



# Overview\* of the ATLAS heavy-ion program

Seminar on Particle Physics  
Phenomenology and Experiments

Dominik Derendarz

March 20, 2023

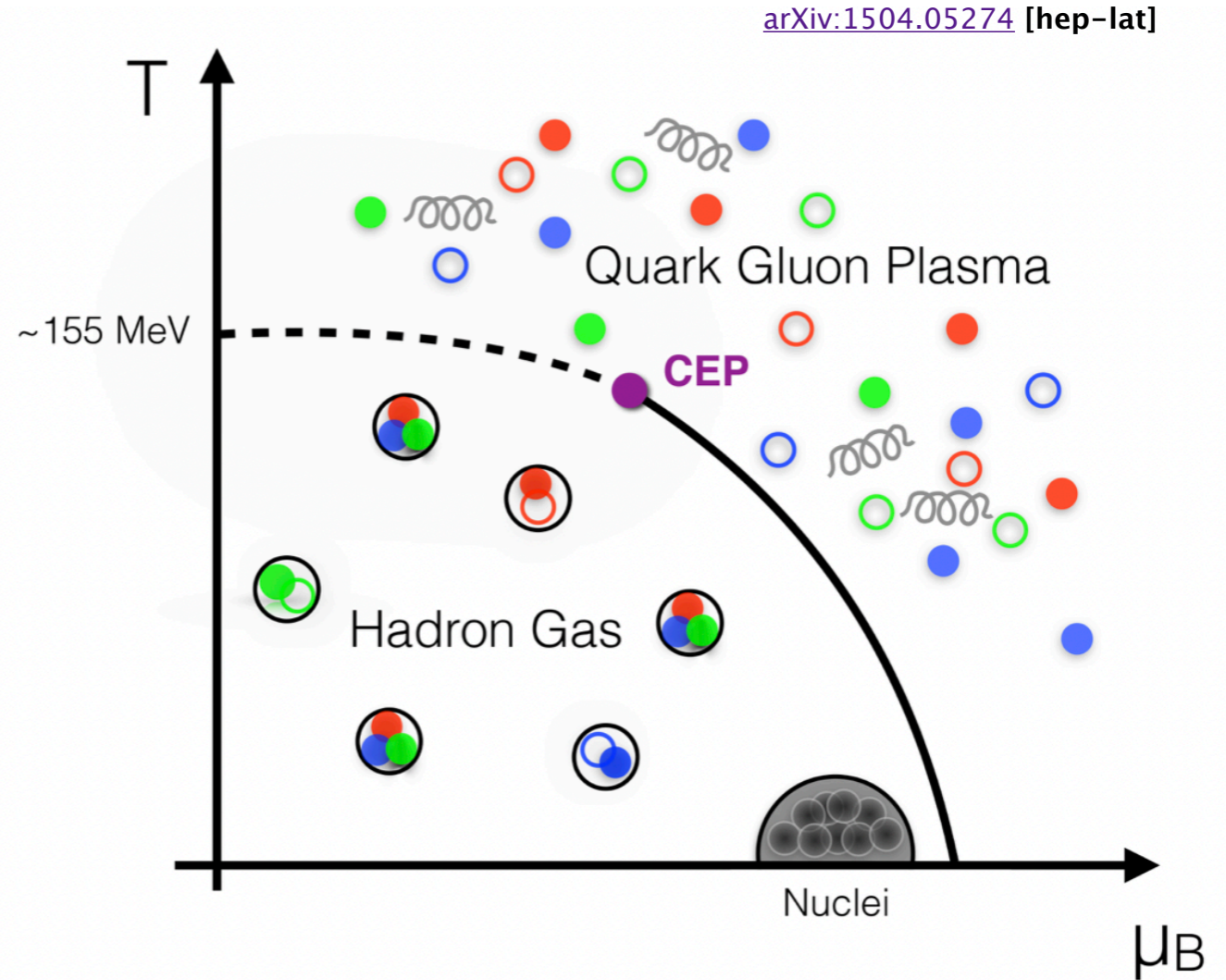
\*personal biased selection

# Heavy ion collisions - QCD laboratory

## Standard Model of Elementary Particles

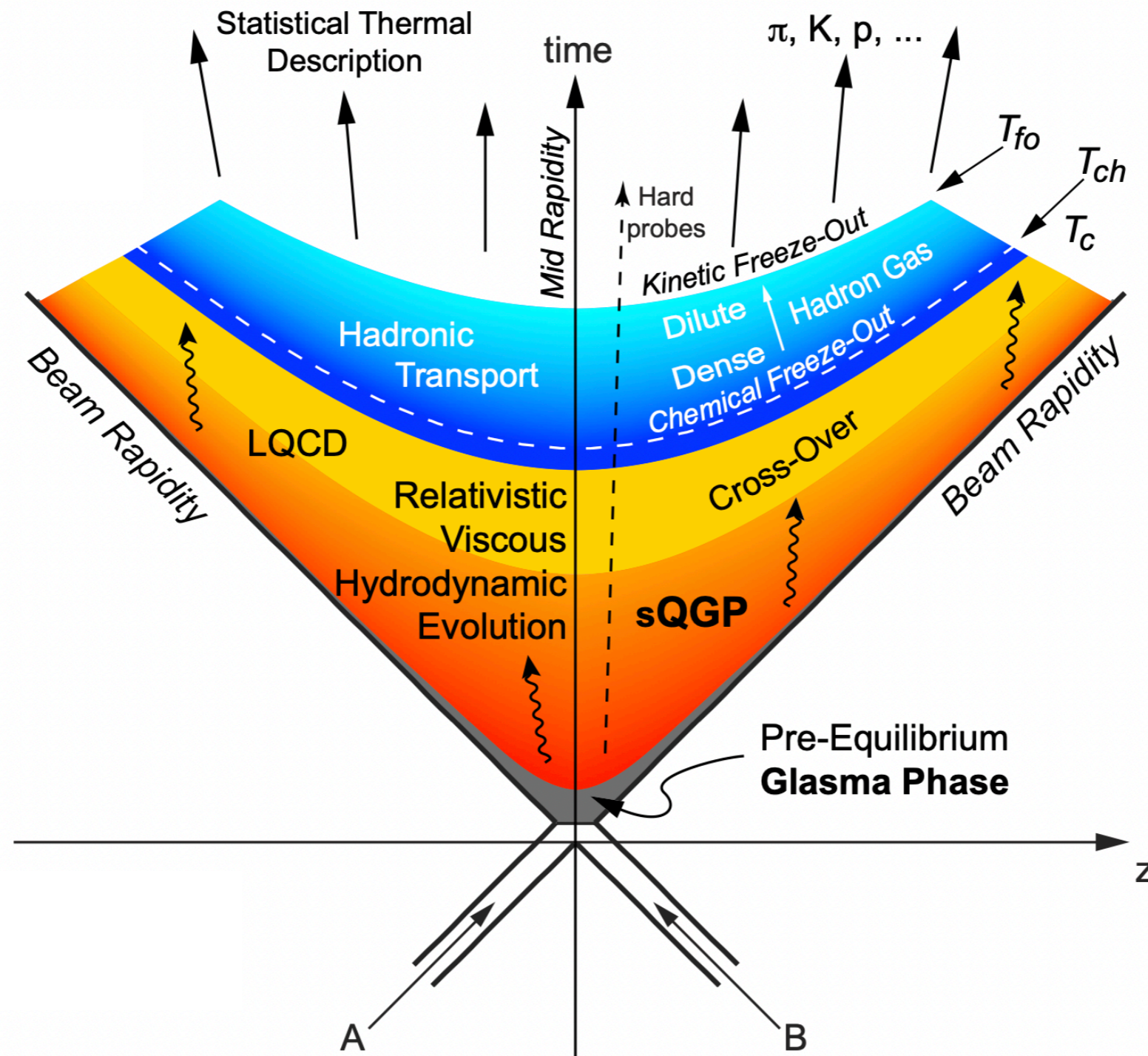
three generations of matter (fermions)			interactions / force carriers (bosons)			
	I	II	III			
QUARKS	mass $\approx 2.2 \text{ MeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ <b>u</b> up	mass $\approx 1.28 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ <b>c</b> charm	mass $\approx 173.1 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ <b>t</b> top	0 0 1 <b>g</b> gluon	SCALAR BOSONS	mass $\approx 124.97 \text{ GeV}/c^2$ 0 0 <b>H</b> higgs
	mass $\approx 4.7 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ <b>d</b> down	mass $\approx 96 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ <b>s</b> strange	mass $\approx 4.18 \text{ GeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ <b>b</b> bottom	0 0 1 <b><math>\gamma</math></b> photon		GAUGE BOSONS VECTOR BOSONS
	mass $\approx 0.511 \text{ MeV}/c^2$ charge $-1$ spin $\frac{1}{2}$ <b>e</b> electron	mass $\approx 105.66 \text{ MeV}/c^2$ charge $-1$ spin $\frac{1}{2}$ <b><math>\mu</math></b> muon	mass $\approx 1.7768 \text{ GeV}/c^2$ charge $-1$ spin $\frac{1}{2}$ <b><math>\tau</math></b> tau	mass $\approx 91.19 \text{ GeV}/c^2$ 0 1 <b>Z</b> Z boson		
mass $< 1.0 \text{ eV}/c^2$ charge 0 spin $\frac{1}{2}$ <b><math>\nu_e</math></b> electron neutrino	mass $< 0.17 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$ <b><math>\nu_\mu</math></b> muon neutrino	mass $< 18.2 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$ <b><math>\nu_\tau</math></b> tau neutrino	mass $\approx 80.39 \text{ GeV}/c^2$ $\pm 1$ 1 <b>W</b> W boson			

## Schematic QCD phase diagram



Phase transition at sufficiently high temperature

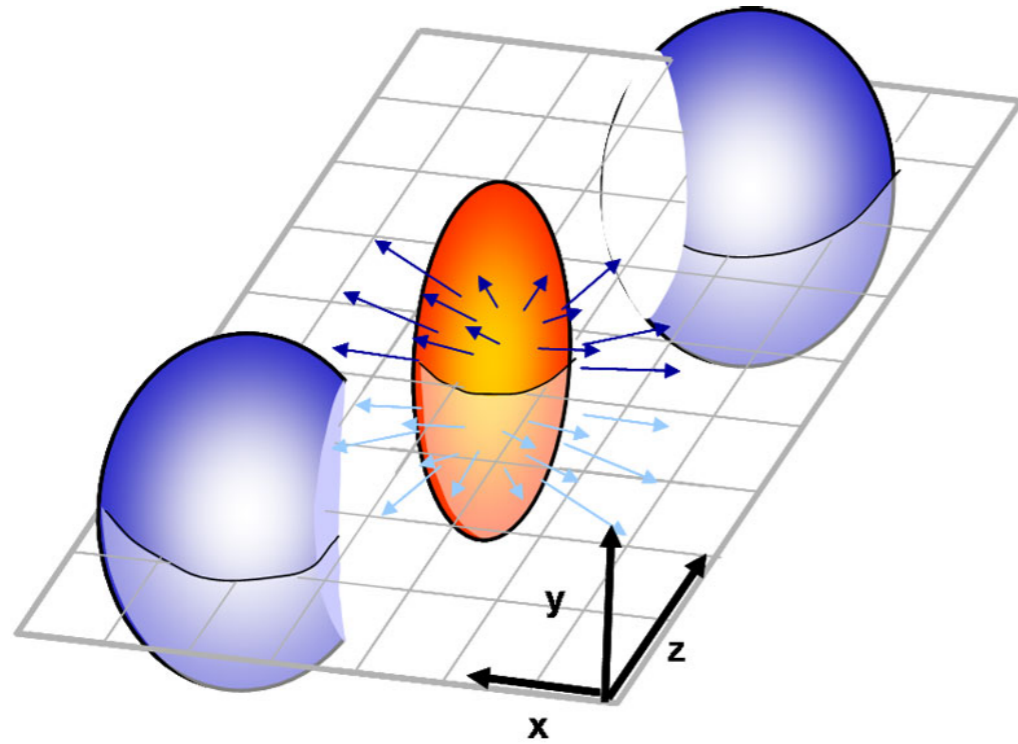
# Heavy ion collisions - QCD laboratory



Heavy ion collisions allows to reach high enough temperature to trigger the phase transition to QGP (for about few fermi or few  $10^{-23}$  s)

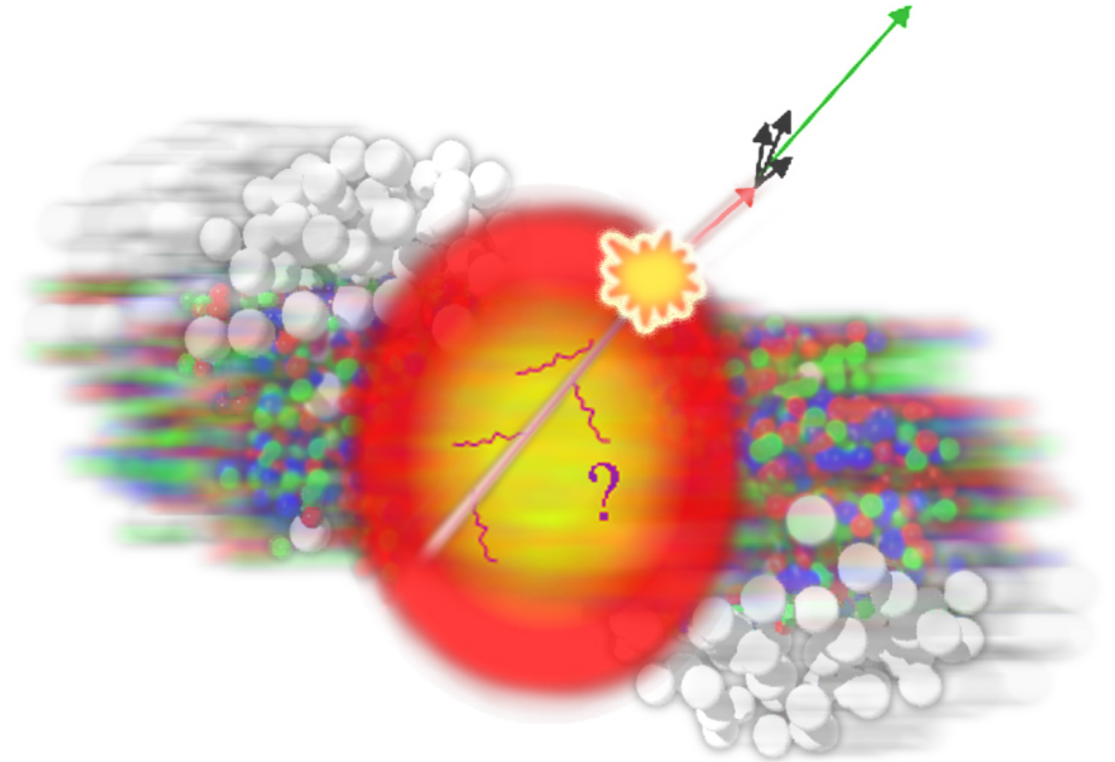
# Signatures of QGP formation

## Collective phenomena



- Initial **spatial** anisotropy is converted to **momentum** anisotropy of final state particles

