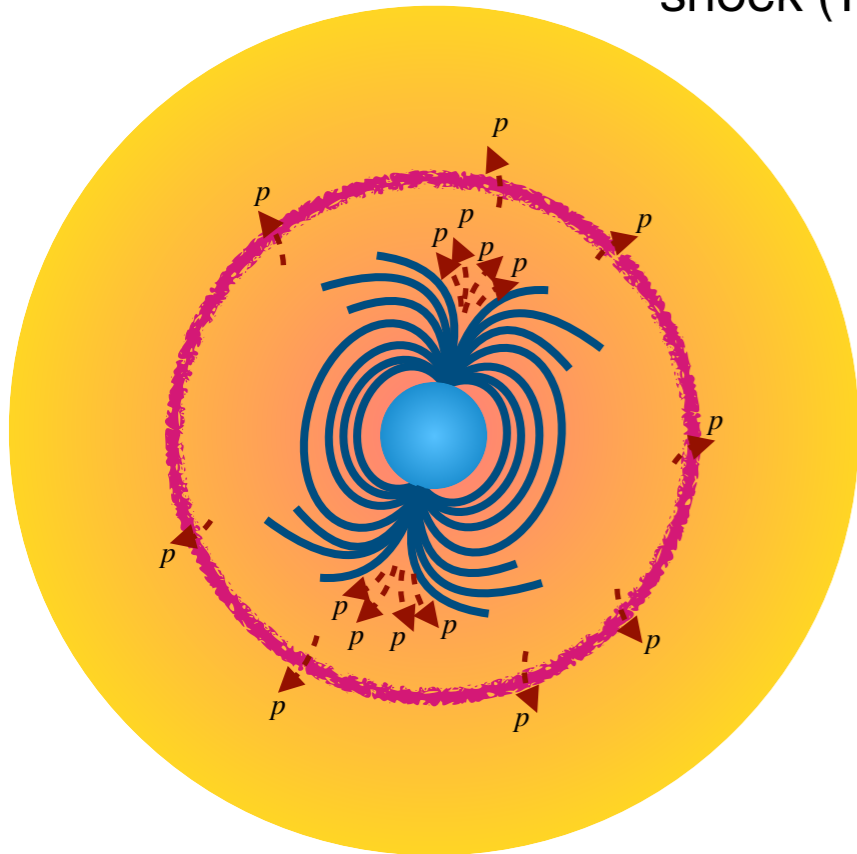


# Cosmic ray (CR) proton acceleration: injection spectra

$$\frac{d\dot{N}_{p,\text{inj}}}{d\varepsilon'_p} = \dot{N}_p^{\text{norm}} Q_p^{\text{inj}}(\varepsilon'_p) = \dot{N}_p^{\text{norm}} \delta(\varepsilon'_p - \varepsilon_p^{\text{cutoff,pc}})$$

$$\varepsilon_p^{\text{cutoff}} = \max \left[ \varepsilon_p^{\text{cutoff,pc}}, \varepsilon_p^{\text{cutoff,TS}} \right]$$

Acceleration sites: Polar cap + Termination shock (TS)



# Cosmic ray (CR) proton acceleration: injection spectra

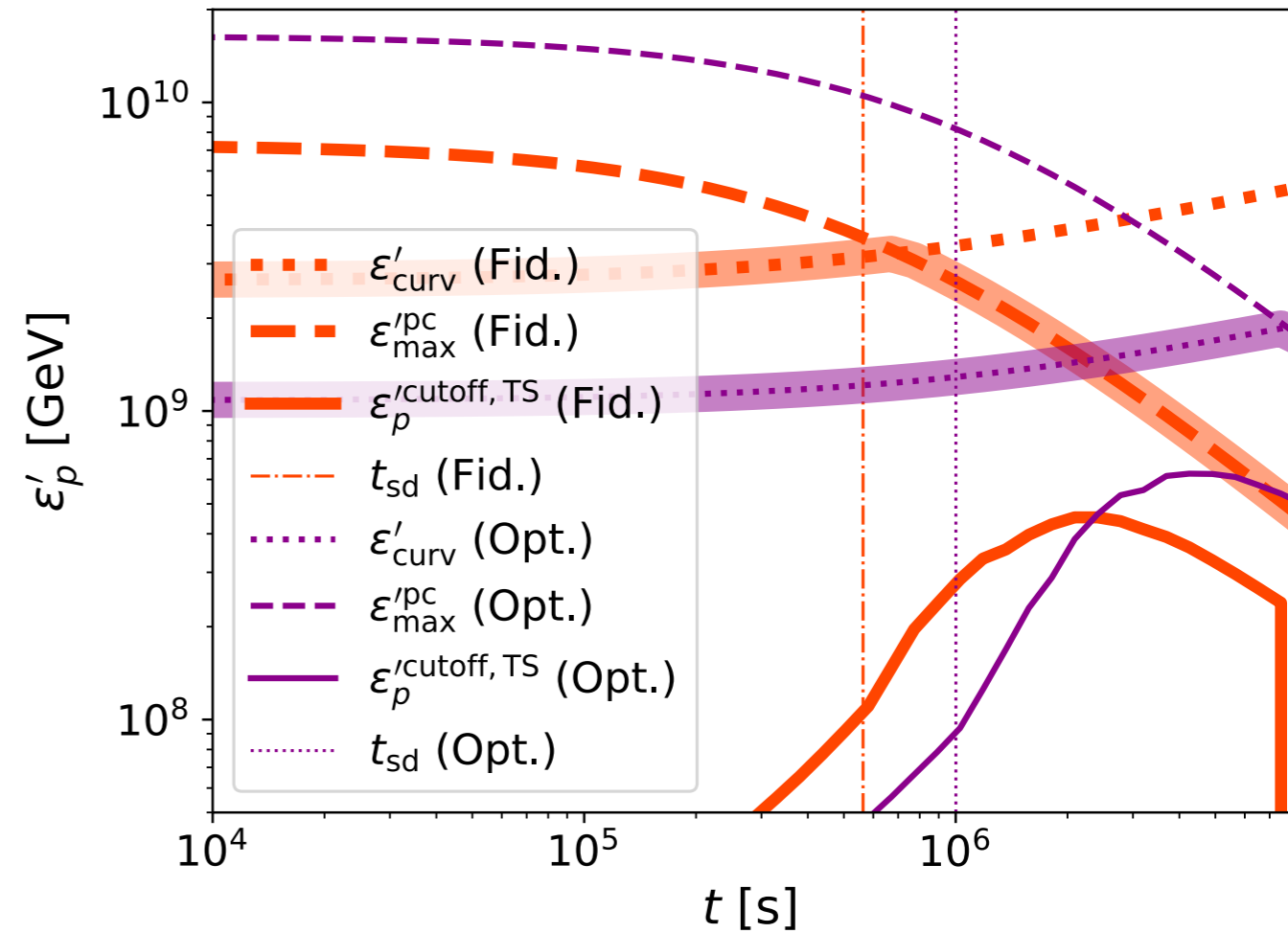
$$\frac{d\dot{N}_{p,\text{inj}}}{d\varepsilon'_p} = \dot{N}_p^{\text{norm}} Q_p^{\text{inj}}(\varepsilon'_p) = \dot{N}_p^{\text{norm}} \delta(\varepsilon'_p - \varepsilon_p^{\text{'cutoff,pc}})$$

$$\varepsilon_p^{\text{'cutoff}} = \max \left[ \varepsilon_p^{\text{'cutoff,pc}}, \varepsilon_p^{\text{'cutoff,TS}} \right]$$

$$\varepsilon_p^{\text{'cutoff,pc}} = \min \left[ \varepsilon_{\text{max}}^{\text{'pc}}, \varepsilon'_{\text{curv}} \right]$$

$$\varepsilon_{\text{max}}^{\text{'pc}} = 4\eta_{\text{gap}}(Ze)B_d \left( \frac{\pi R_*}{cP} \right)^2 R_*$$

$$\varepsilon'_{\text{curv}} = \gamma_p m_p c^2 = \left[ \frac{3m_p^4 c^8 B_d R_{\text{curv}}^2}{2e} \right]^{1/4}$$



# Cosmic ray (CR) proton acceleration: injection spectra

$$\frac{d\dot{N}_{p,\text{inj}}}{d\varepsilon'_p} = \dot{N}_p^{\text{norm}} Q_p^{\text{inj}}(\varepsilon'_p) = \dot{N}_p^{\text{norm}} \delta(\varepsilon'_p - \varepsilon_p^{\text{'cutoff,pc}})$$