## Jet quenching

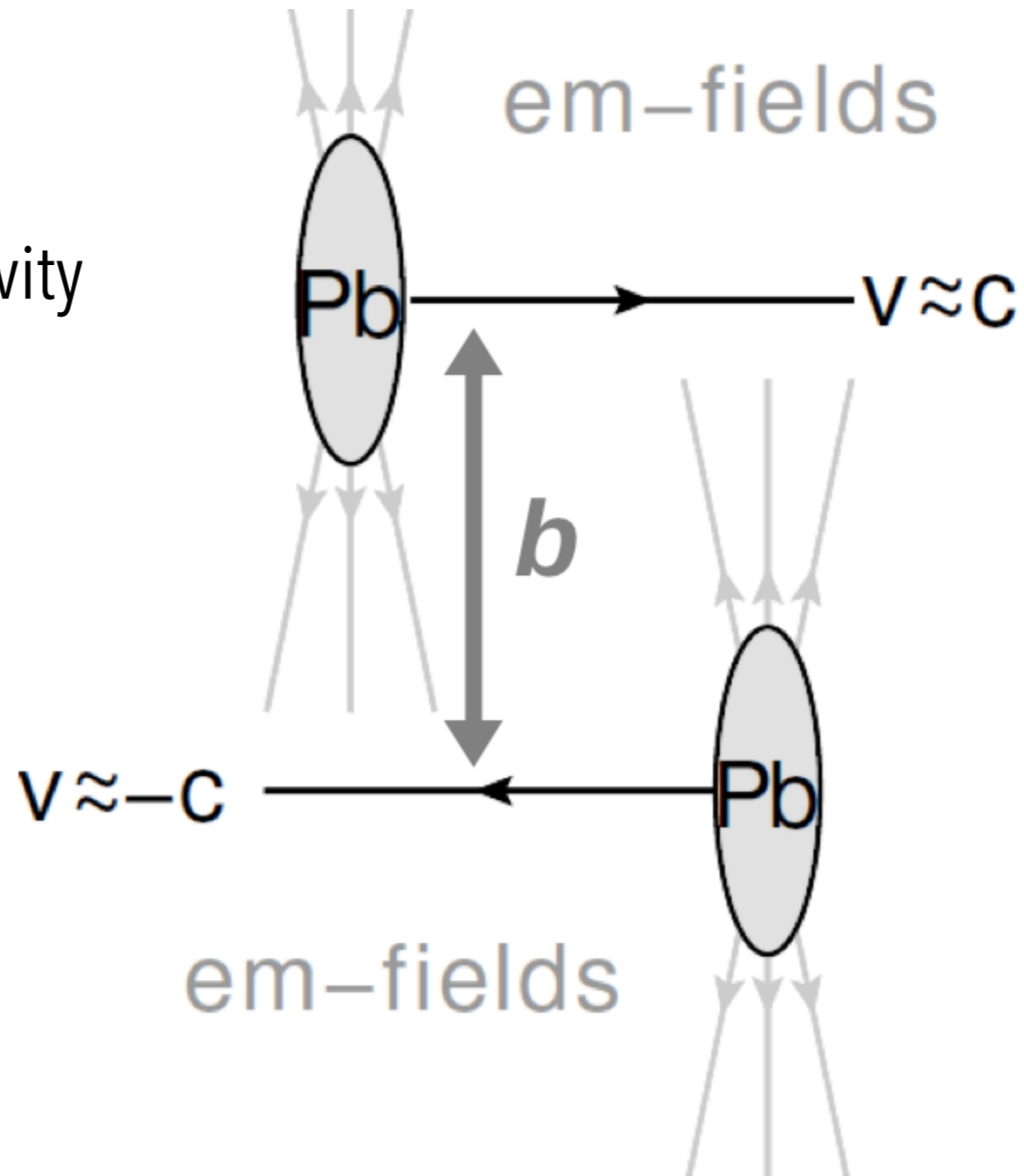


- Energy loss of high  $p_T$  parton due to the interactions with the QGP medium

# Heavy ion collisions - QED laboratory

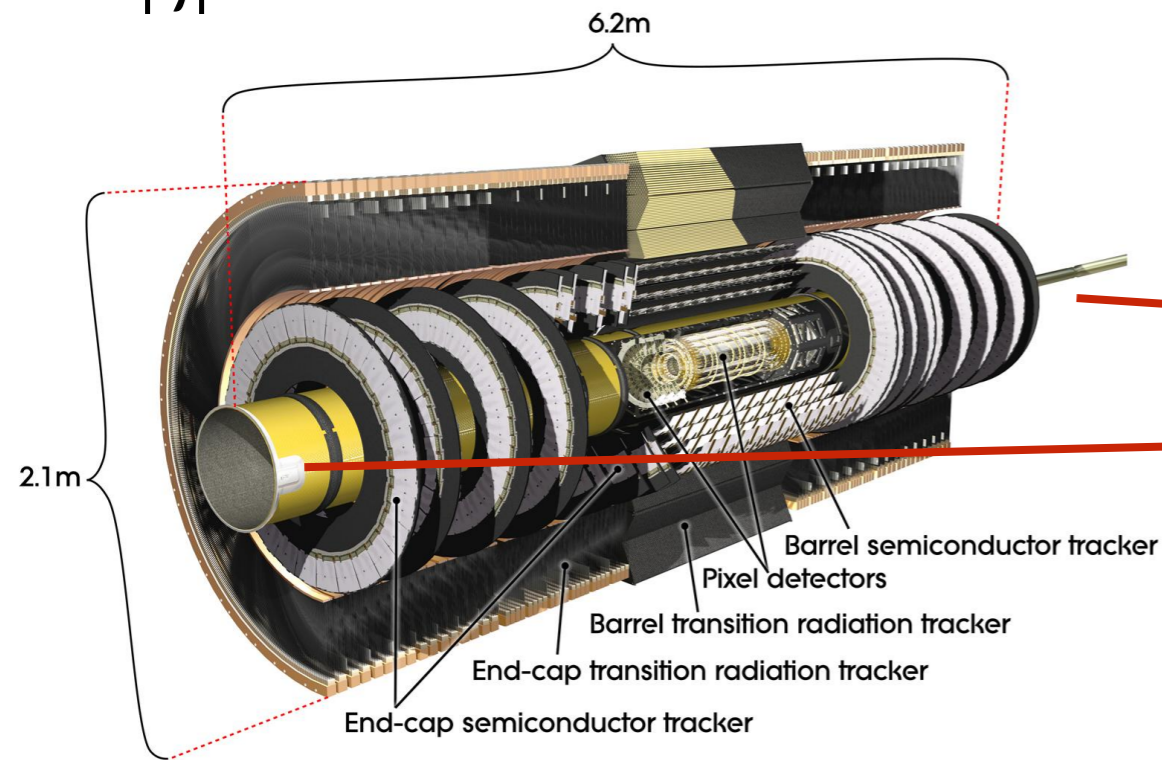
## Photon induced processes

QED laboratory with sensitivity for new physics.

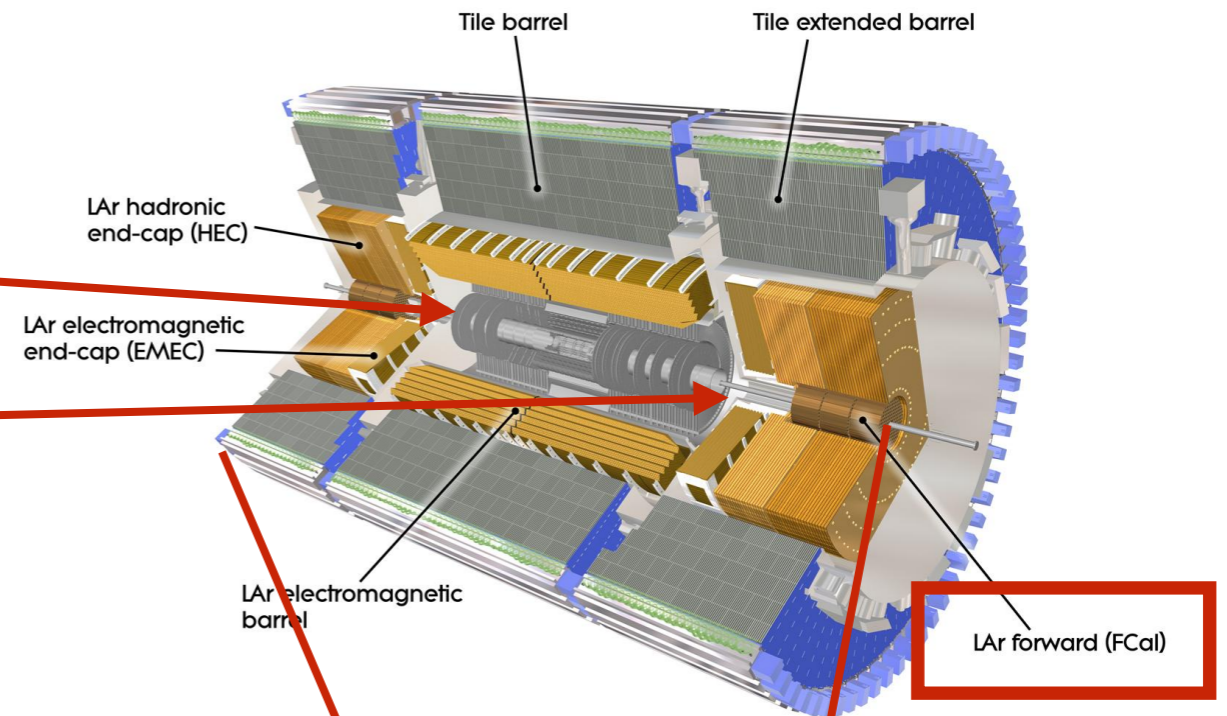


# ATLAS detector

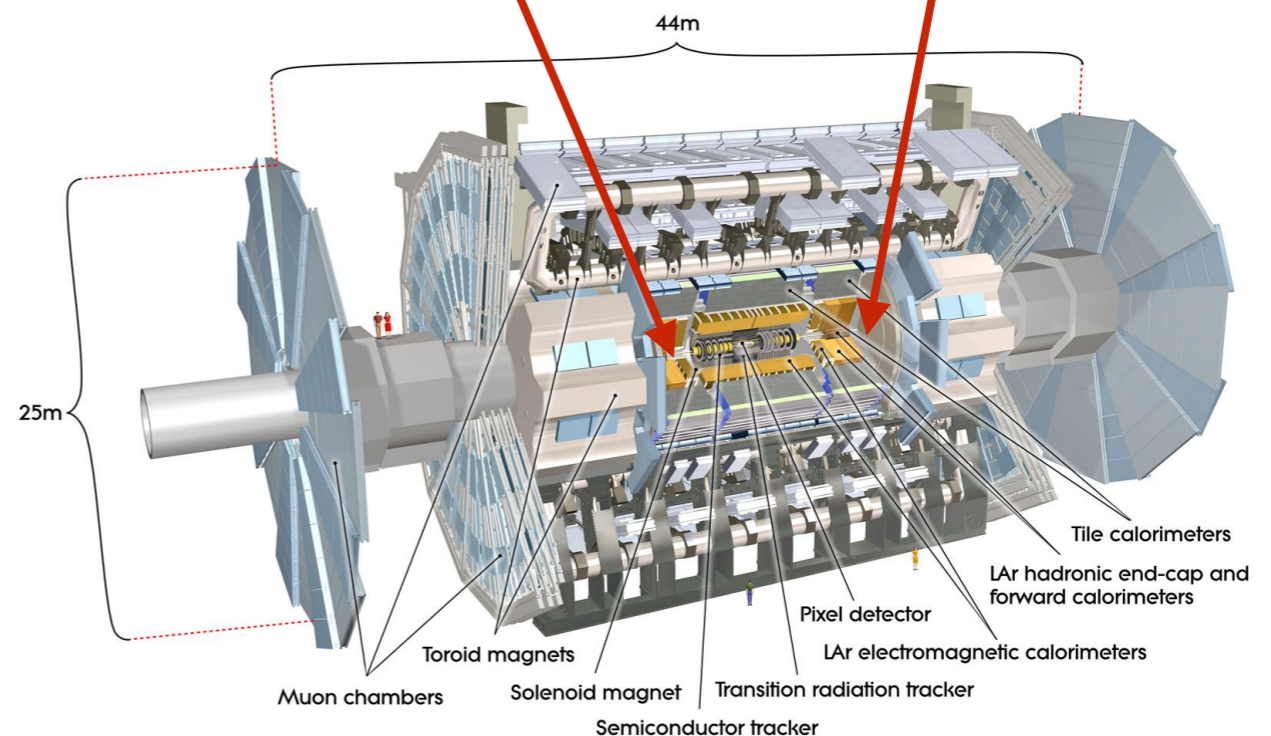
Tracker  $|\eta| < 2.5$



Calorimeters  $|\eta| < 4.9$



- And forward detectors located far far away from the interaction point ZDC (140m), AFP & ALFA (~240m)



# Heavy ion datasets

	System	Year	$\sqrt{s_{NN}}$ [TeV]	$L_{int}$
<b>Run1</b>	Pb+Pb	2010	2.76	$7 \mu\text{b}^{-1}$
	Pb+Pb	2011	2.76	$0.14 \text{ nb}^{-1}$
	pp	2012	8	$19.4 \text{ fb}^{-1}$
	pp	2013	2.76	$4 \text{ pb}^{-1}$
	p+Pb	2013	5.02	$29 \text{ nb}^{-1}$
<b>Run2</b>	pp	2015	5.02	$28 \text{ pb}^{-1}$
	Pb+Pb	2015	5.02	$0.49 \text{ nb}^{-1}$
	p+Pb	2016	5.02	$0.5 \text{ nb}^{-1}$
	p+Pb	2016	8.16	$0.16 \text{ pb}^{-1}$
	Xe+Xe	2017	5.44	$3 \mu\text{b}^{-1}$
	pp	2017	5.02	$270 \text{ pb}^{-1}$
	Pb+Pb	2018	5.02	$1.76 \text{ nb}^{-1}$

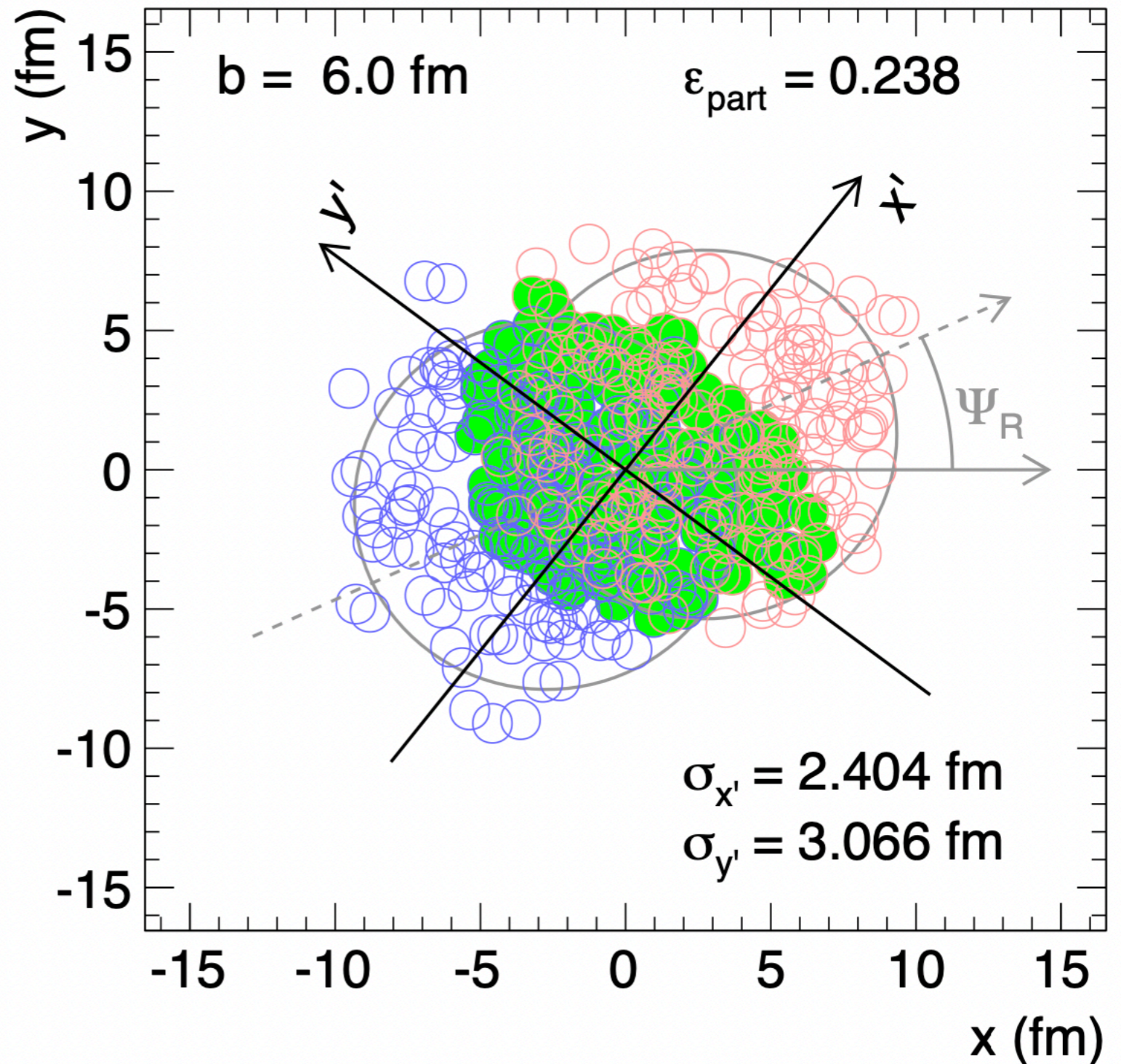
For **Run3** (2022-2024) expected:

- $\sim 6 \text{ nb}^{-1}$  of Pb+Pb
- p+Pb  $\sim X \text{ nb}^{-1}$
- Short pilot run with 0+0 and p+0

# Centrality of heavy ion collision

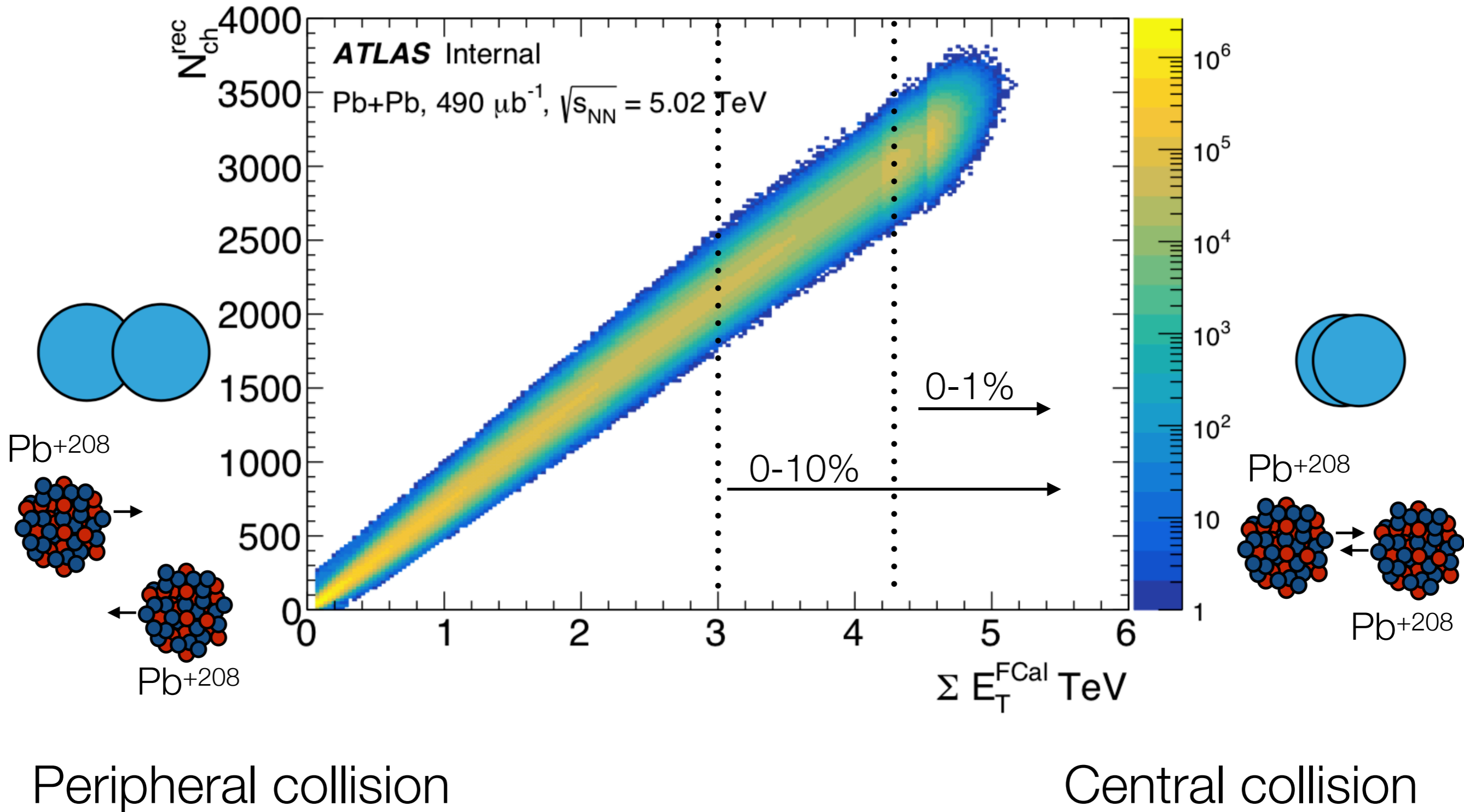
[arXiv:1204.1409](https://arxiv.org/abs/1204.1409) [nucl-ex]

- One nucleon can interact many times due to the thickness of the other nuclei that he see
- $\langle N_{\text{part}} \rangle$ ,  $\langle N_{\text{coll}} \rangle$  estimated based on MC Glauber fits to data





# Centrality of heavy ion collision



# Bare eye jet quenching from 2010

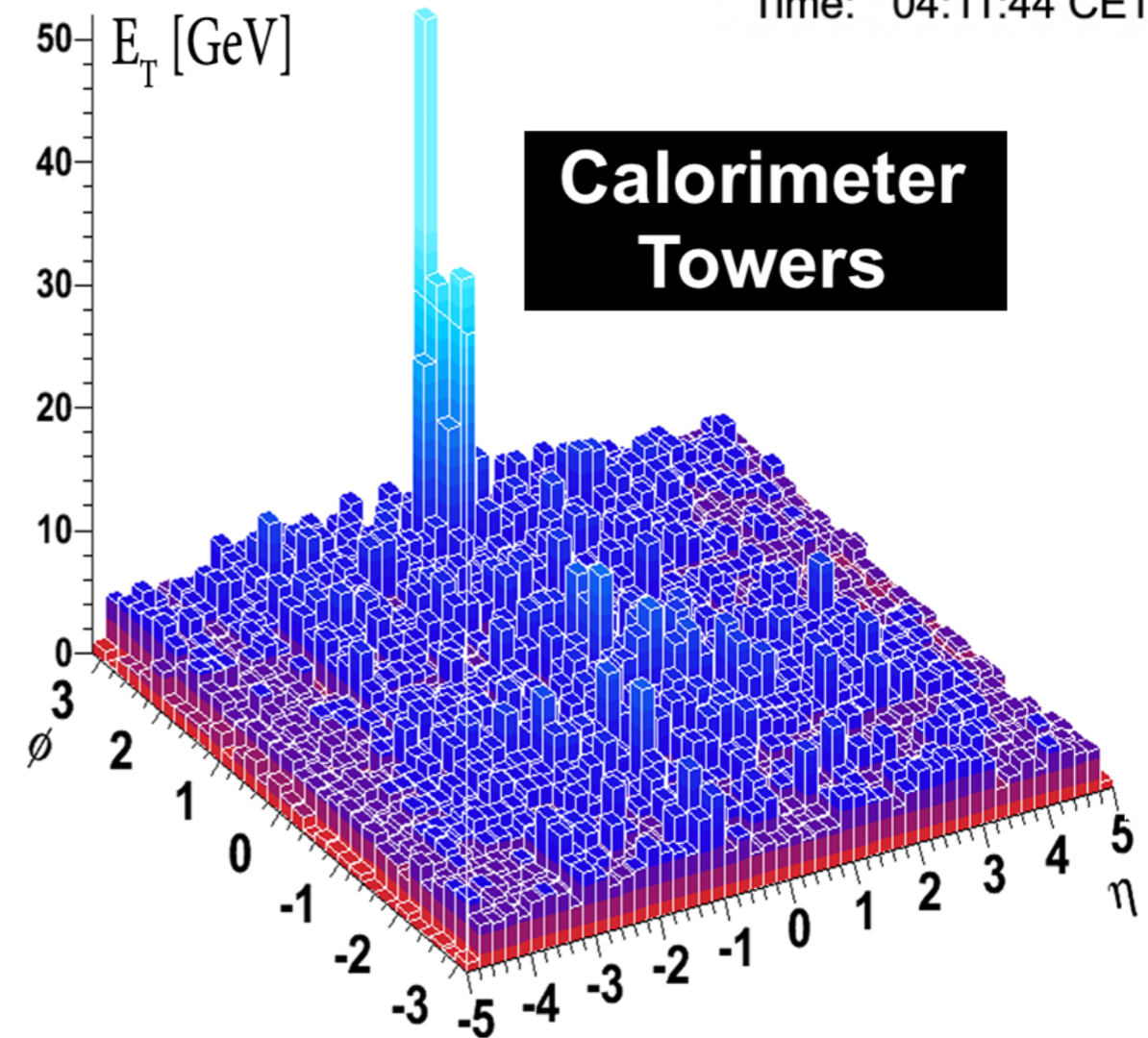
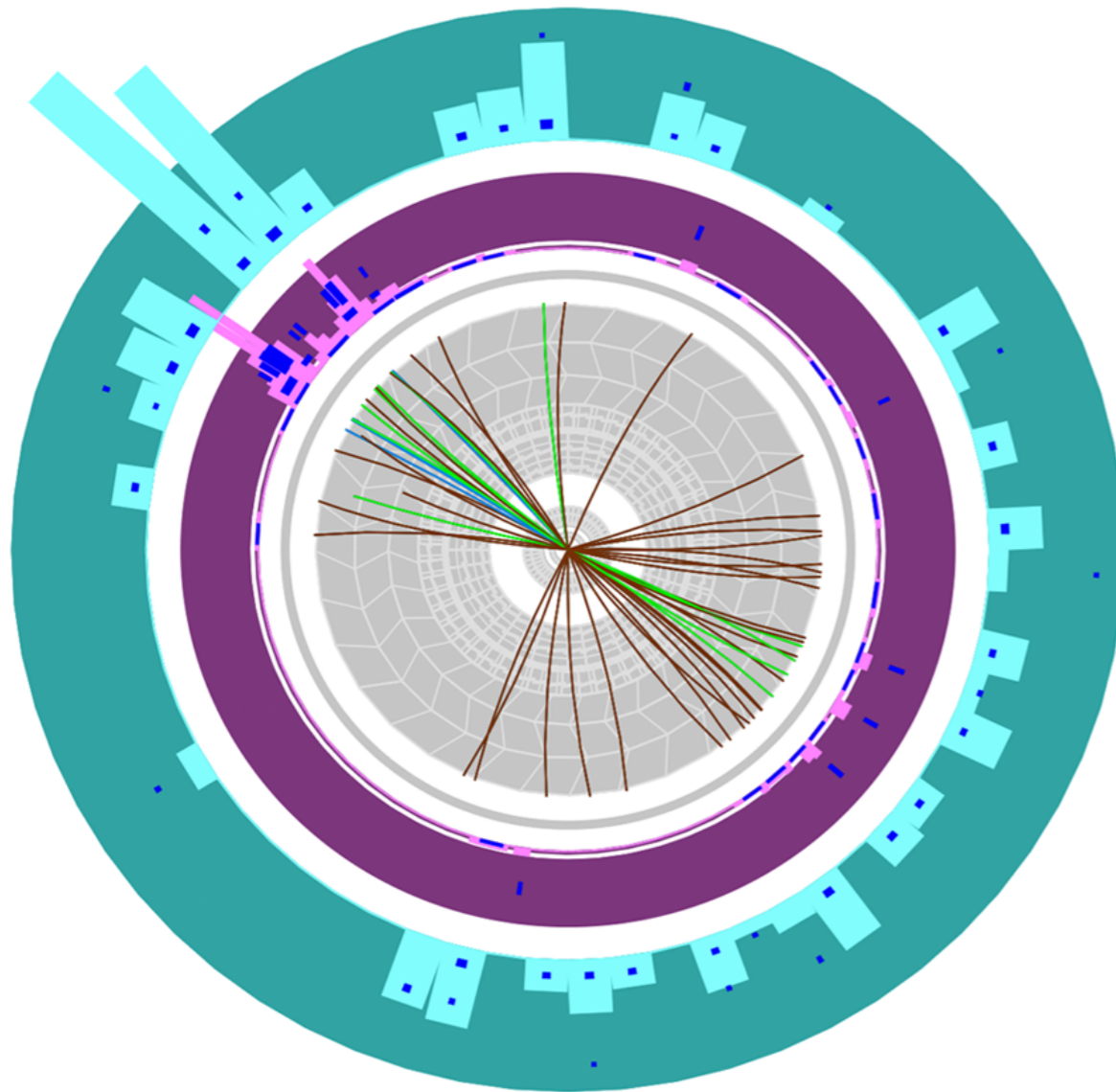
## ATLAS

Run: 169045

Event: 1914004

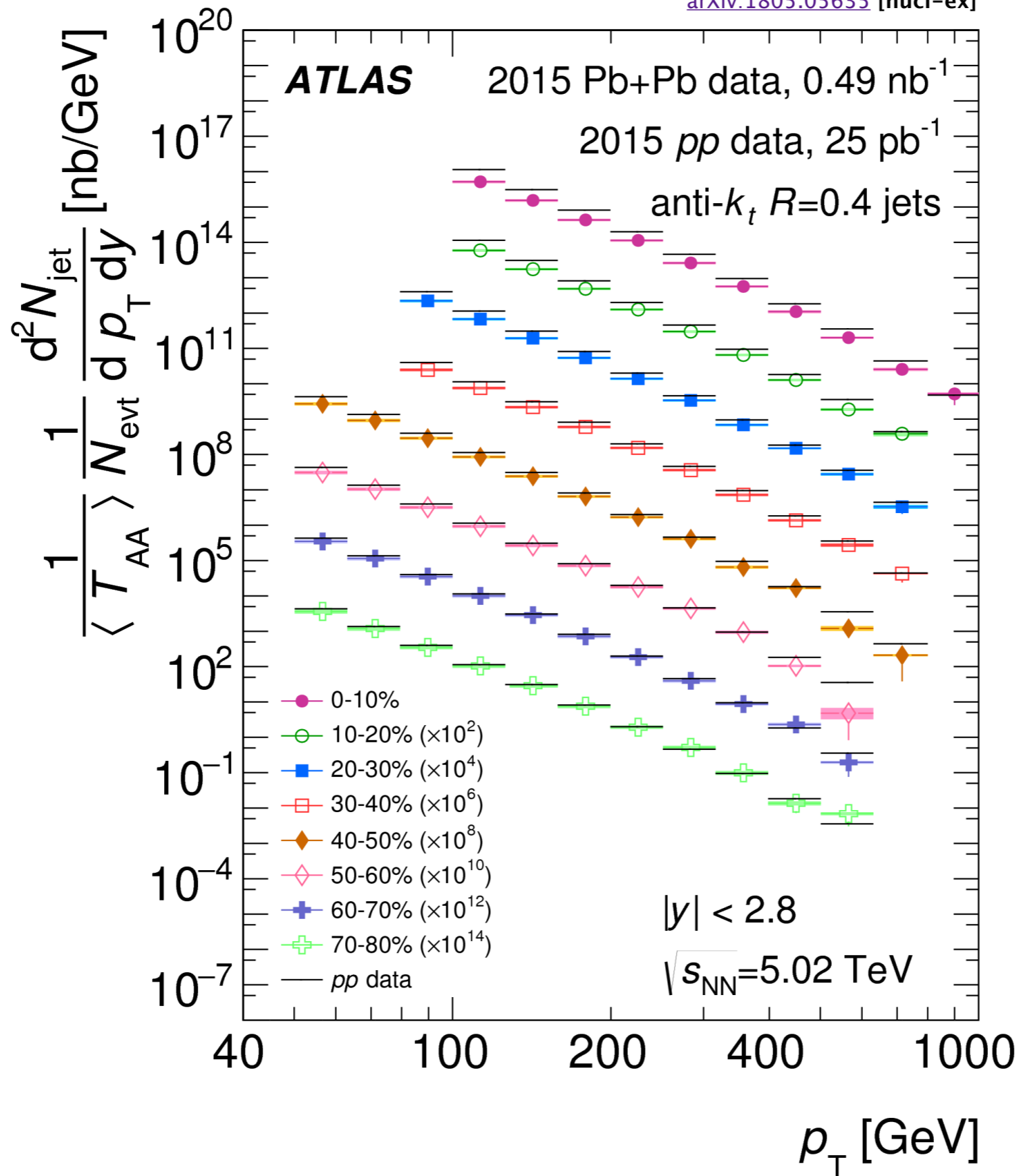
Date: 2010-11-12

Time: 04:11:44 CET



# Inclusive jet spectra

arXiv:1805.05635 [nucl-ex]



# Nuclear modification factor $R_{AA}$

$$R_{AA} = \frac{\frac{1}{N_{\text{evt}}} \left. \frac{d^2 N_{\text{jet}}}{dp_T dy} \right|_{\text{cent}}}{\langle T_{AA} \rangle \left. \frac{d^2 \sigma_{\text{jet}}}{dp_T dy} \right|_{pp}}$$

Scaled Pb+Pb

Reference p+p

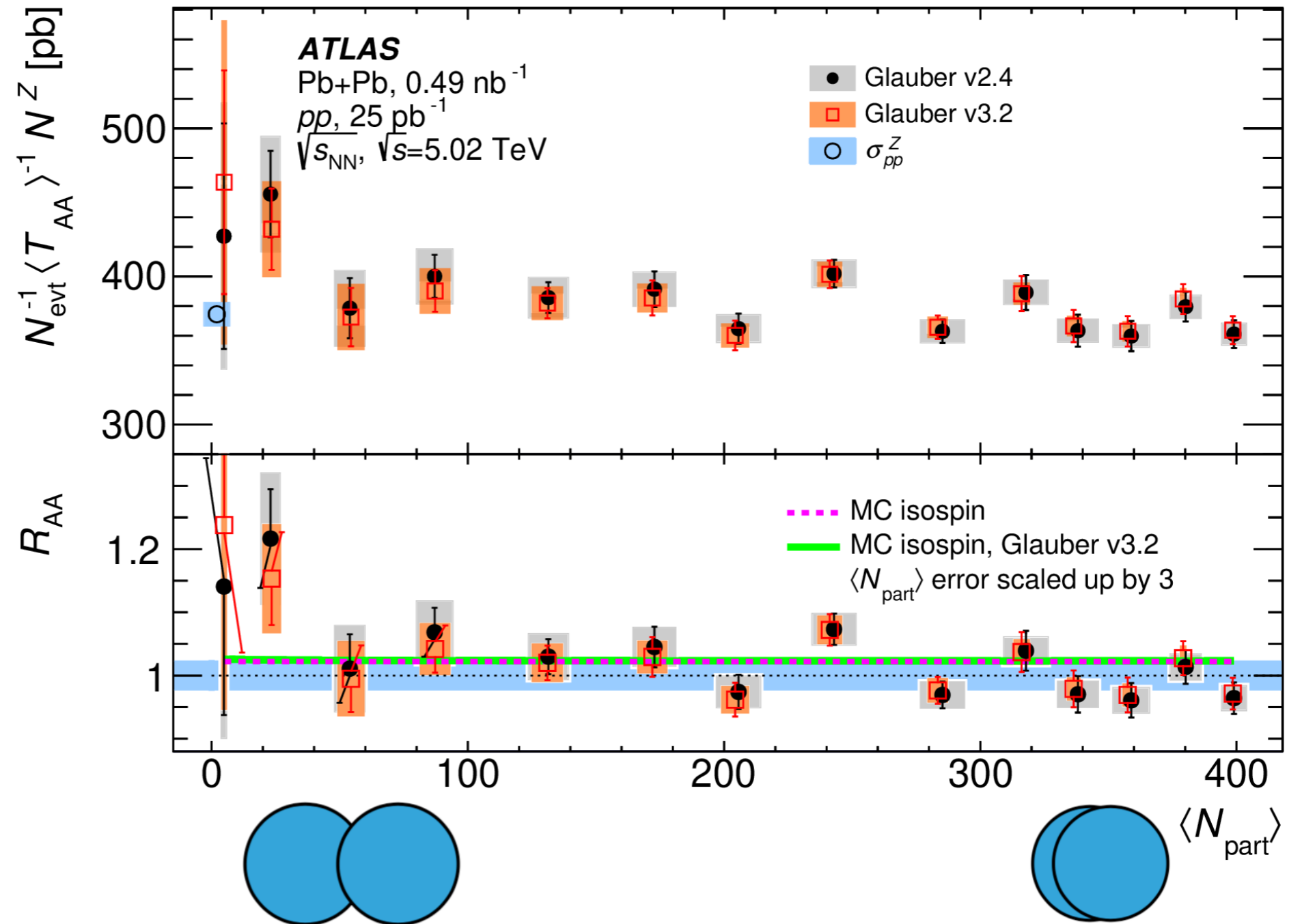
Where nuclear overlap function  $\langle T_{AA} \rangle$

$$\langle T_{AB} \rangle_f = \langle N_{\text{coll}} \rangle_f / \sigma_{\text{inel}}^{\text{NN}}$$

calculated in the Glauber MC approach ([arXiv:nucl-ex/0701025](https://arxiv.org/abs/nucl-ex/0701025))

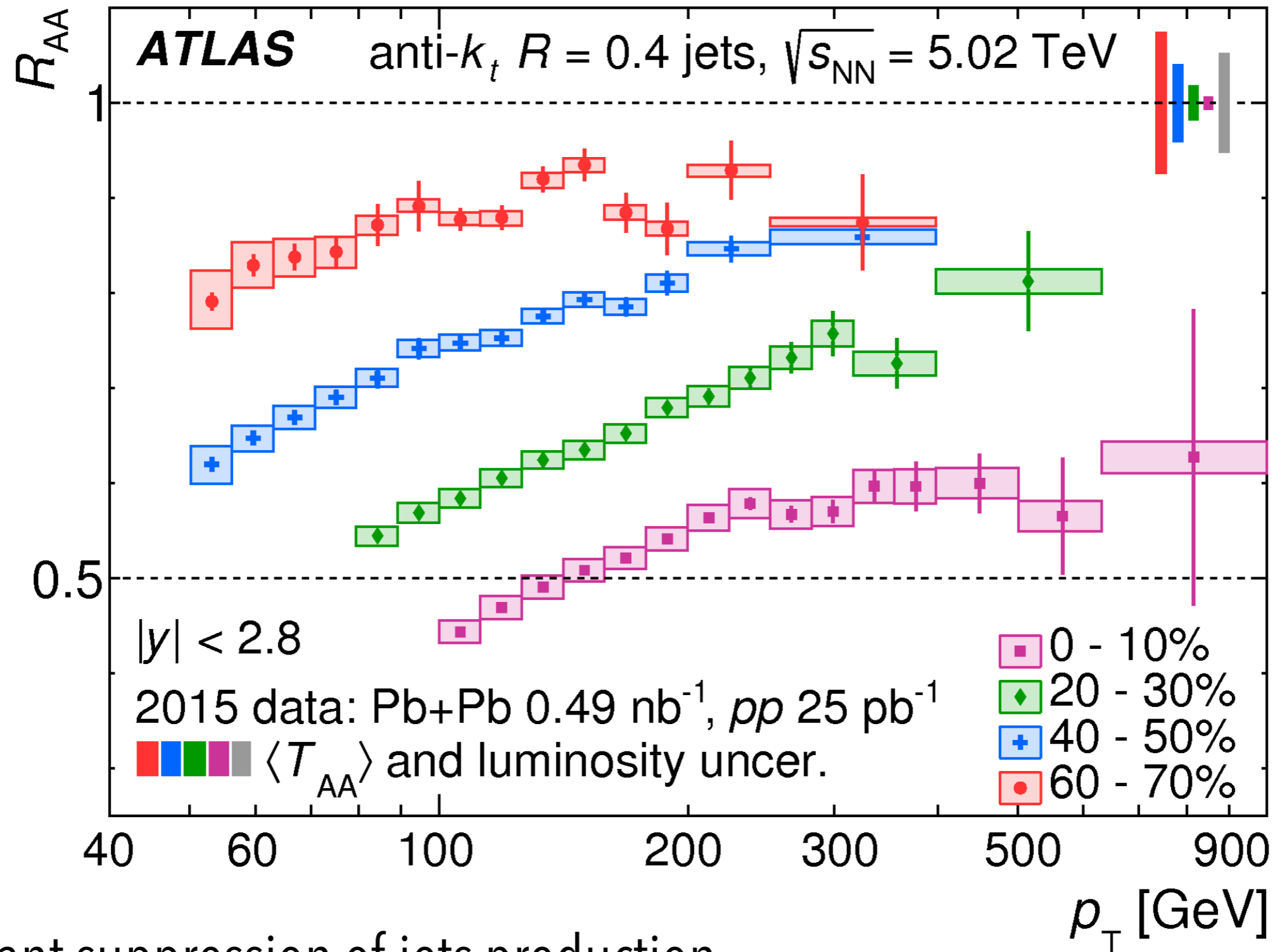
# Calibration of the $T_{AA}$ scaling

arXiv:1910.13396 [nucl-ex]



Probes that are not interacting with the QGP medium ( $W, Z, \gamma$ ) follows the  $\langle T_{AA} \rangle$  scaling.