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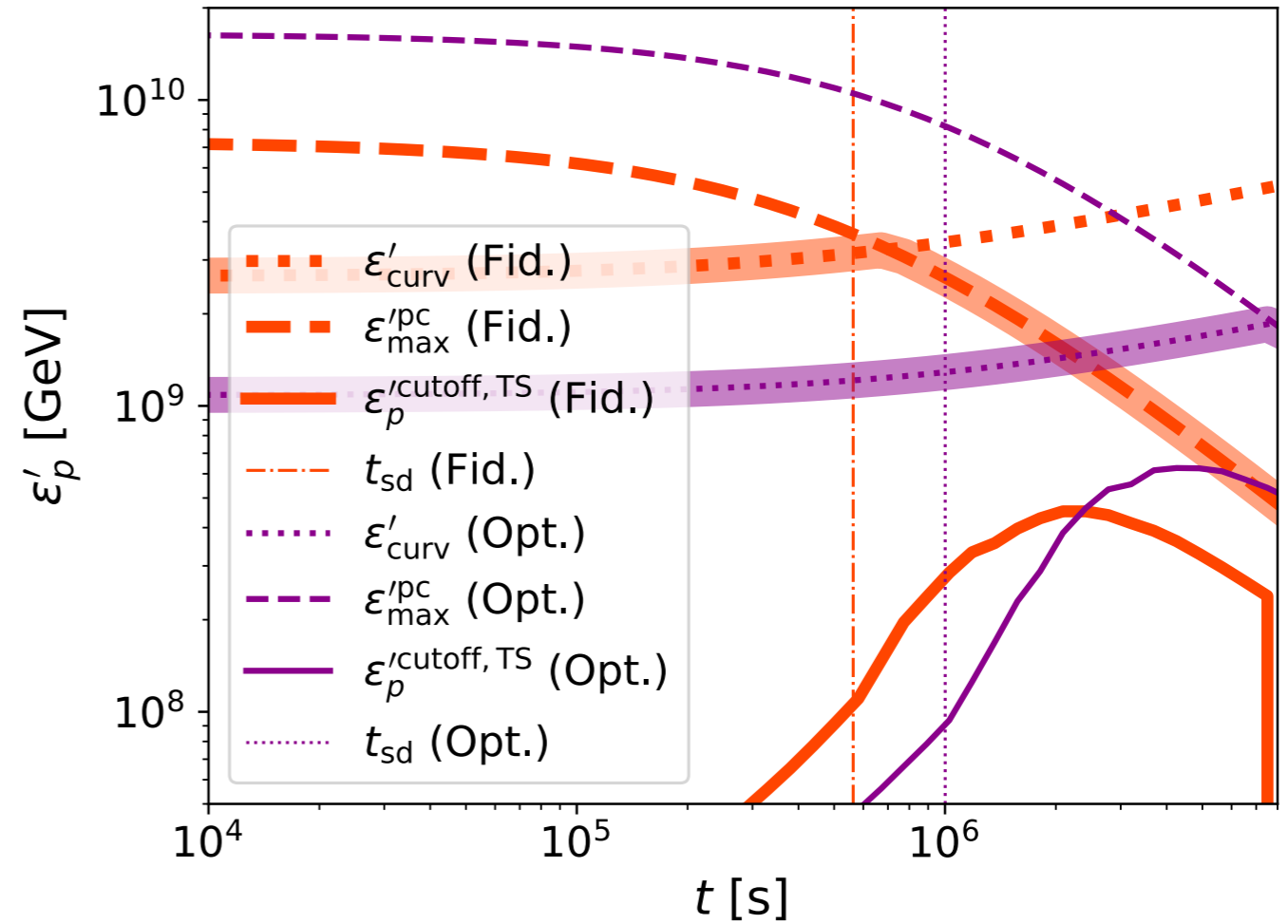
$$t_{\text{acc}}^{\prime-1} = t_{\text{loss}}^{\prime-1}$$

$$t_{\text{acc}}^{\prime-1} = \eta_{\text{acc}} \varepsilon'_p / (ZecB'_{\text{neb}})$$

$$t_{\text{loss}}^{\prime-1} = t_{\text{esc}}^{\prime-1} + t_{\text{cool}}^{\prime-1}$$

$$t_{\text{esc}}^{\prime-1} = \max \left[ R(t)^2 / D_c(\varepsilon'_p), R(t) / c \right]$$

$$t_{\text{cool}}^{\prime-1} = t_{pp}^{\prime-1} + t_{p\gamma}^{\prime-1} + t_{\text{sync}}^{\prime-1} + t_{\text{BH}}^{\prime-1} + t_{\text{dyn}}^{\prime-1}$$



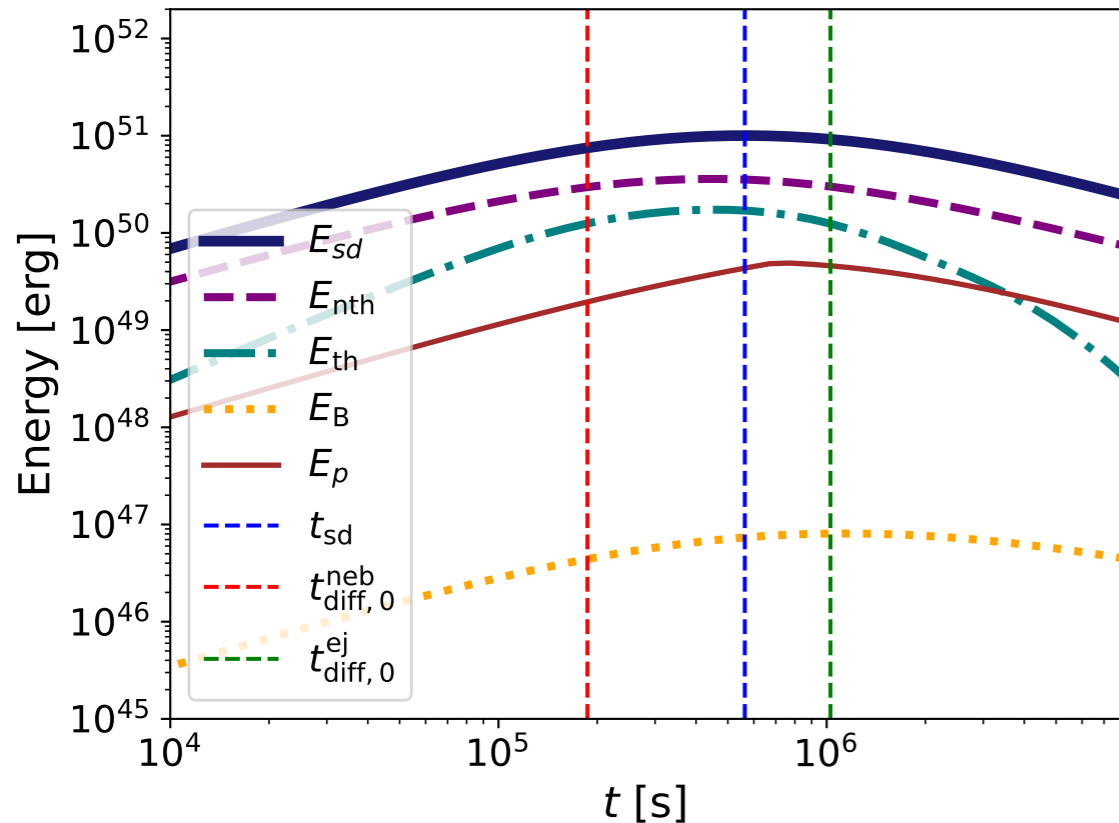
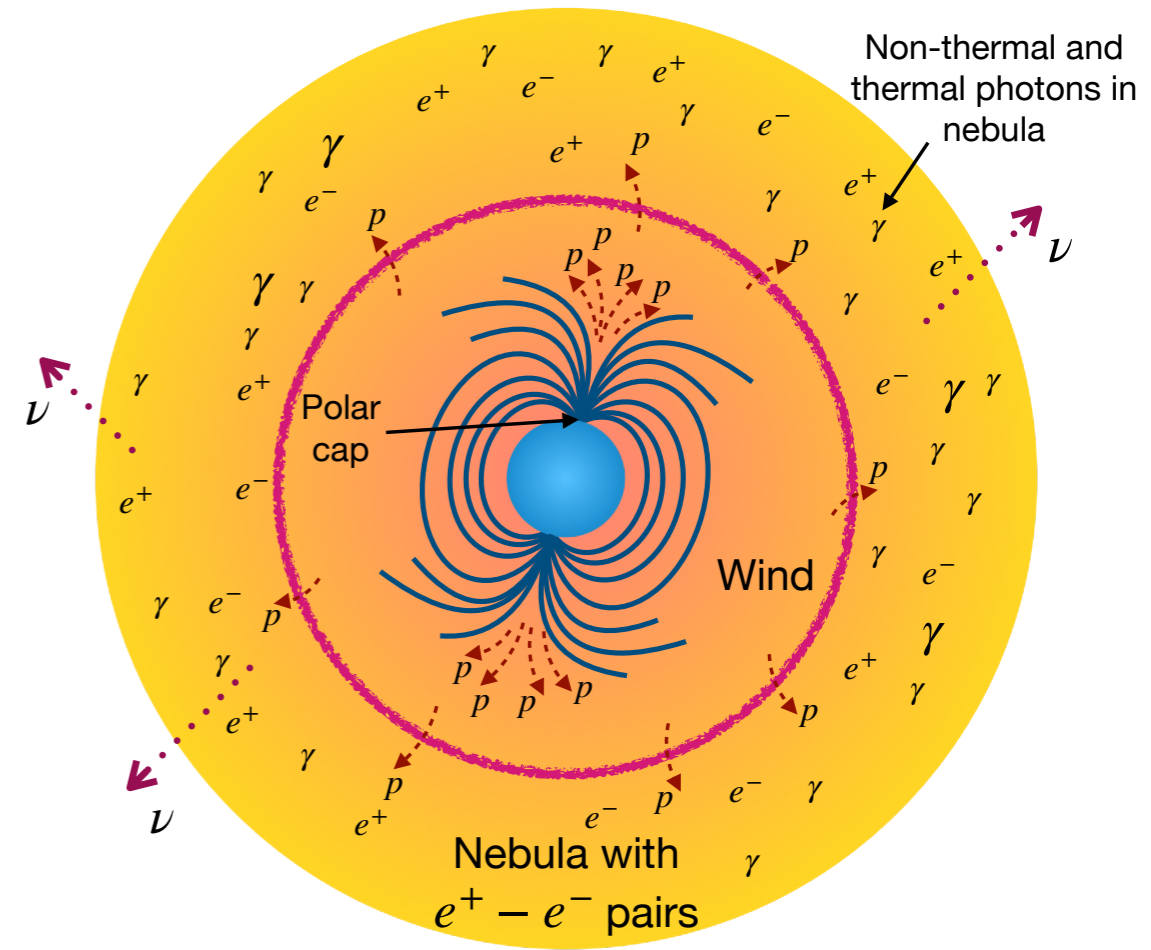
# Cosmic ray (CR) proton acceleration