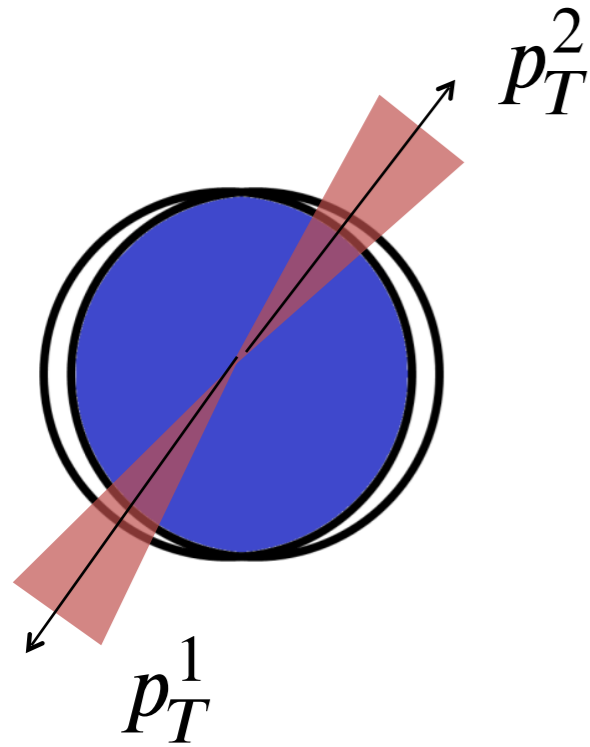
arXiv:1805.05635 [nucl-ex]



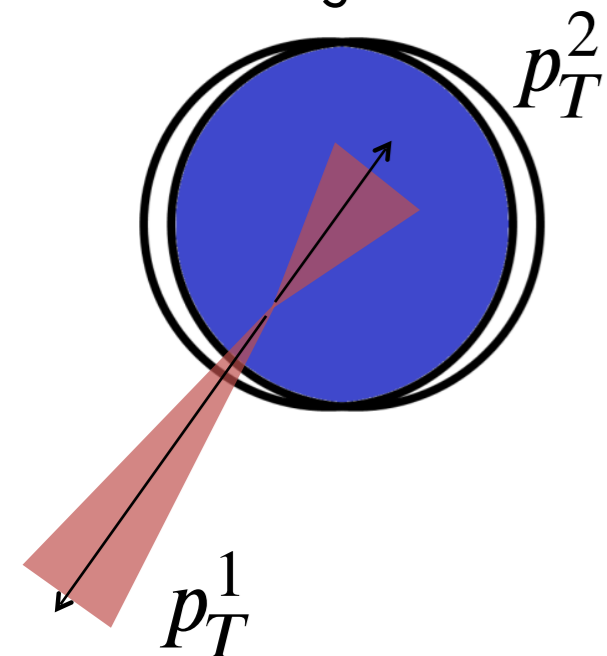
Significant suppression of jets production  
in the entire kinematic range

# Dijet asymmetry

Balanced pair of jets produced near the centre of the collisions zone



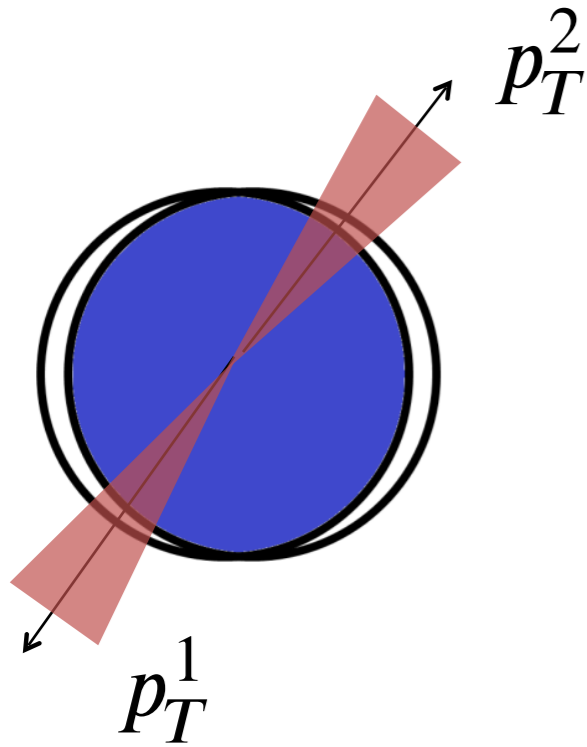
Unbalanced pair of jets produced near the edge of the collisions zone



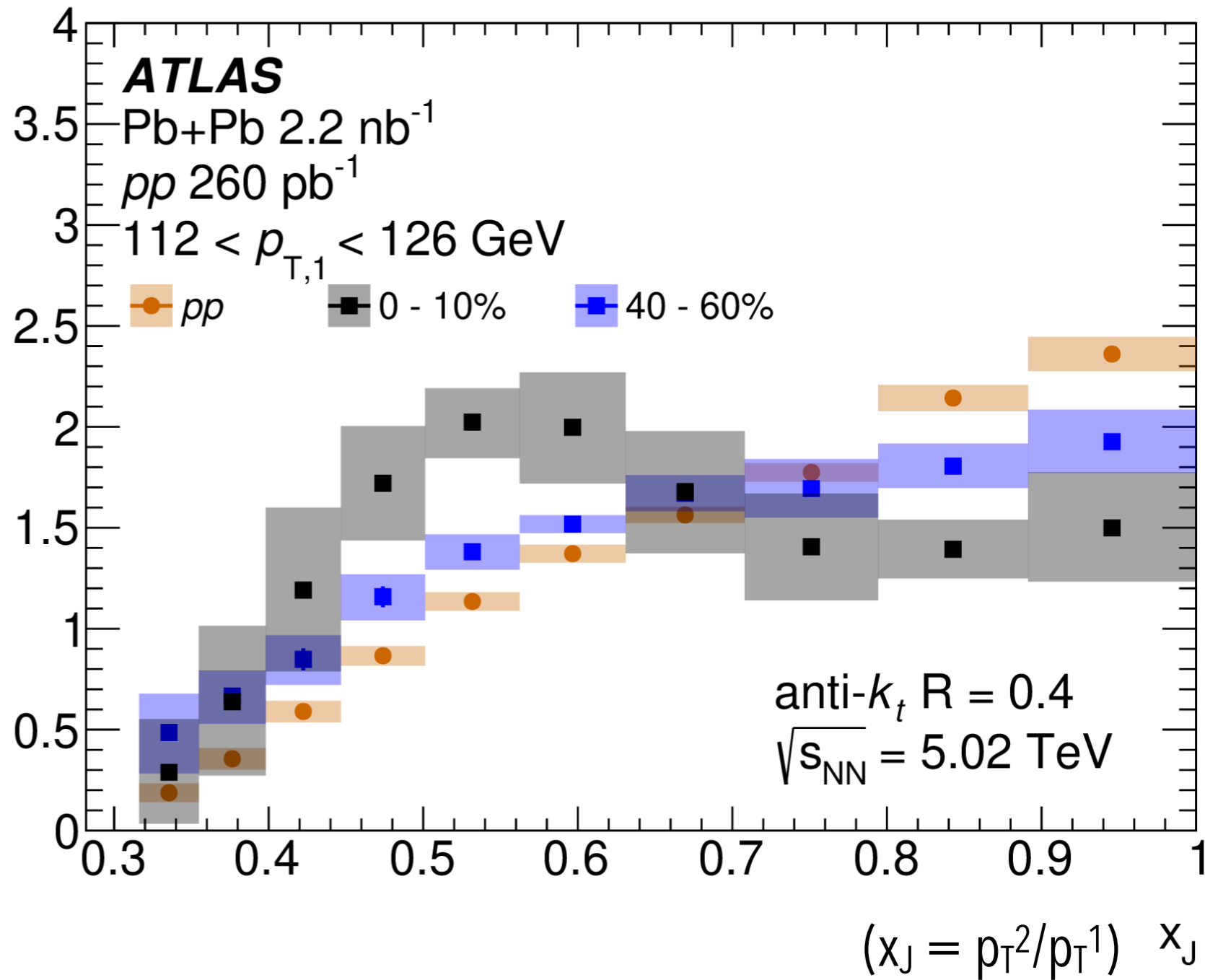
# Dijet asymmetry

[arXiv:2205.00682](https://arxiv.org/abs/2205.00682) [nucl-ex]

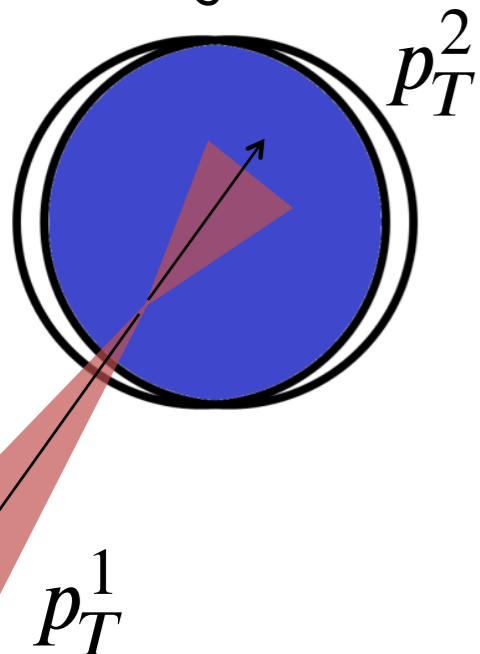
Balanced pair of jets produced near the centre of the collisions zone



$$\frac{1}{N_{\text{pair}}} \frac{dN_{\text{pair}}}{dx_J}$$



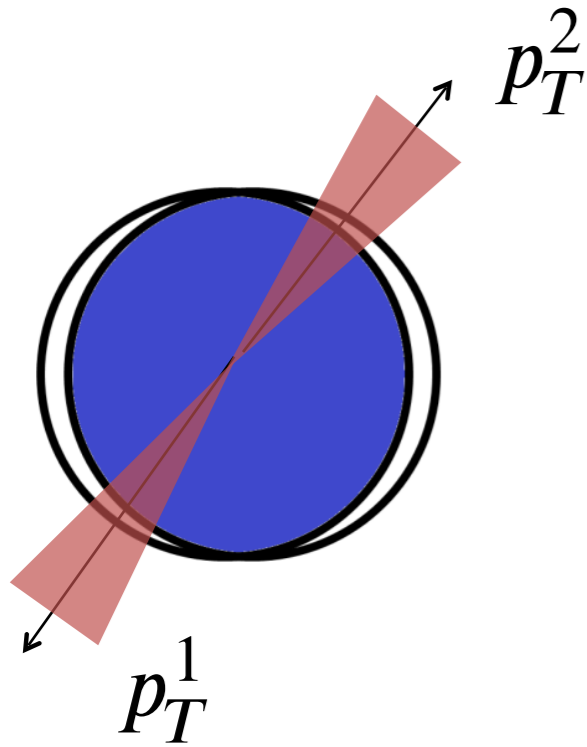
Unbalanced pair of jets produced near the edge of the collisions zone



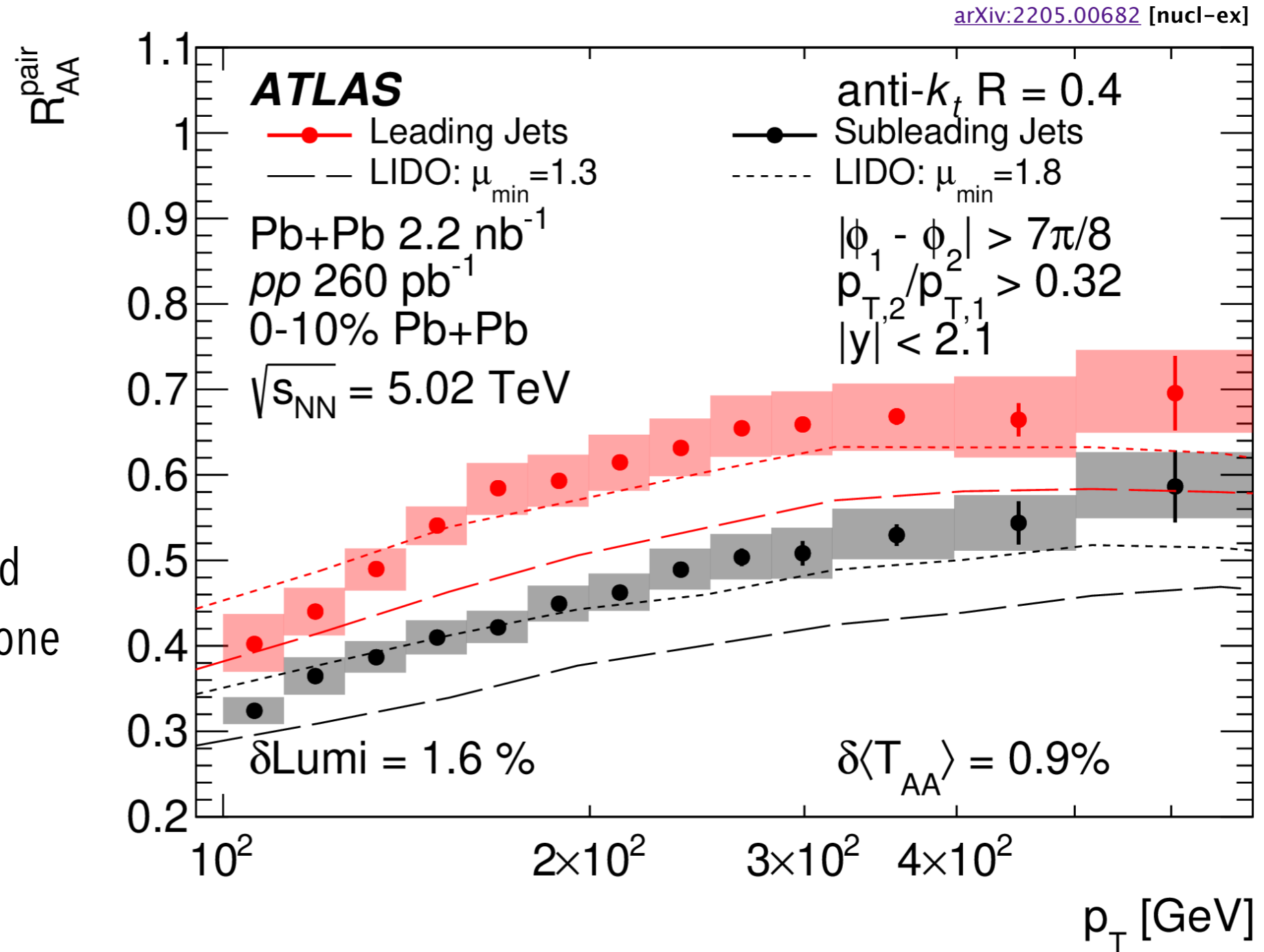
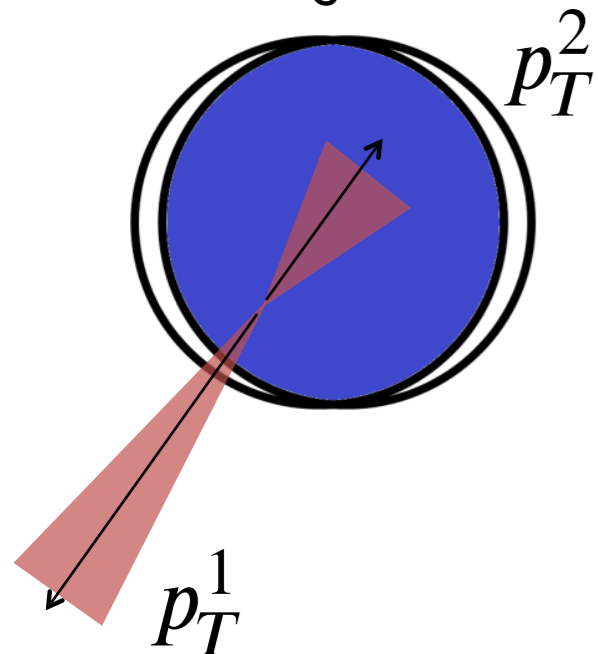


# Dijet asymmetry

Balanced pair of jets produced near the centre of the collisions zone

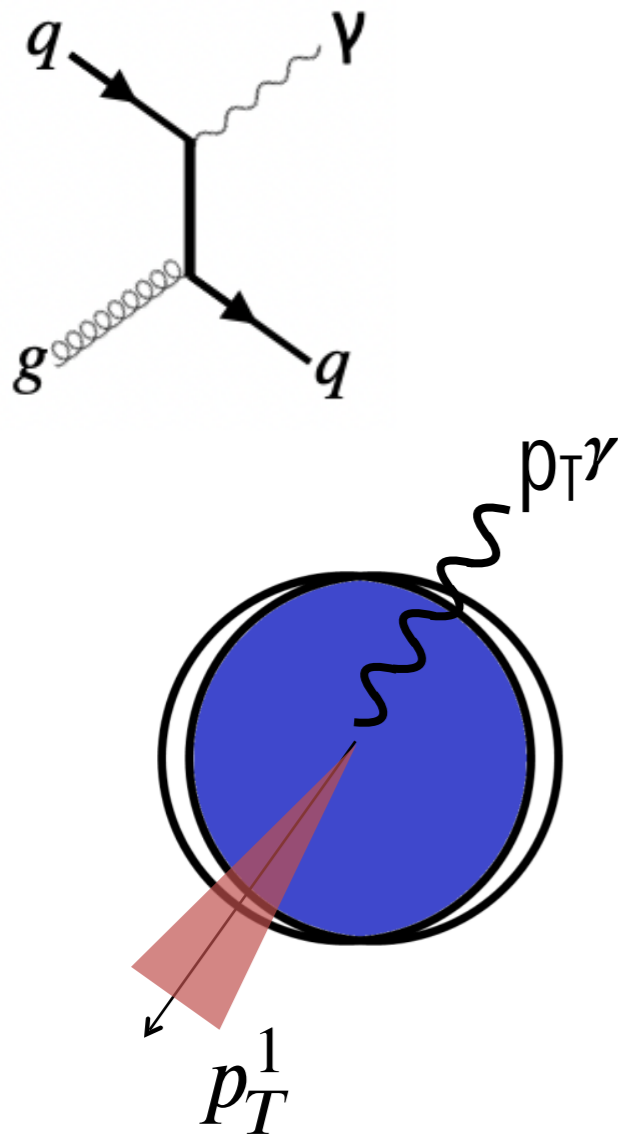


Unbalanced pair of jets produced near the edge of the collisions zone



# Photon tagged jet $R_{AA}$

<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/HION-2022-14/>



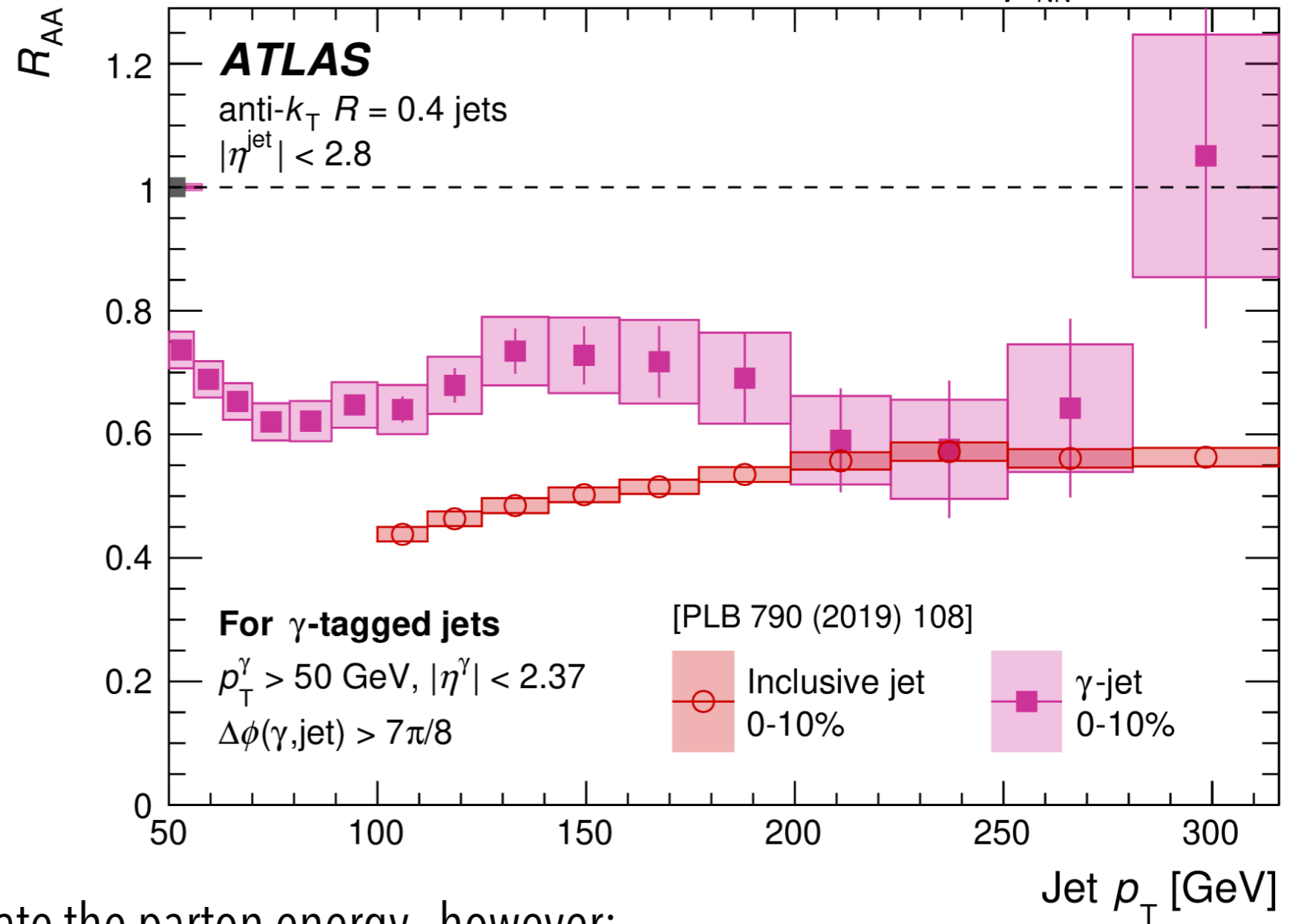
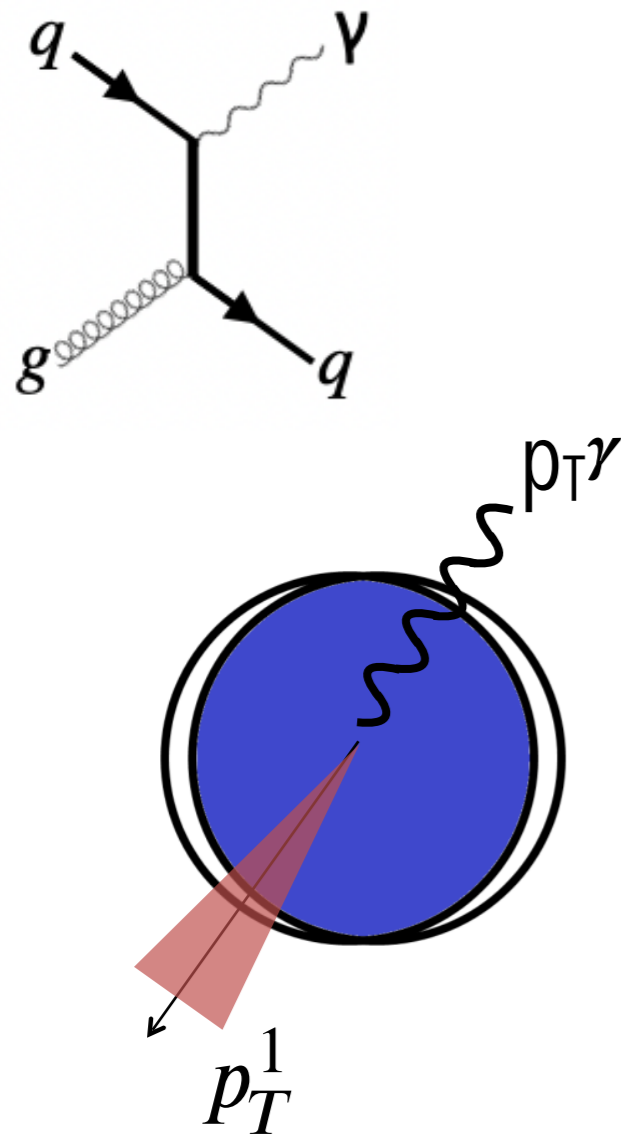
Photons and Z's can help to calibrate the parton energy, however:

- Rare process compared to the dijets production
- Flavour fraction differs compared to dijets

# Photon tagged jet $R_{AA}$

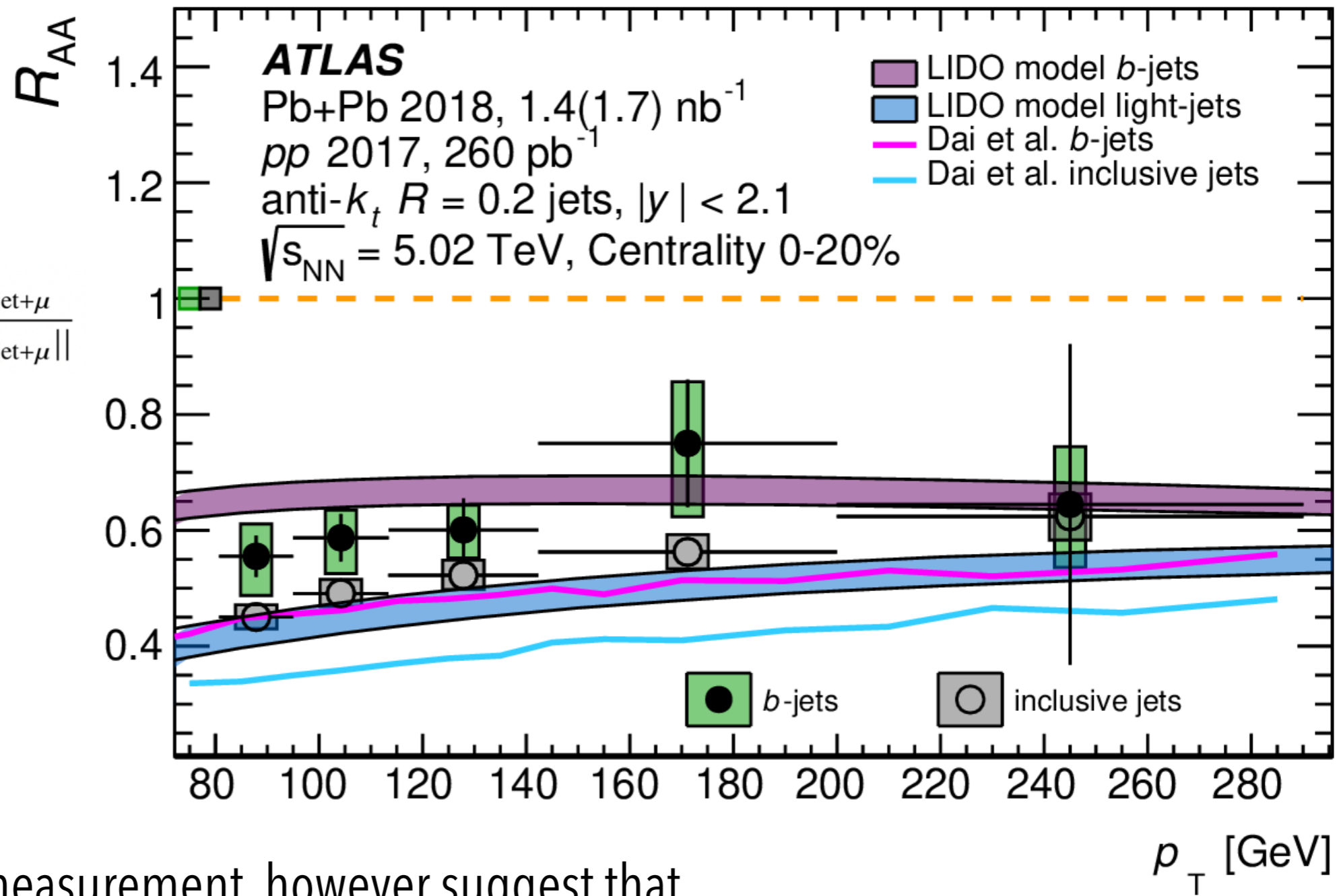
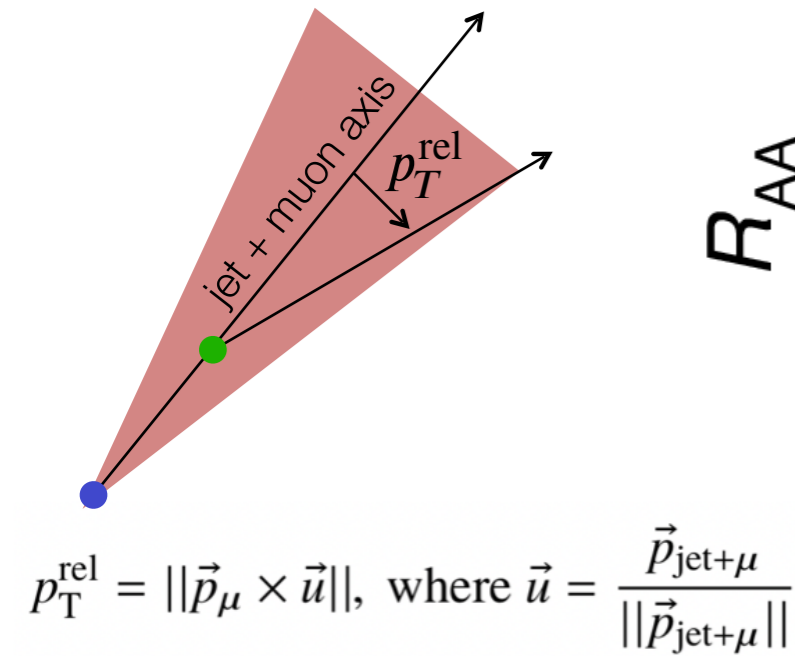
<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/HION-2022-14/>

2018 Pb+Pb 1.7 nb<sup>-1</sup>, 2017 pp 260 pb<sup>-1</sup>,  $\sqrt{s_{NN}} = 5.02$  TeV



Photons and Z's can help to calibrate the parton energy, however:

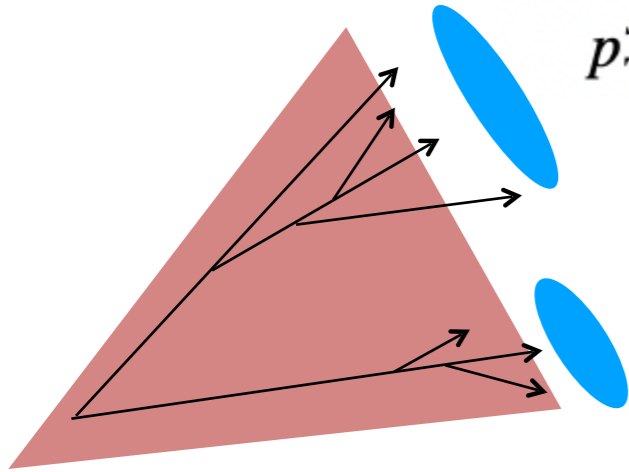
- Rare process compared to the dijets production
- Flavour fraction differs compared to dijets - **sample dominated by quark initiated jets**



Particularly difficult measurement, however suggest that ***b-jets are less*** suppressed than **inclusive jets**

# Energy loss dependence on the jet substructure

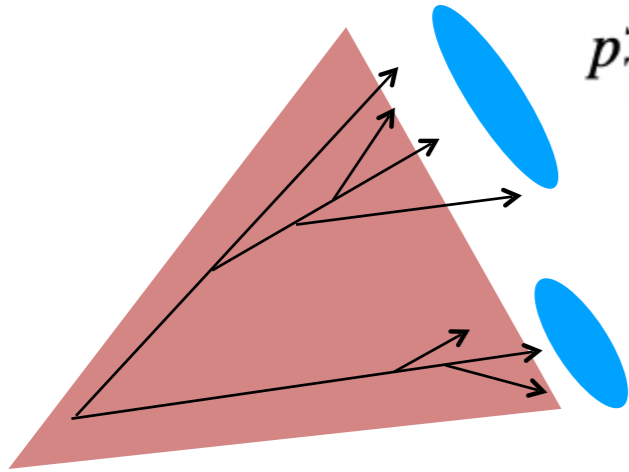
$$\frac{\min(p_T^{sj_1}, p_T^{sj_2})}{p_T^{sj_1} + p_T^{sj_2}} > z_{\text{cut}} \left( \frac{\Delta R_{12}}{R} \right)^\beta$$



- The soft drop grooming procedure ( $\beta = 0$  and  $\mathbf{z}_{\text{cut}} = 0.2$ ) is applied to determine opening angle of the hardest splitting ( $\mathbf{r}_{\mathbf{g}}$ )
- Jets re-clustered using Track-CaloClusters as constituents to improve angular resolution

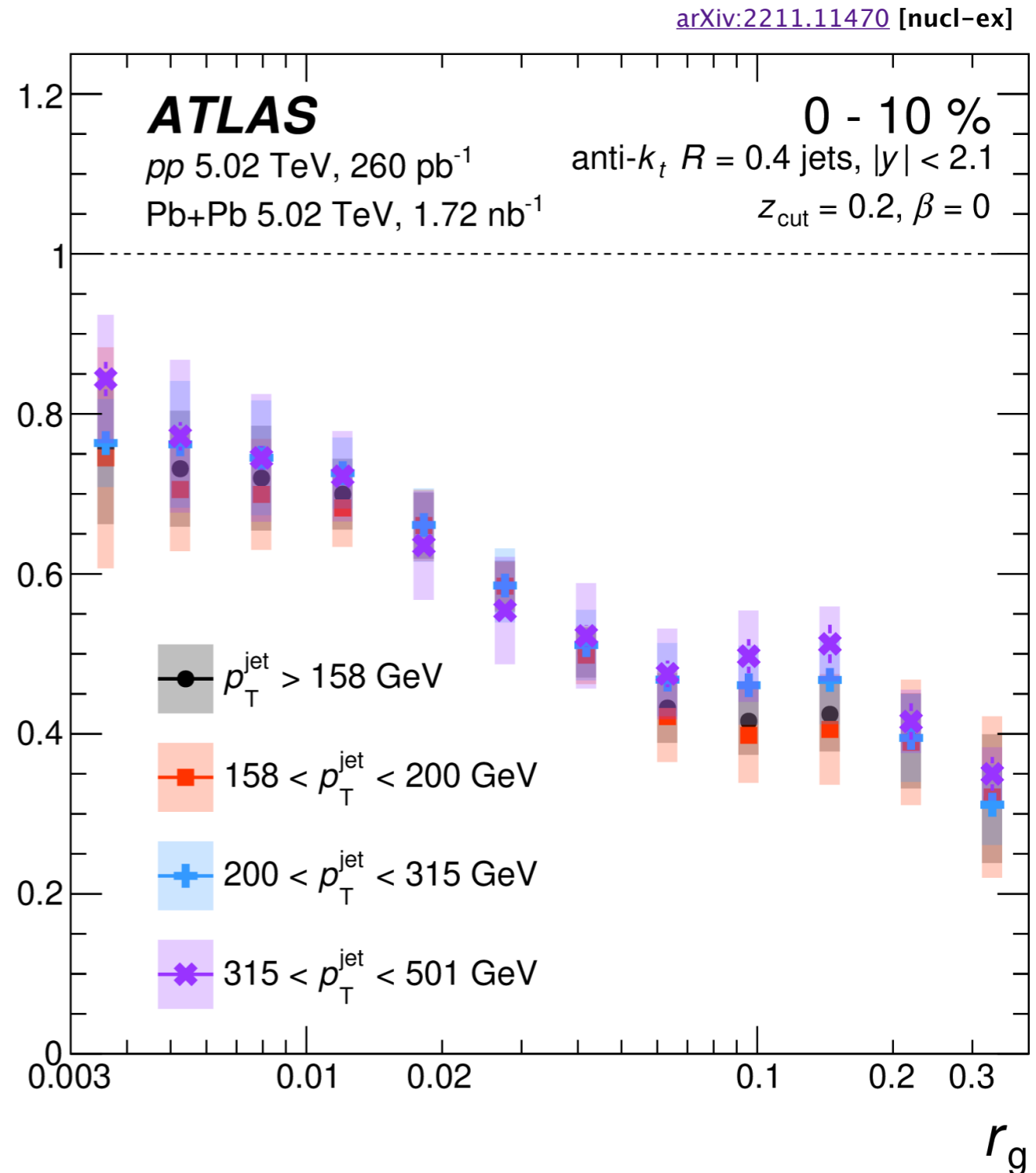
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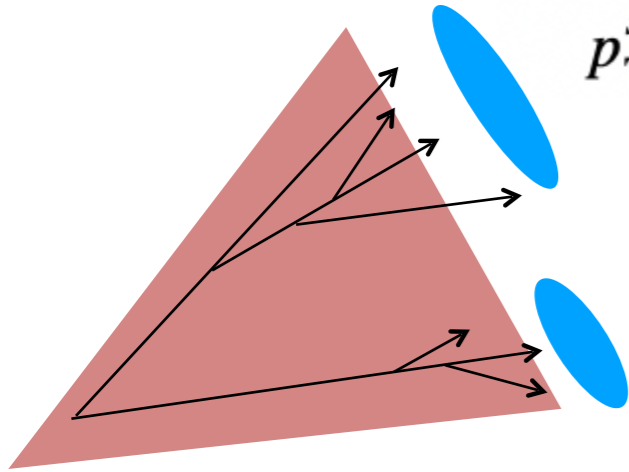
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$R_{AA}$