Compute steady state CR spectrum by solving the transport equation

$$\frac{d}{d\varepsilon'_p} \left( -\frac{\varepsilon'_p}{t'_{\text{cool}}} \frac{dN_p}{d\varepsilon'_p} \right) = \frac{d\dot{N}_{p,\text{inj}}}{d\varepsilon'_p} - \frac{1}{t'_{\text{esc}}} \frac{dN_p}{d\varepsilon'_p}$$

$$\frac{dN_p}{d\varepsilon'_p} = \frac{t'_{\text{cool}}}{\varepsilon'_p} \int_{\varepsilon'_p}^{\infty} d\tilde{\varepsilon}_p \dot{N}_{p,\text{inj}}(\tilde{\varepsilon}_p) \exp\left(-\mathcal{G}(\varepsilon'_p, \tilde{\varepsilon}_p)\right)$$

$$\mathcal{G}(\varepsilon_1, \varepsilon_2) = \int_{\varepsilon_1}^{\varepsilon_2} \frac{t'_{\text{cool}}}{t'_{\text{esc}}} \frac{d\tilde{\varepsilon}_p}{\tilde{\varepsilon}_p}$$



$$E_p = \dot{N}_{p,\text{inj}} \varepsilon_p^{\text{cutoff,pc}} t$$

# Cosmic ray (CR) proton acceleration

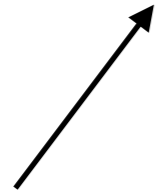
Compute steady state CR spectrum by solving the transport equation



This along with the photon field spectrum gives the neutrino fluences

$e^+ - e^-$  spectra

$$\frac{dN}{d\gamma_e} \sim \begin{cases} \gamma_e^{-1.5}, & \gamma_e \leq \gamma_{e,br} \\ \gamma_e^{-2.5}, & \gamma_e > \gamma_{e,br} \end{cases}$$

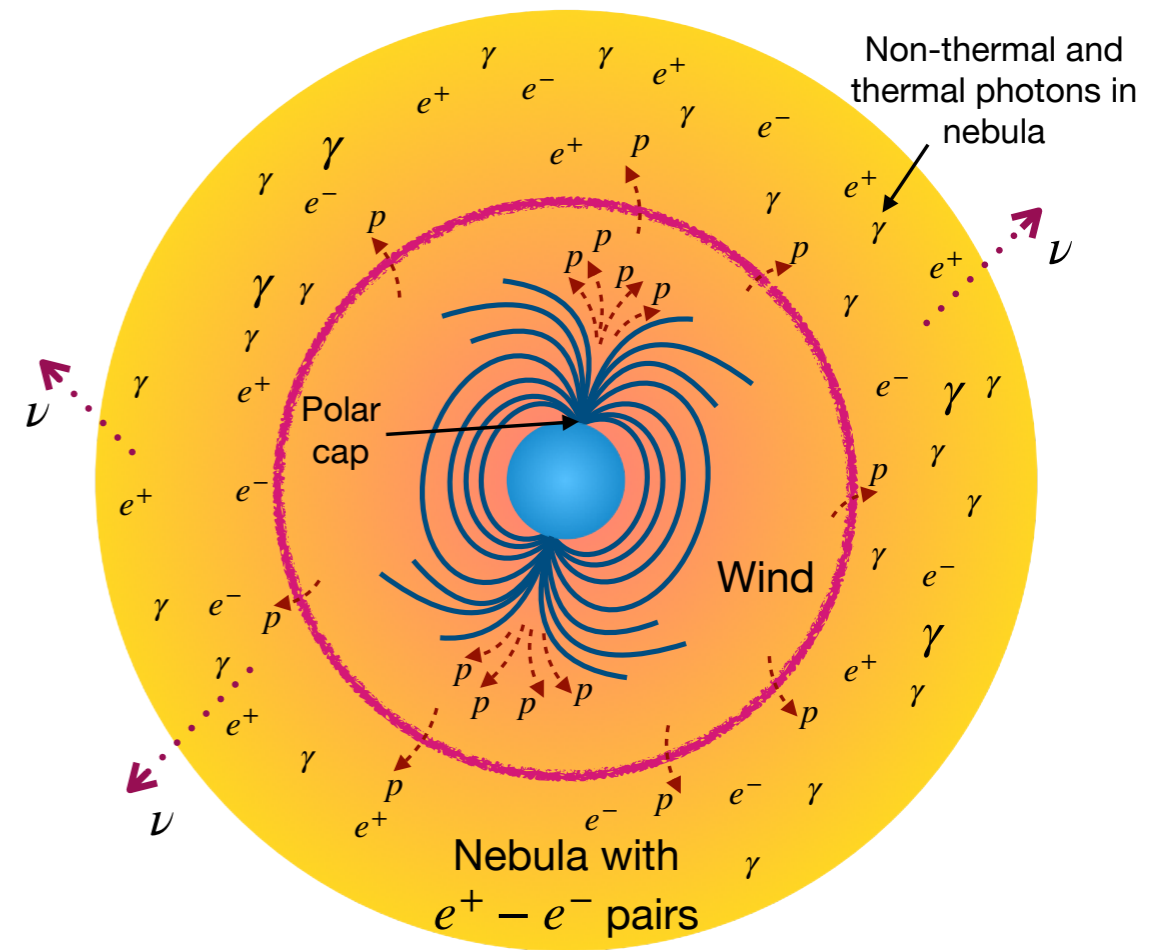


Electron break Lorentz factor

For galactic PWNe:

$$\gamma_{e,br} \sim 10^5 - 10^6$$

Pair injection at upstream of TS -> decreased wind velocity and hence lower  $\gamma_{e,br}$ . We choose  $\gamma_{e,br} = 10^3$



# Cosmic ray (CR) proton acceleration

Compute steady state CR spectrum by solving the transport equation



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Electron break Lorentz factor