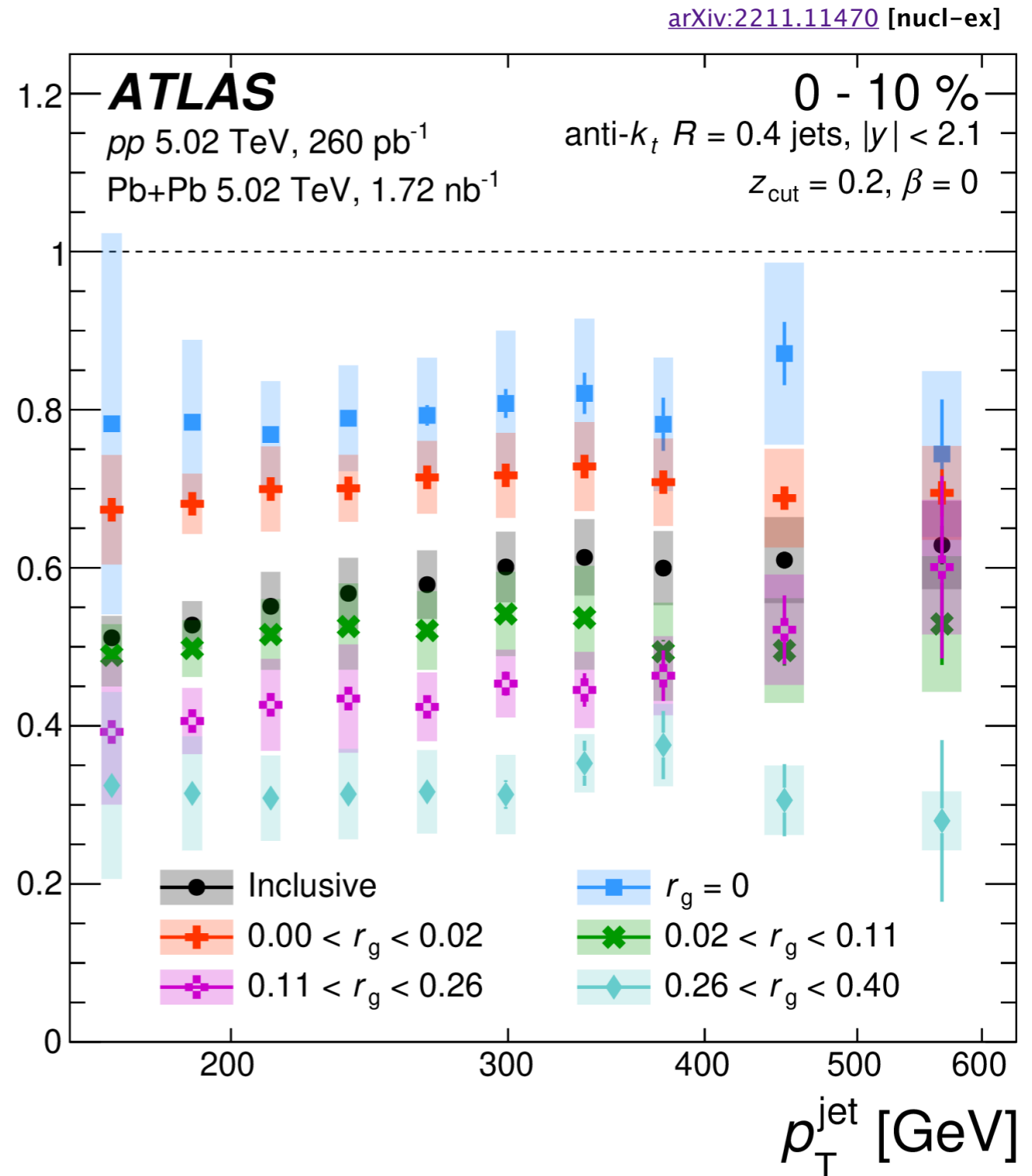
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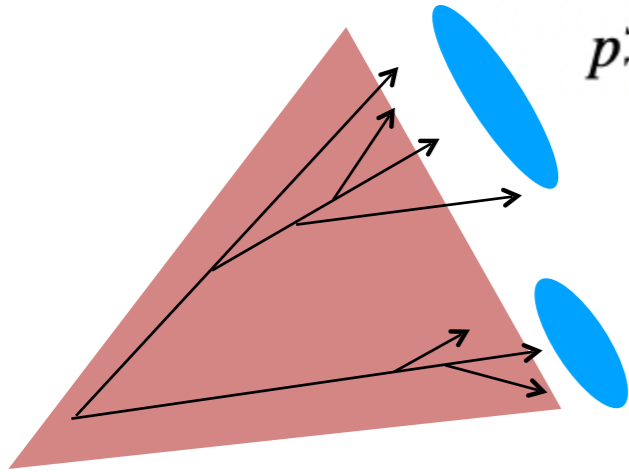
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$R_{AA}$



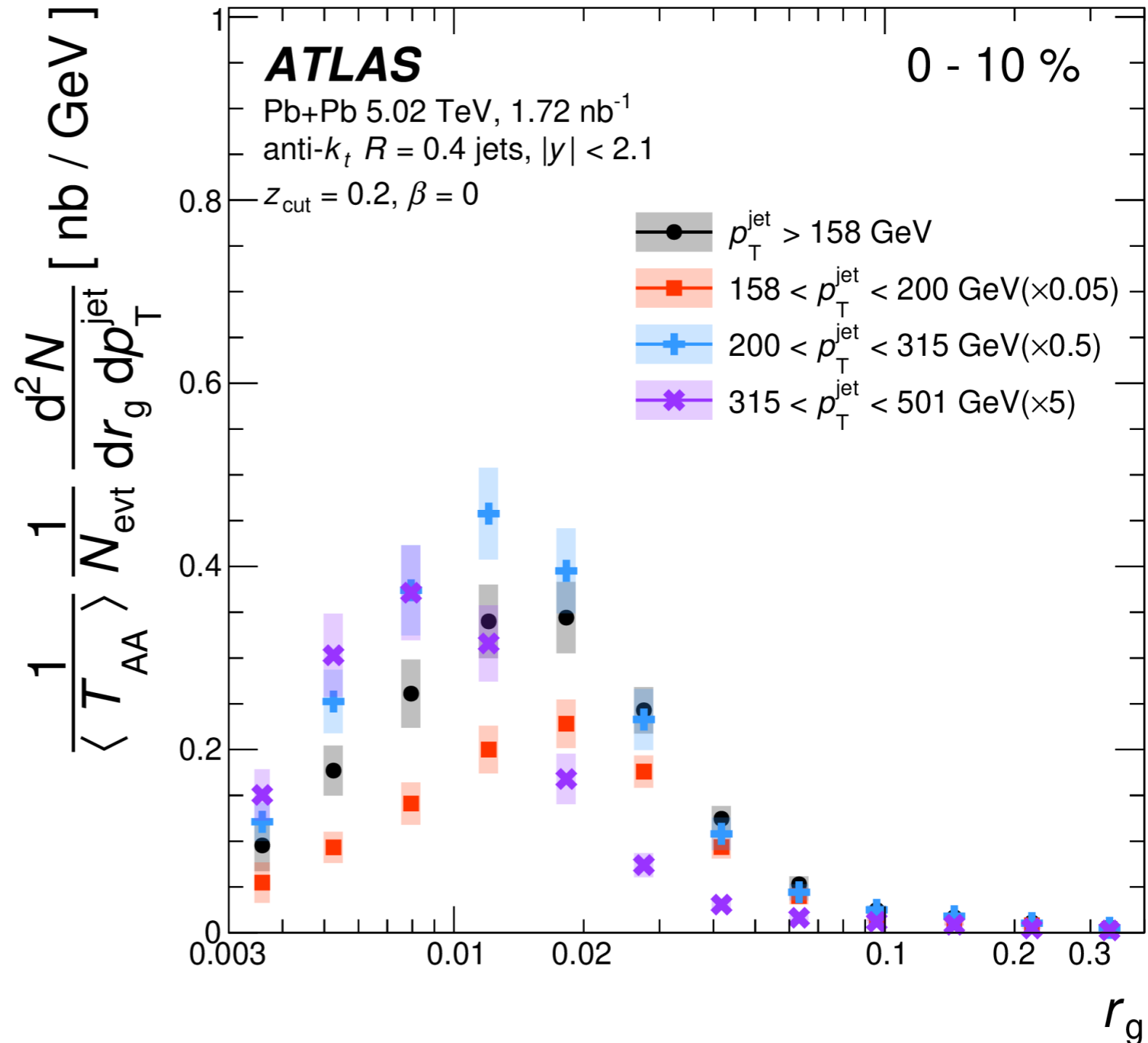
# Energy loss dependence on the jet substructure

$$\frac{\min(p_T^{sj1}, p_T^{sj2})}{p_T^{sj1} + p_T^{sj2}} > z_{\text{cut}} \left( \frac{\Delta R_{12}}{R} \right)^\beta$$



arXiv:2211.11470 [nucl-ex]

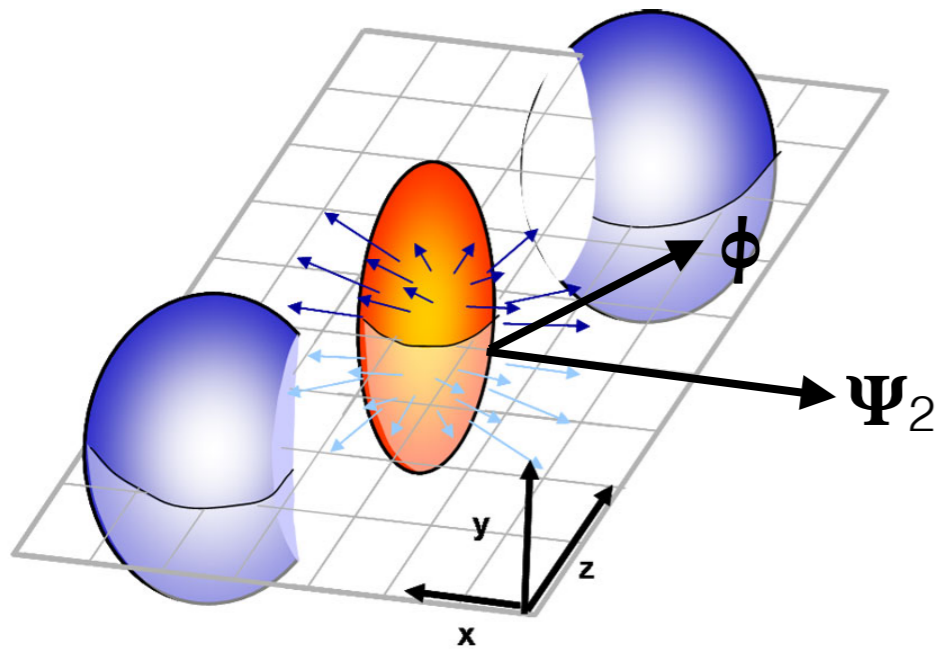
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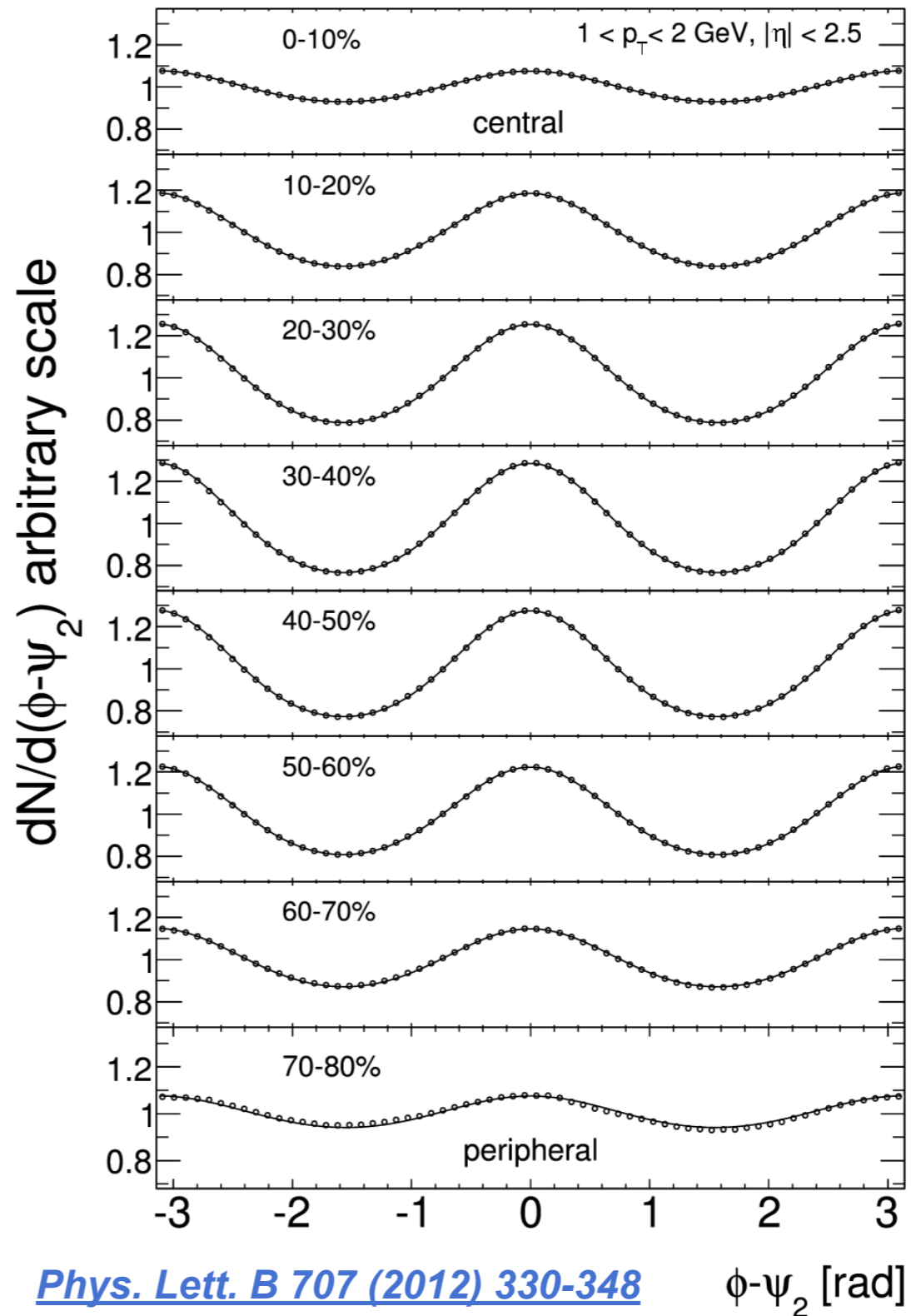
# Collectivity in heavy-ion collisions

Geometry driven  $v_2$



$$\frac{dN}{d\phi} \propto 1 + 2v_N \cos[N(\phi - \Psi_N)]$$

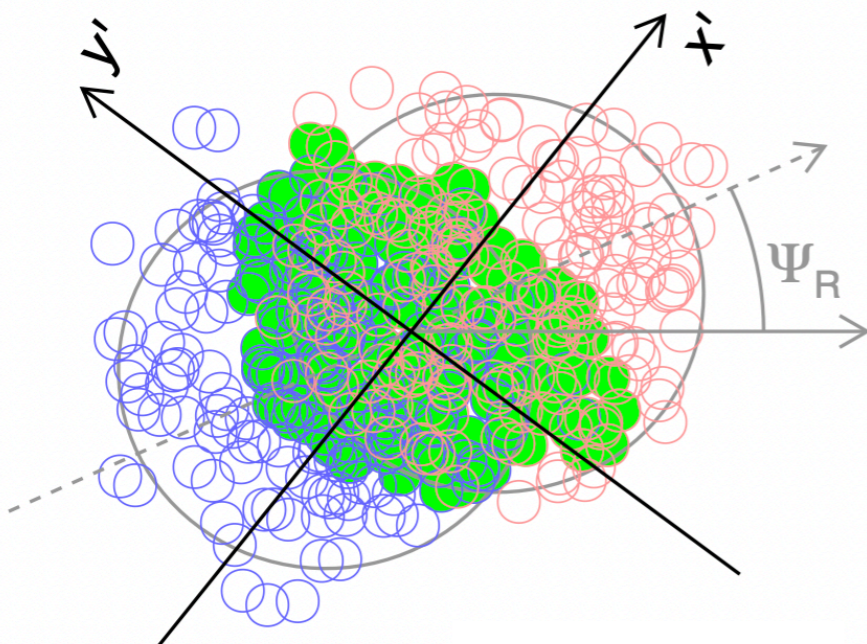
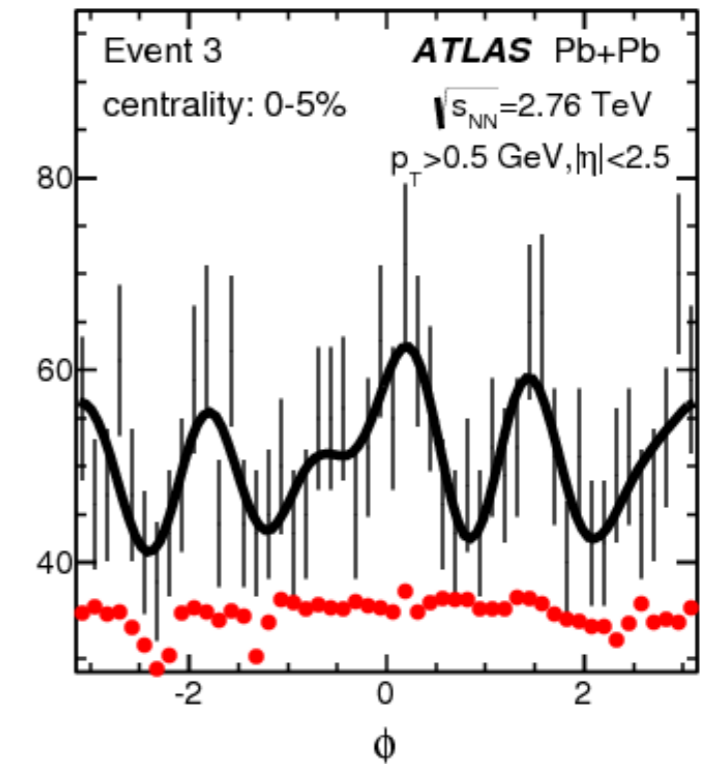
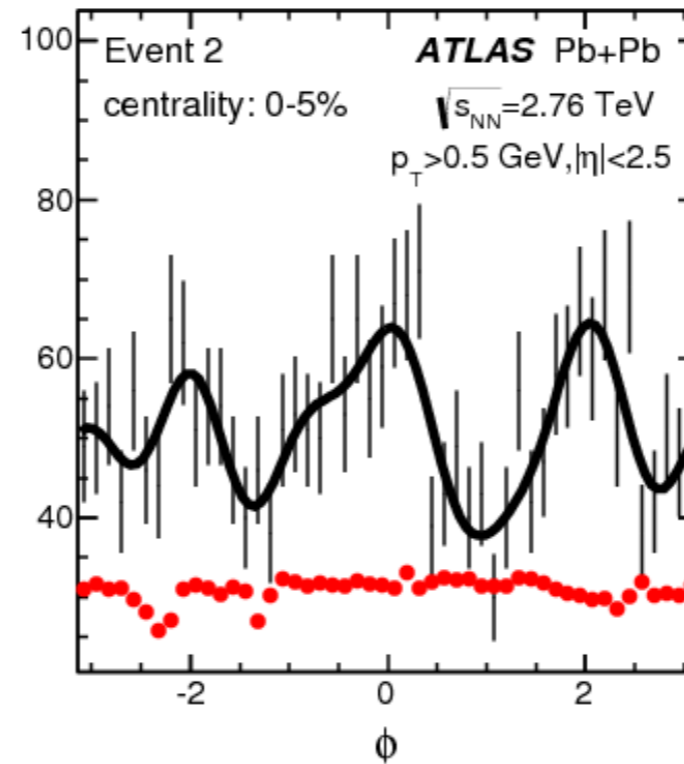
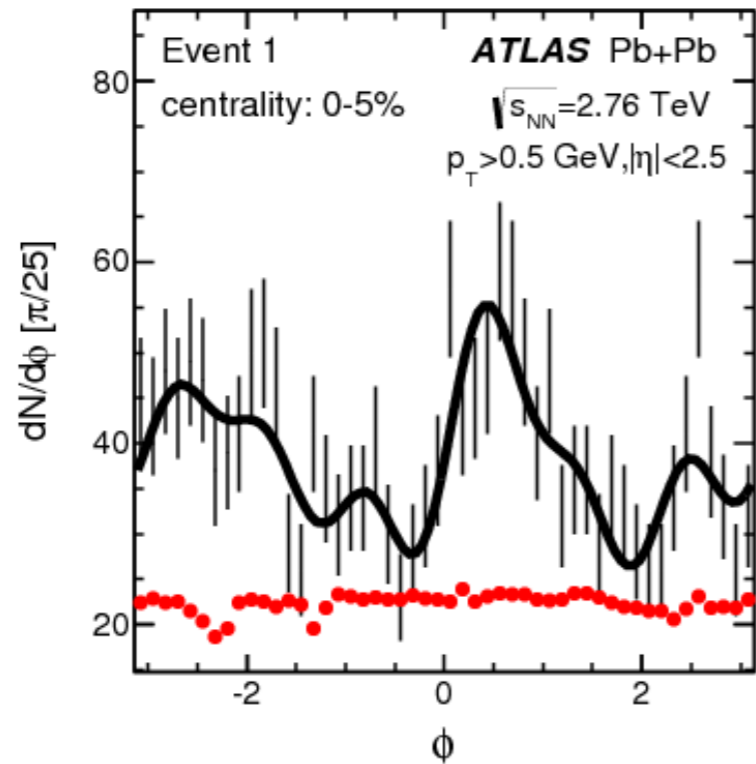
ATLAS Pb+Pb  $\sqrt{s_{NN}} = 2.76$  TeV  $L_{int} = 7 \mu\text{b}^{-1}$



# Collectivity in heavy-ion collisions

$$\frac{dN}{d\phi} \propto 1 + 2v_N \cos[N(\phi - \Psi_N)]$$

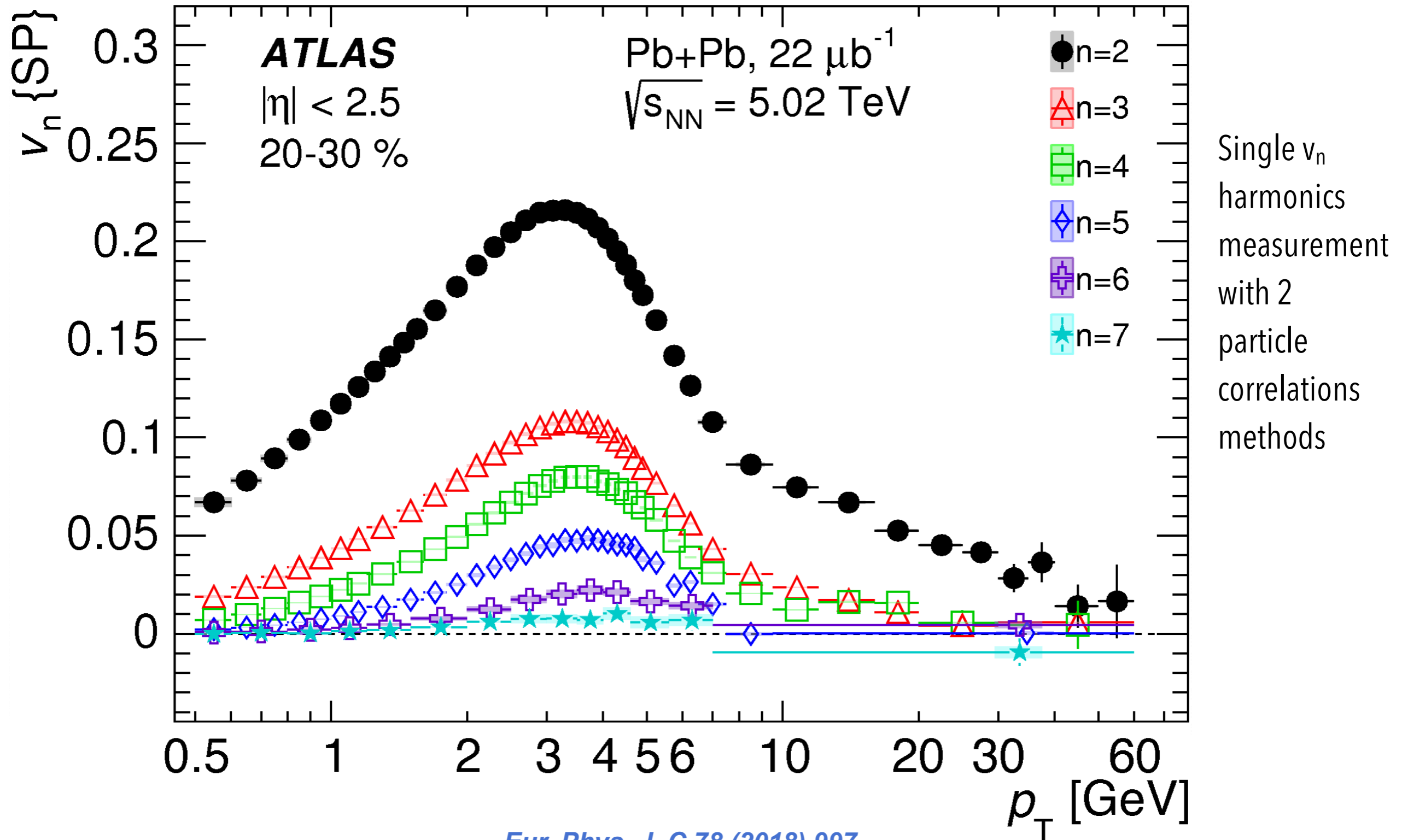
[JHEP 11 \(2013\) 183](#)



Fluctuation driven higher order harmonics  $v_n$  ( $n > 2$ )

# Collectivity in heavy-ion collisions

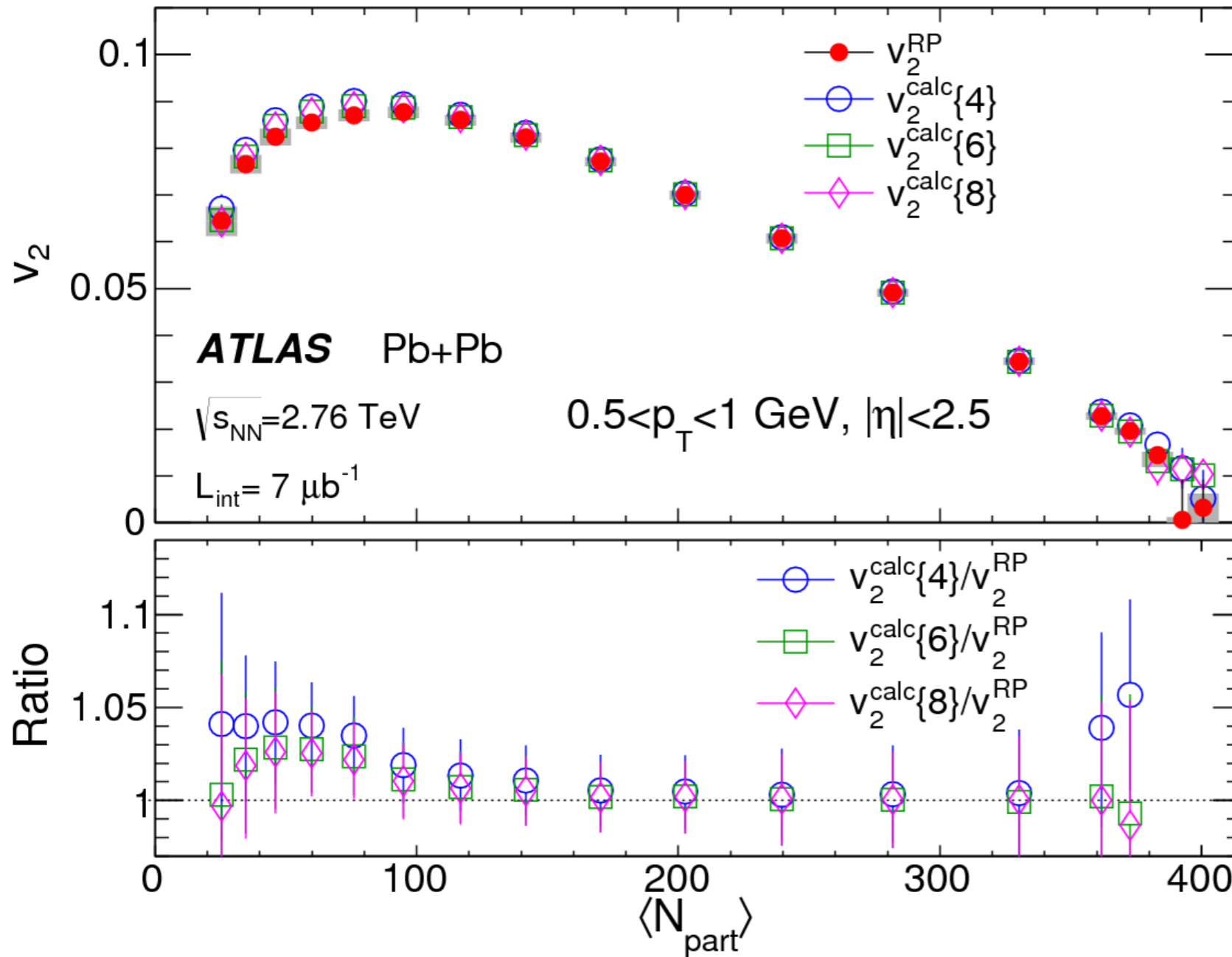
$$\frac{dN}{d\phi} \propto 1 + 2v_N \cos[N(\phi - \Psi_N)]$$



[Eur. Phys. J. C 78 \(2018\) 997](#)

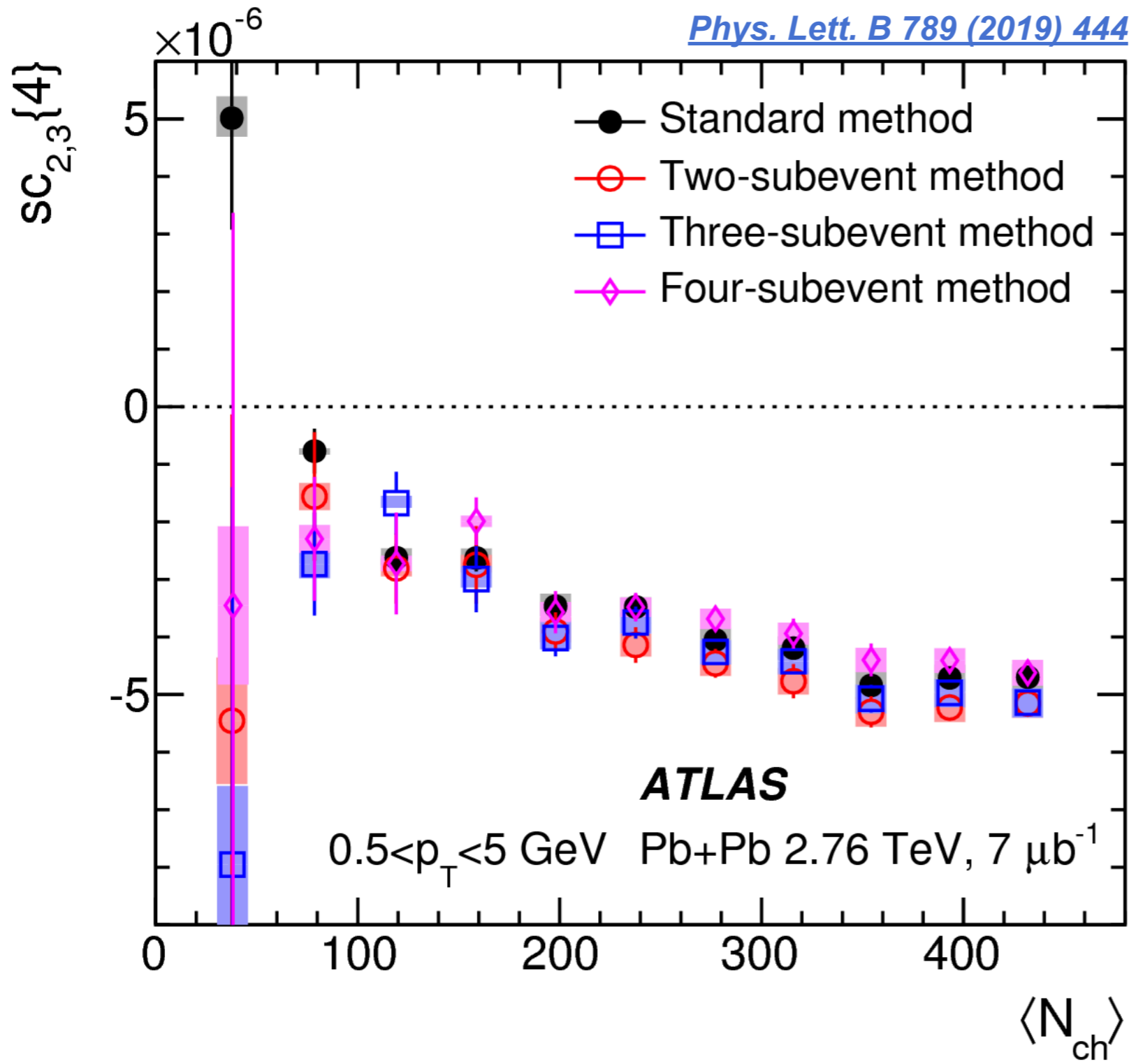
# Evolution of collectivity measurements

*JHEP 11 (2013) 183*



Multi particle cumulants

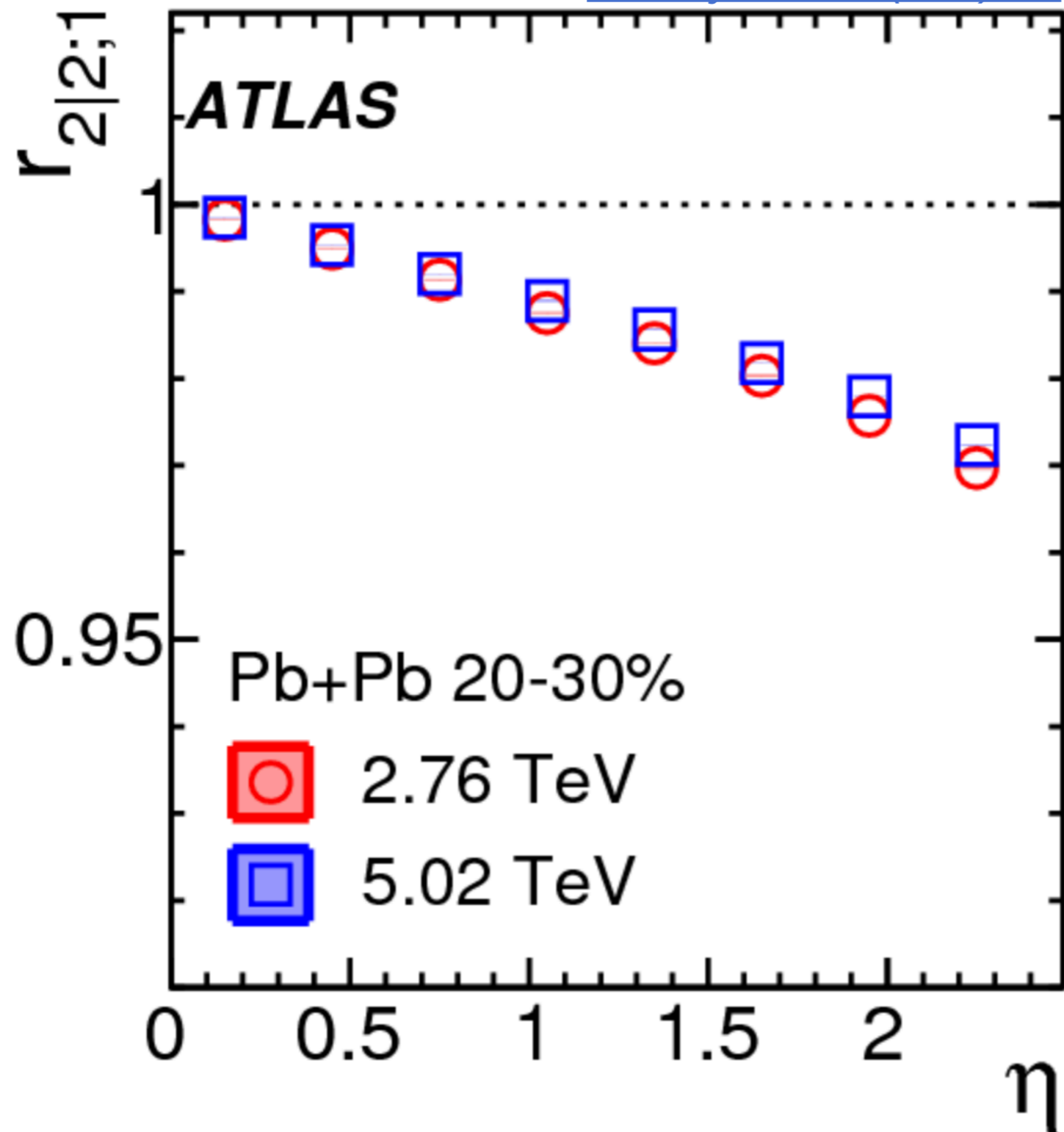
# Evolution of collectivity measurements