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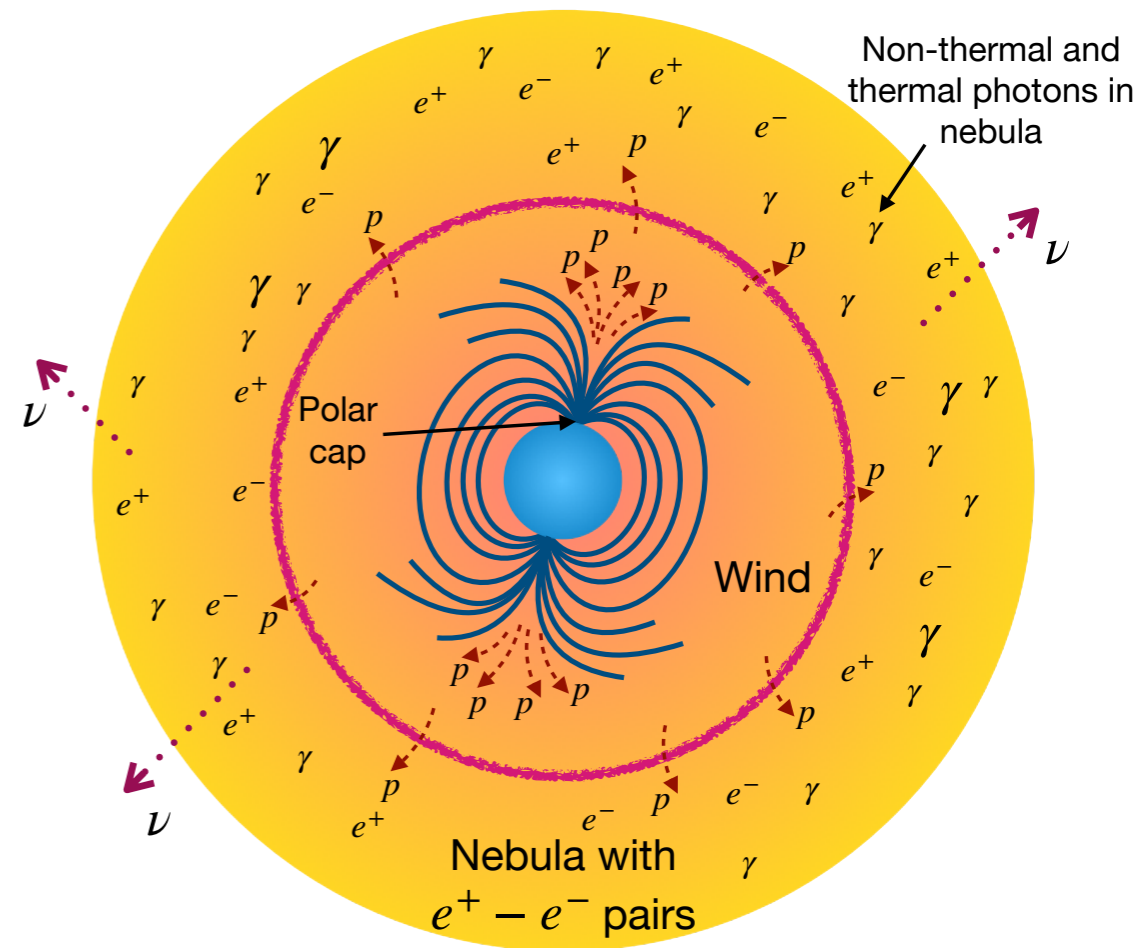
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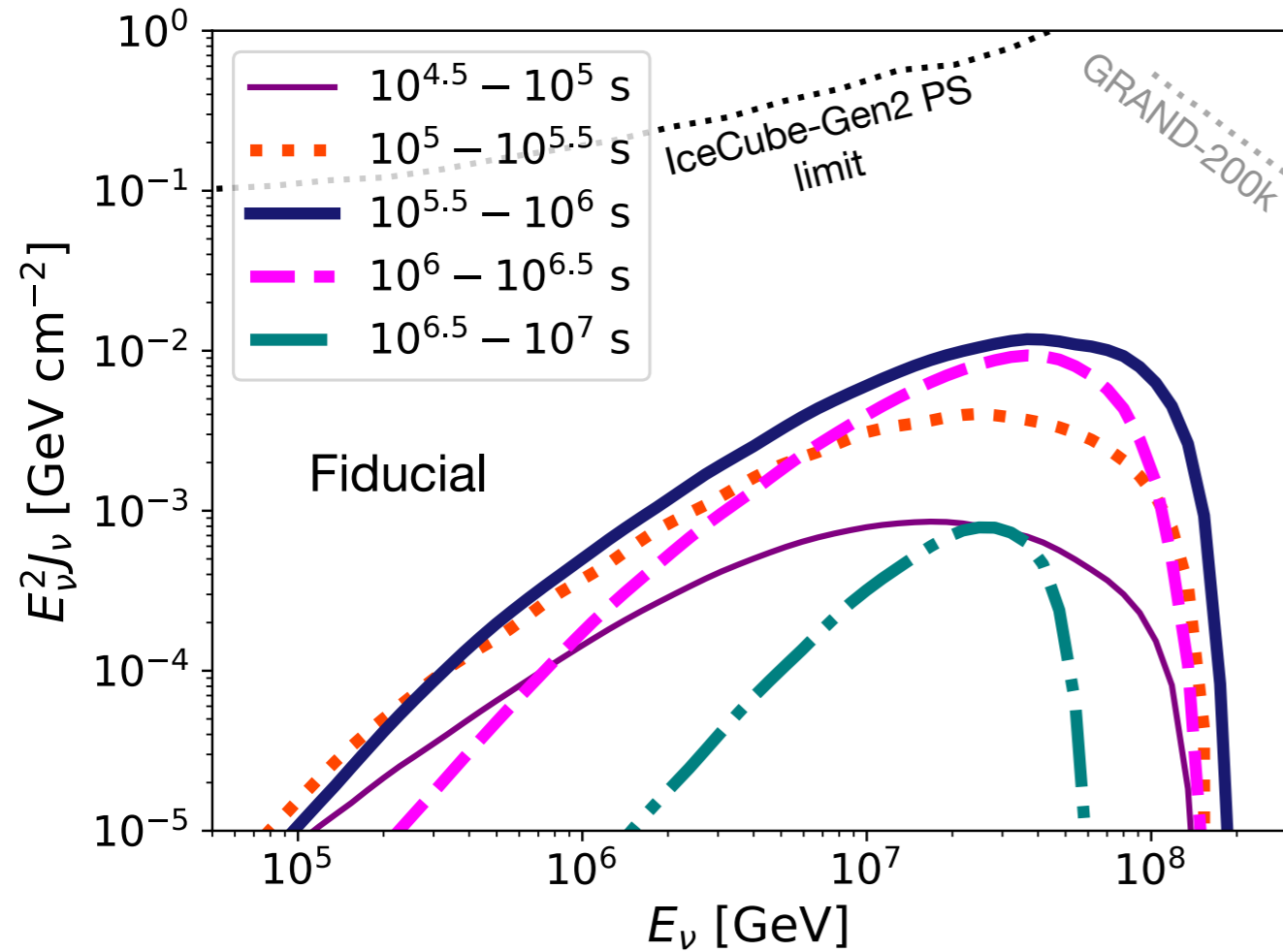
Transport Equation

Synchrotron, inverse Compton,  $\gamma\gamma$  processes

EM cascades



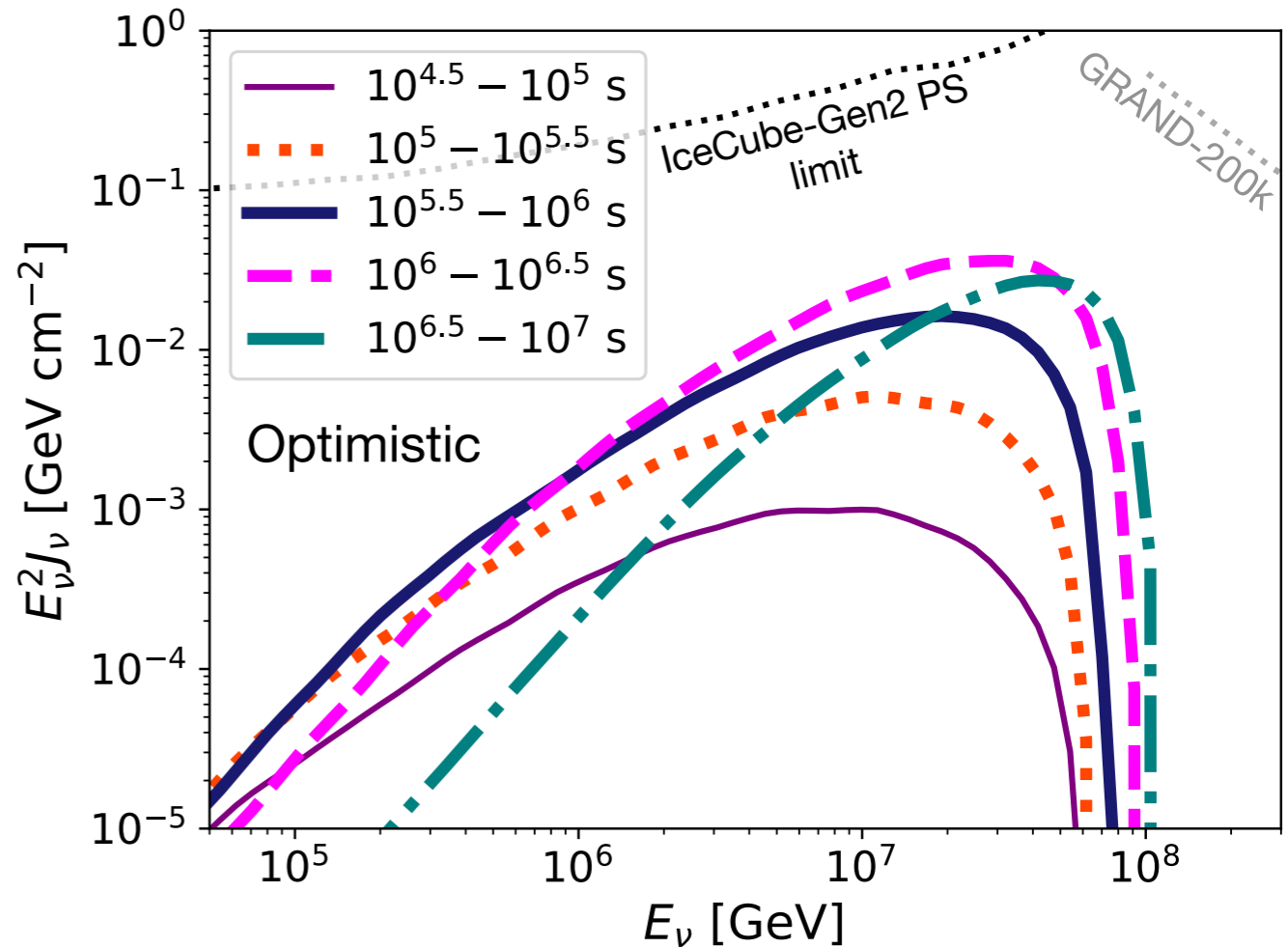
# The money plot: Neutrino fluences (takeaway)



Peak fluence:  $\sim 1 \times 10^{-2} \text{ GeV cm}^{-2}$

Neutrino energy:  $\sim 10^7 \text{ GeV} - 10^8 \text{ GeV}$

Peak fluence  $\sim 10^{5.5} \text{ s} - 10^6 \text{ s}$  post-merger

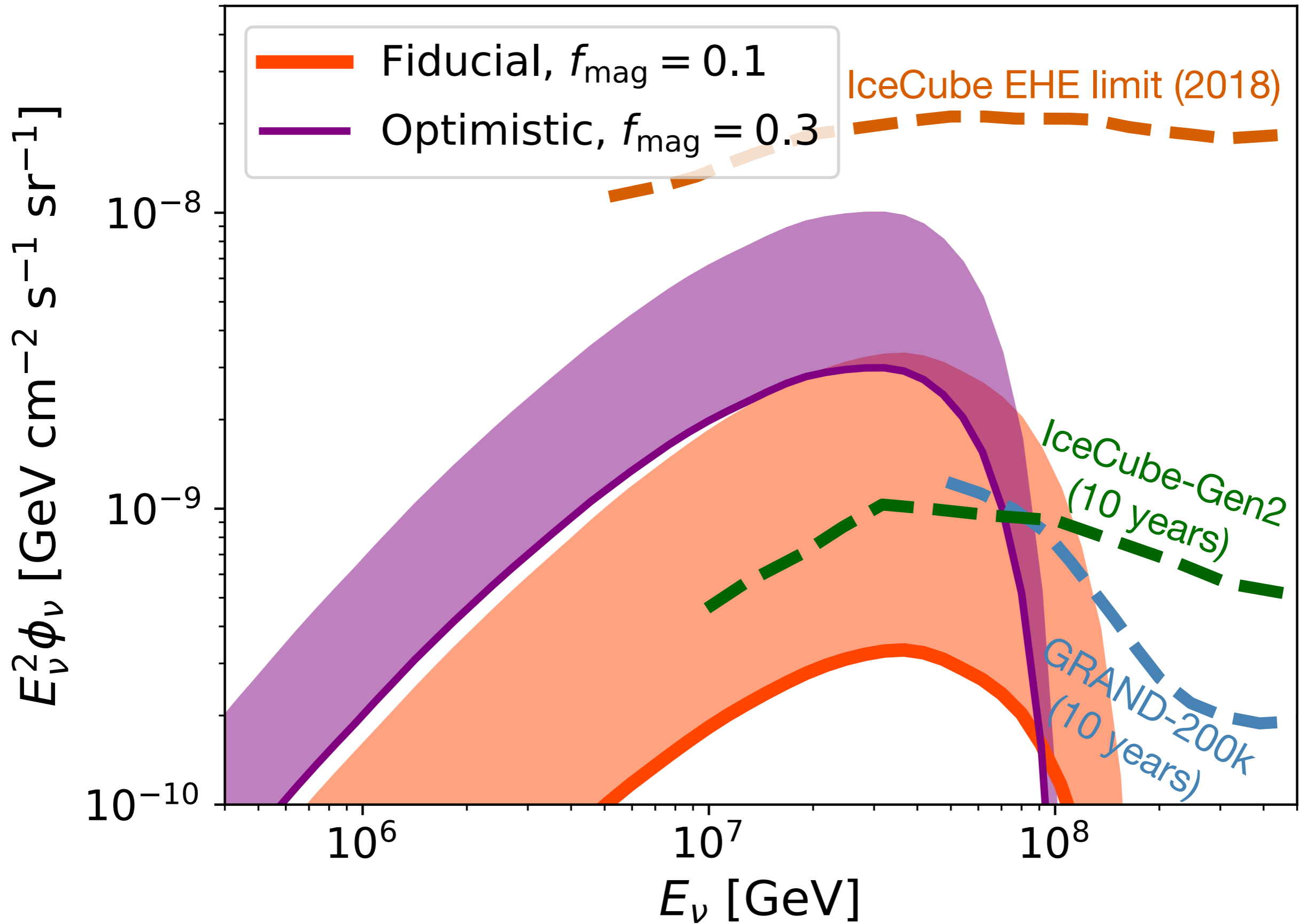


Peak fluence:  $\sim 4 \times 10^{-2} \text{ GeV cm}^{-2}$

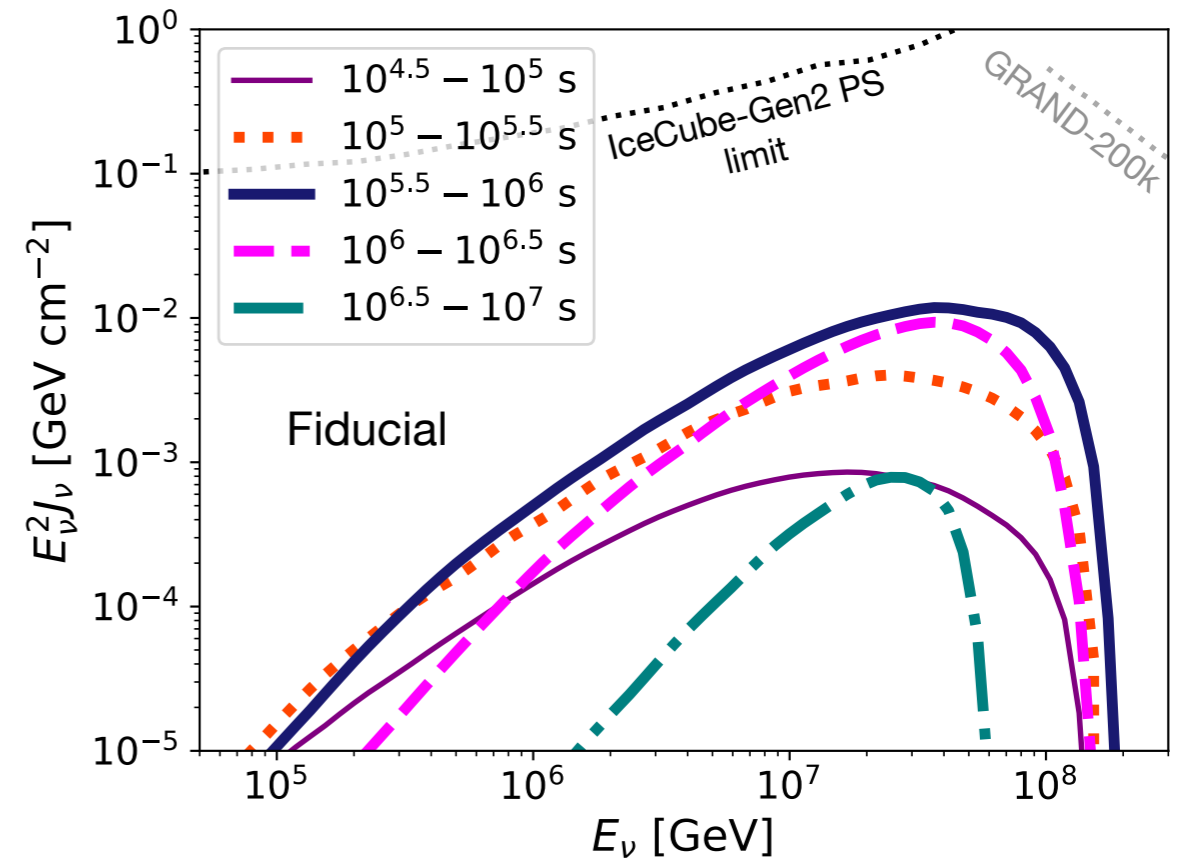
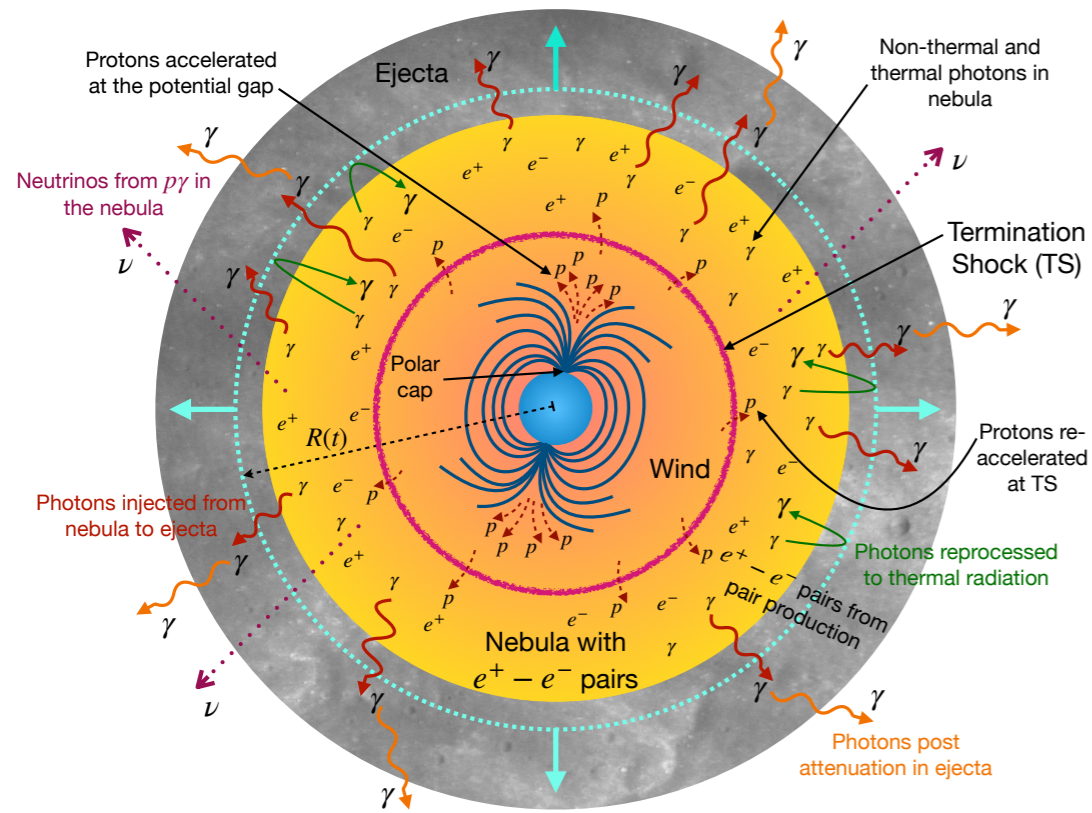
Peak fluence  $\sim 10^6 \text{ s} - 10^{6.5} \text{ s}$  post-merger

$d_L = 40 \text{ Mpc}$

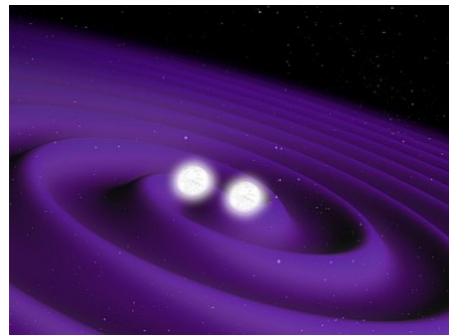
# Diffuse neutrino flux



# Takeaways



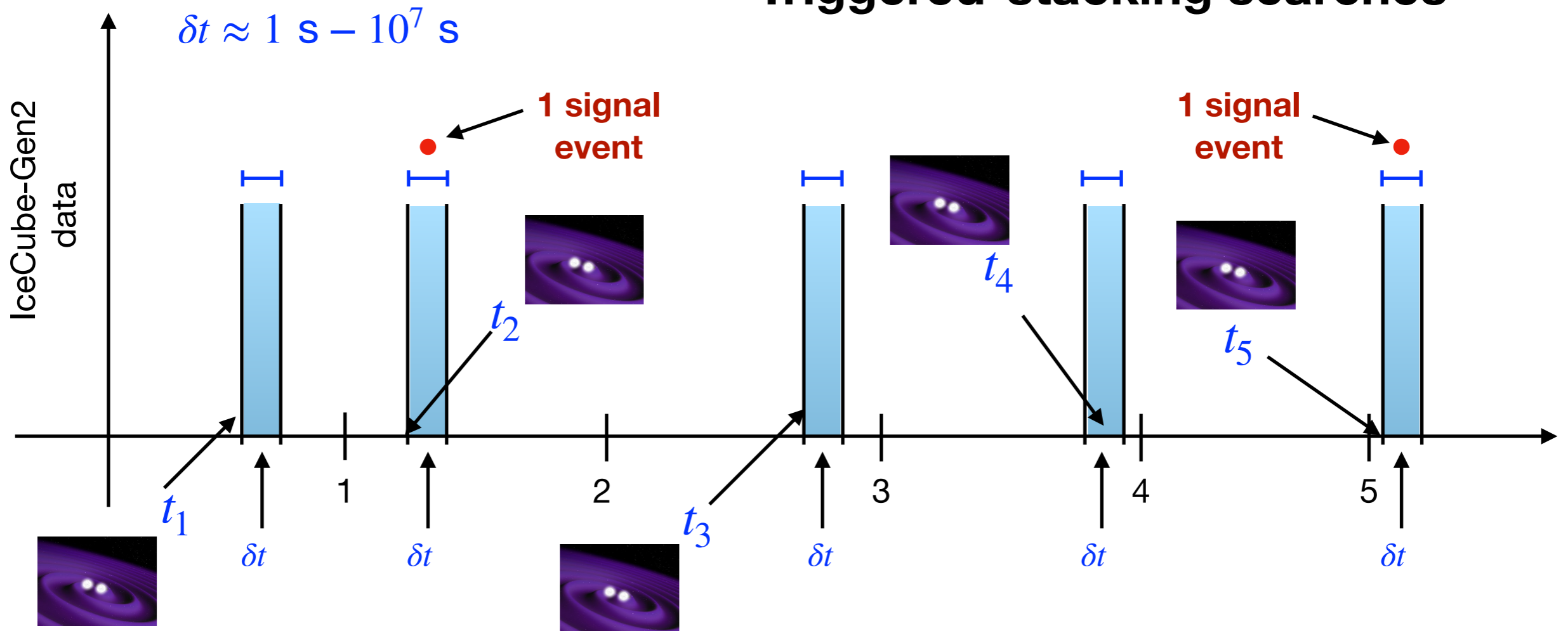
# Detection strategy: triggered stacking search



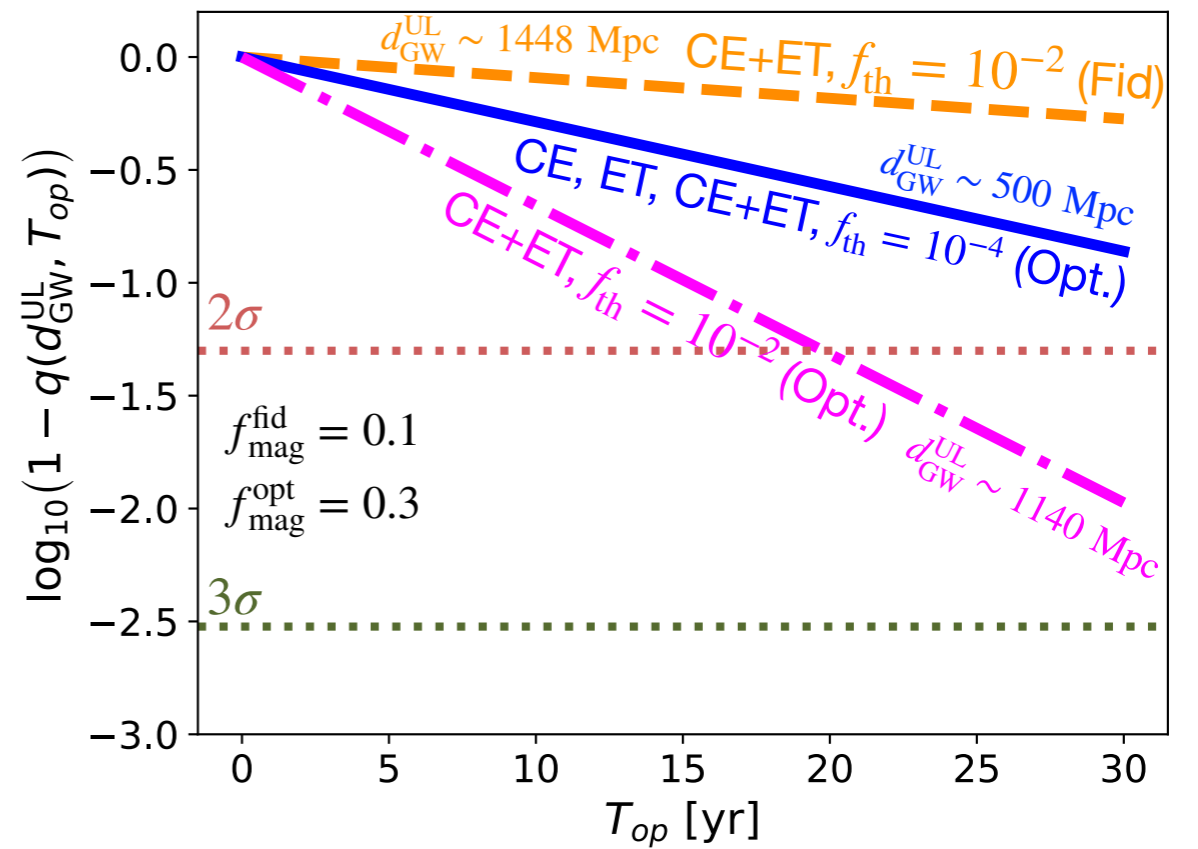
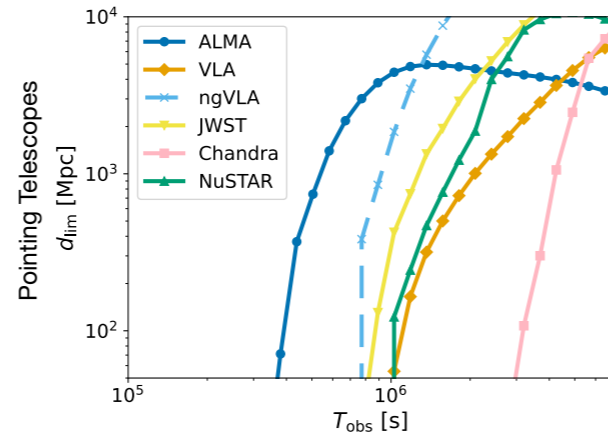
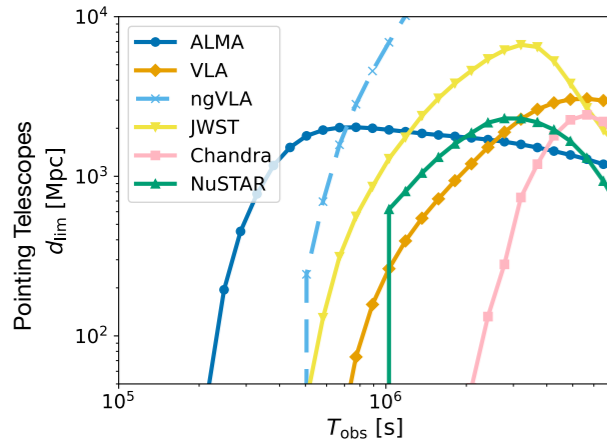
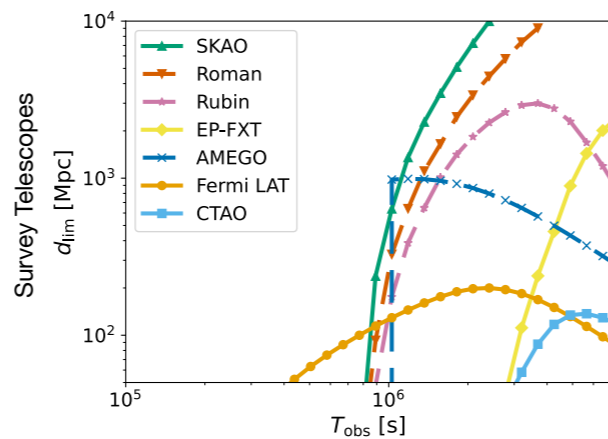
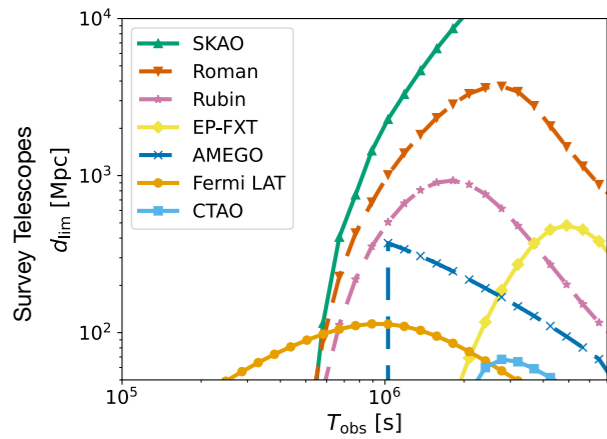
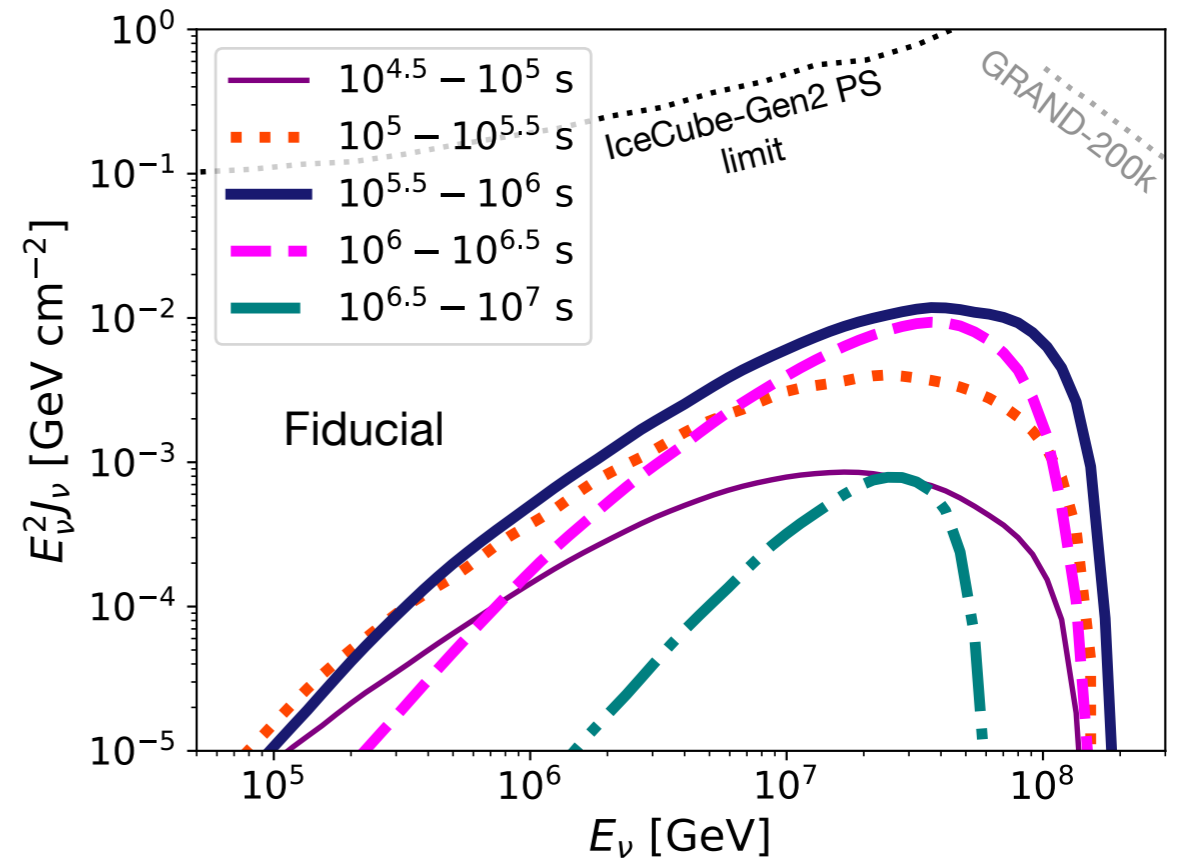
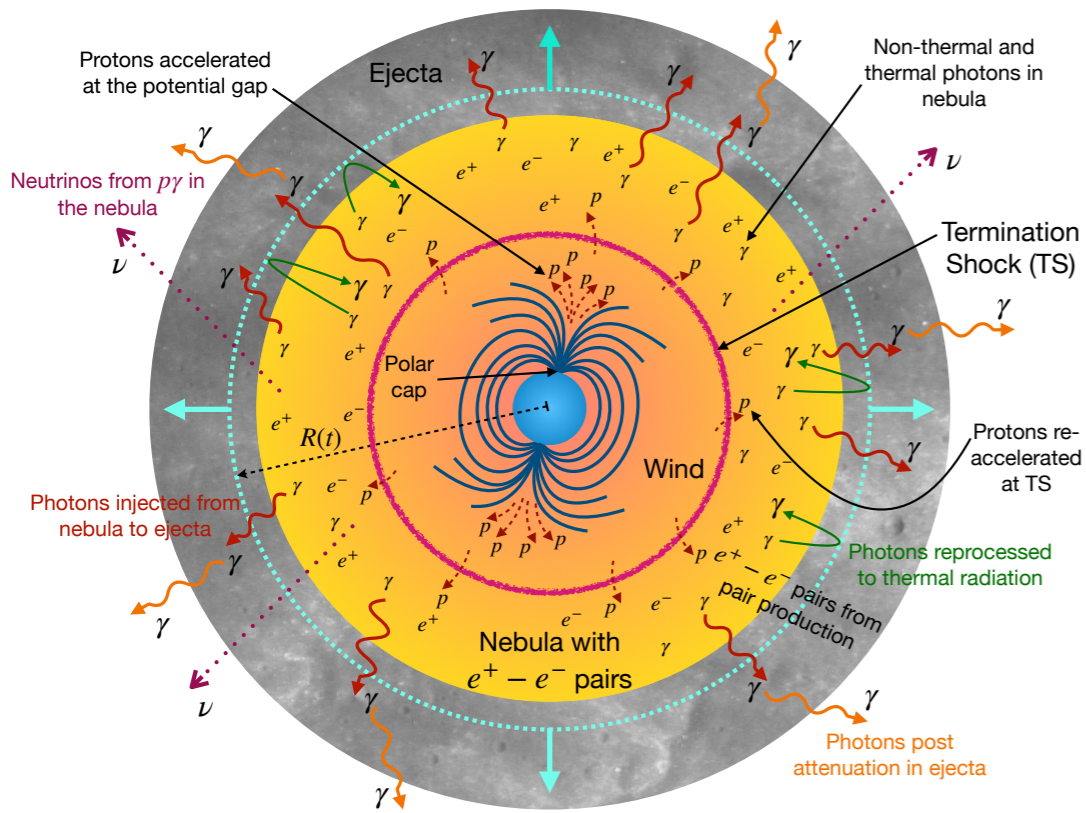
Trigger from next-gen GW detectors (ET, CE)

Neutrinos in IceCube-Gen 2

## Triggered-stacking searches



# Takeaways



# Future prospects: the big picture



**Fate of BNS merger remnants** (in particular stable remnants)

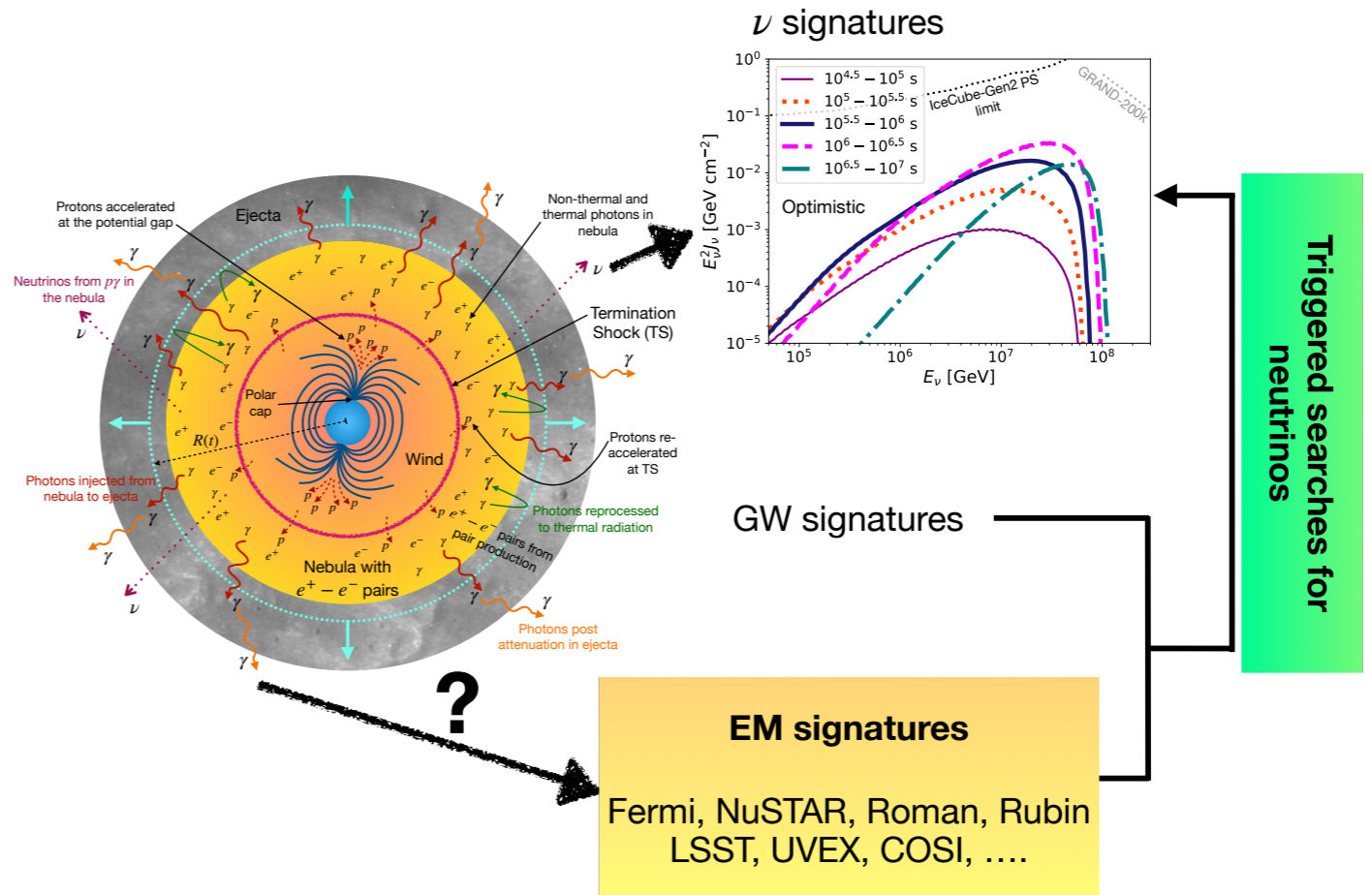
**Understanding:** (a) **Sites:** polar cap vs. equatorial current sheets and (b) **Mechanisms:** turbulence, magnetic reconnection, of **particle acceleration**

Magnetar origins of **superluminous supernovae (SLSNe)** and radio emissions from SLSNe, like, PTF10hgi

**Magnetar boosted kilonovae** (broadband emission from short GRBs, like, GRB 200522A) - detectable unto 100 Mpc

**Probing physics beyond the Standard Model:**

$10^{13}$  G –  $10^{15}$  G magnetic fields (axions, ....)



**EM-triggered searches for neutrinos**

**Thank You!**