Correlation between harmonics of different order with symmetric cumulants (here anti-correlation between  $v_2$  and  $v_3$ )

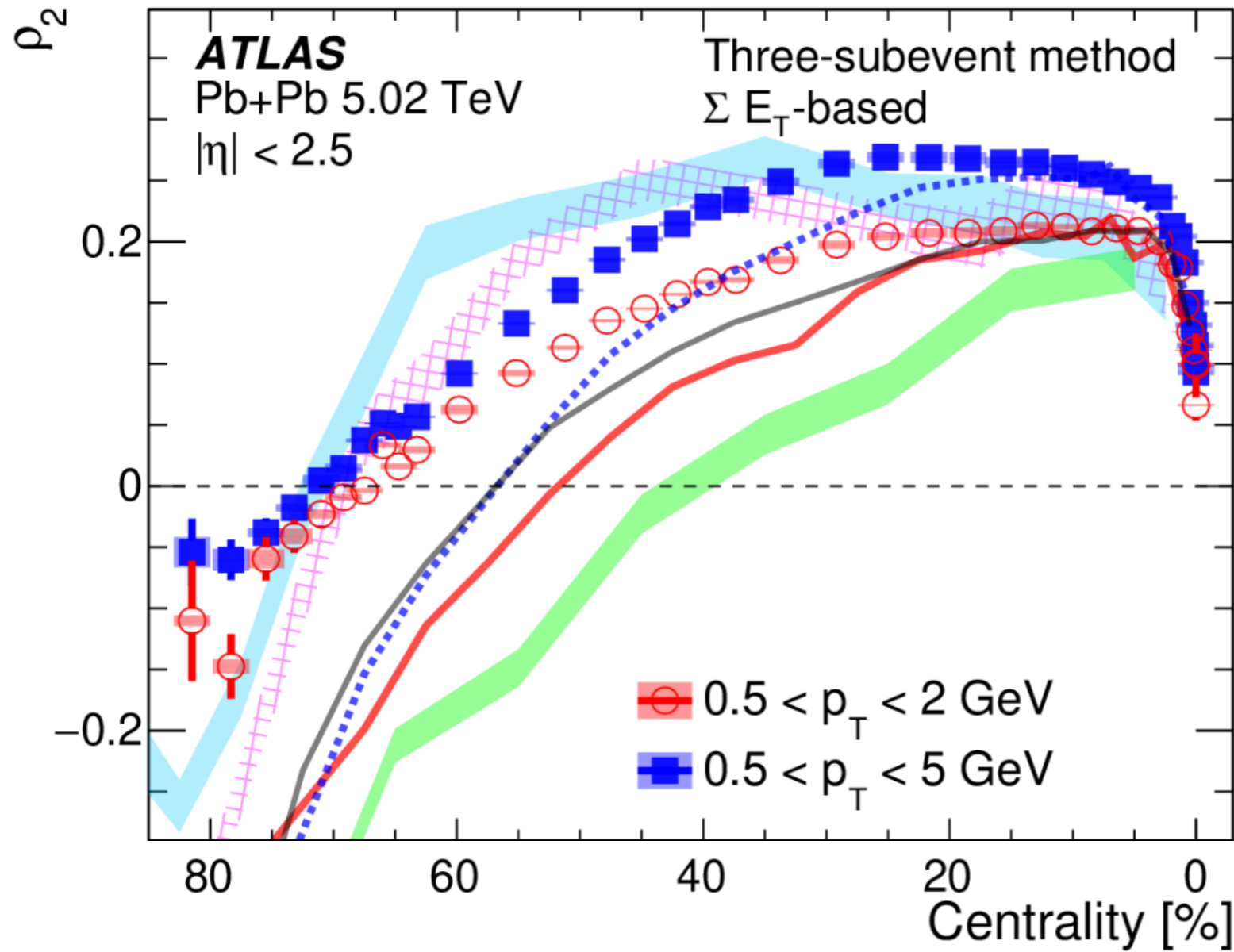
# Evolution of collectivity measurements

*Eur. Phys. J. C 76 (2018) 142*



# Evolution of collectivity measurements

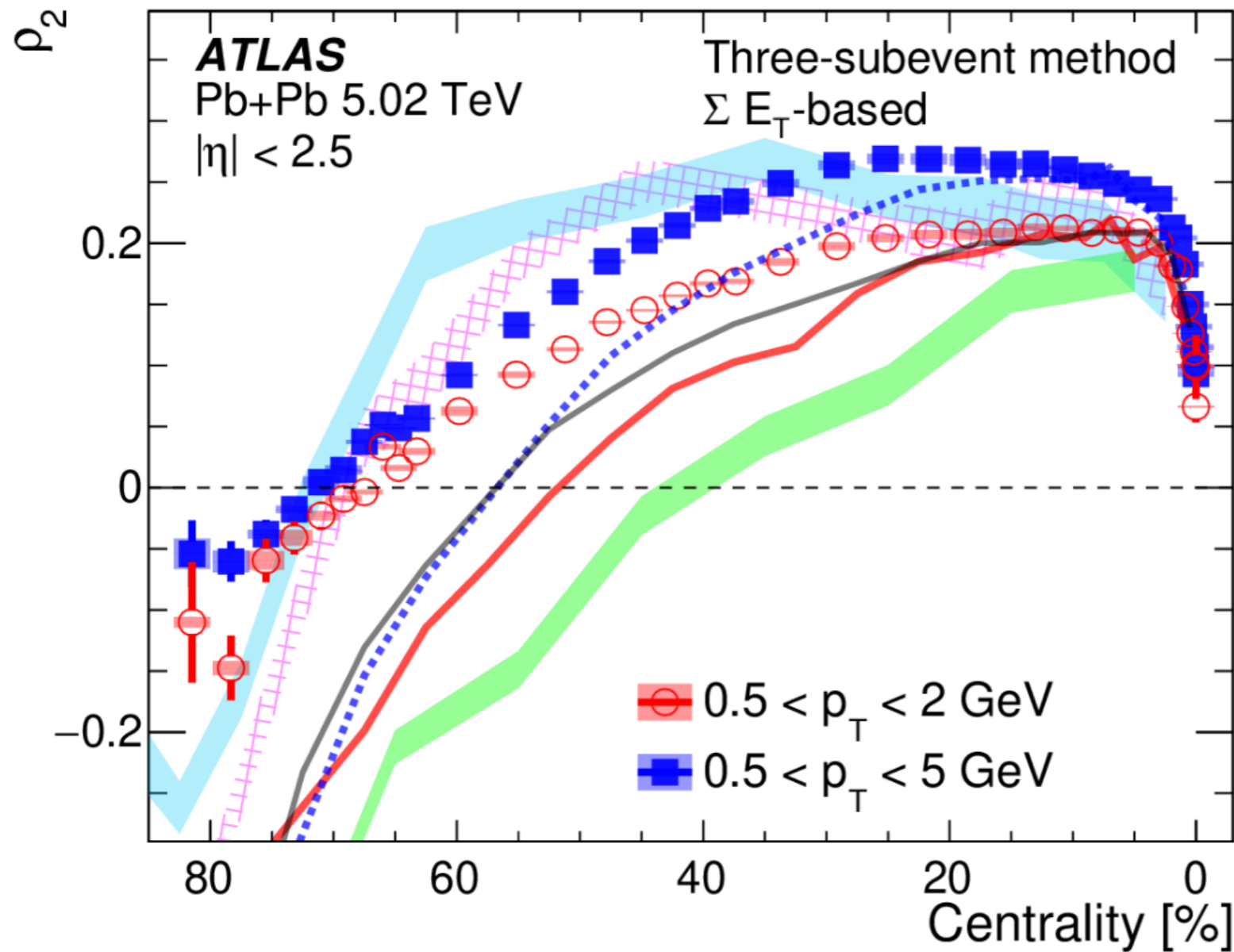
arXiv:2205.00039 [nucl-ex]



Correlation between  $v_n$  harmonics and mean  $p_T$  in the event (here for  $v_2$ )

# Evolution of collectivity measurements

arXiv:2205.00039 [nucl-ex]

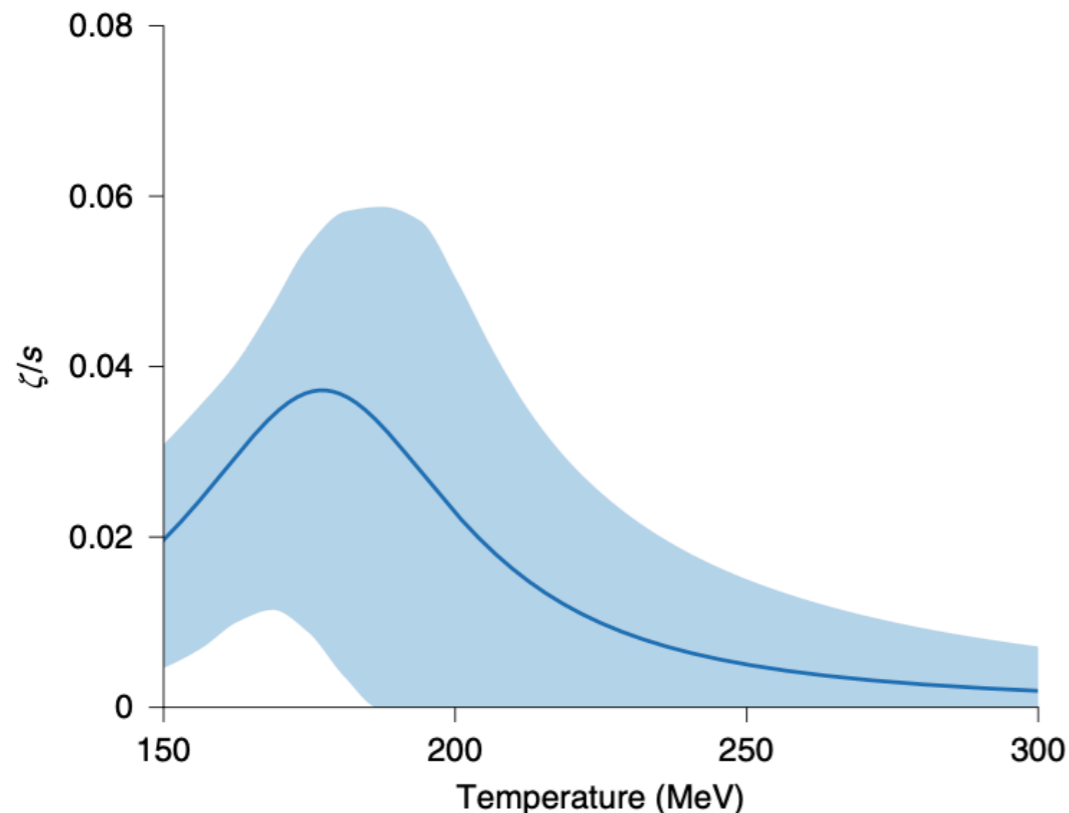
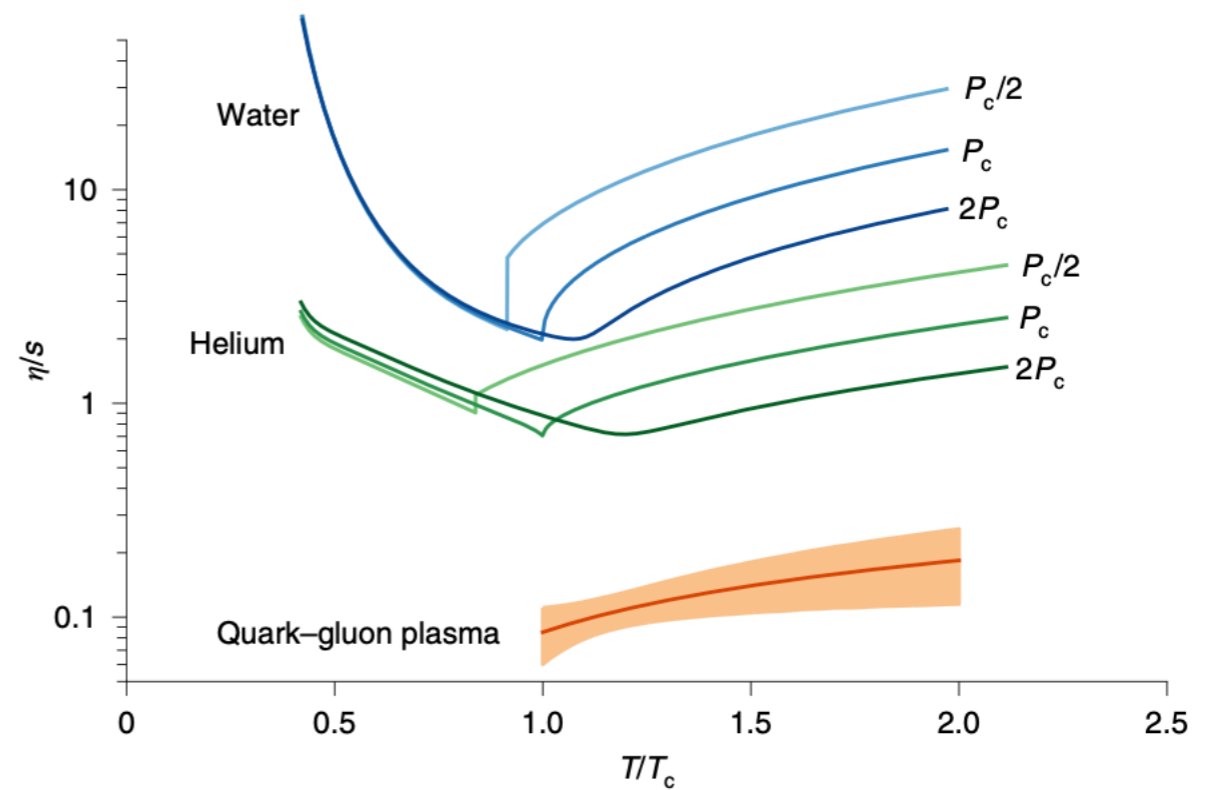
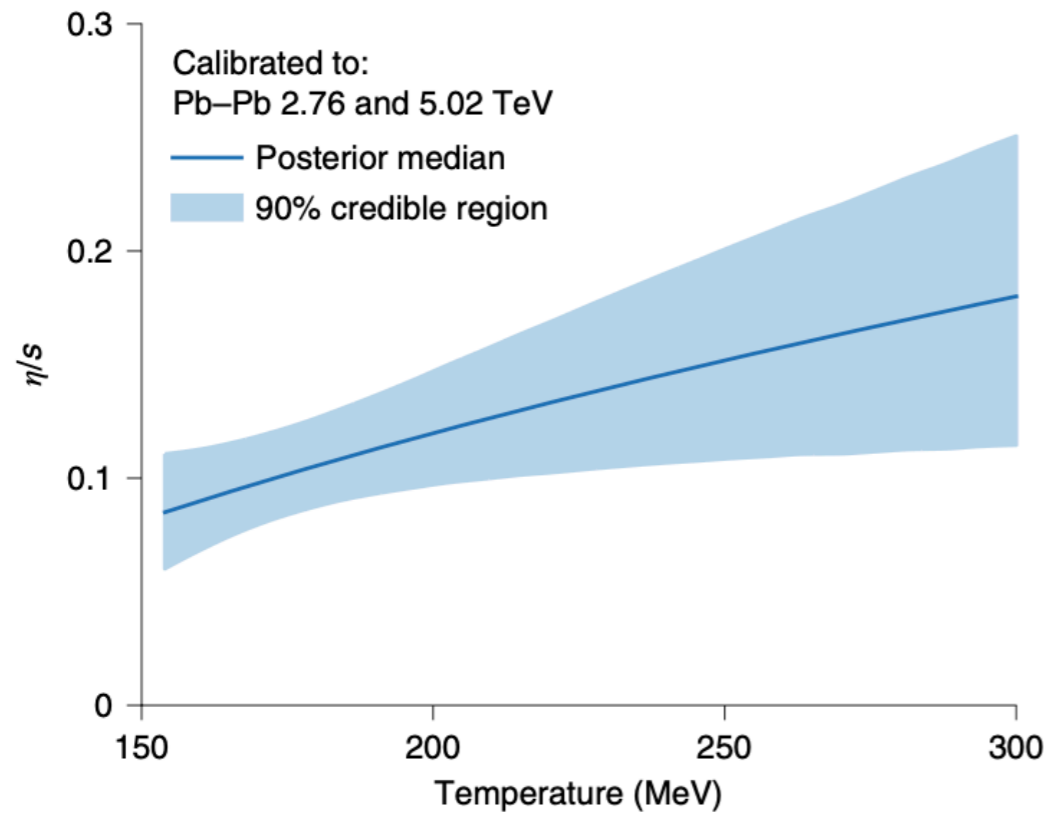


Correlation between  $v_n$  harmonics and mean  $p_T$  in the event (here for  $v_2$ )

**Disclaimer** - by no means this is not complete list, it is only meant to show increasing complexity



# Constraining QGP parameters



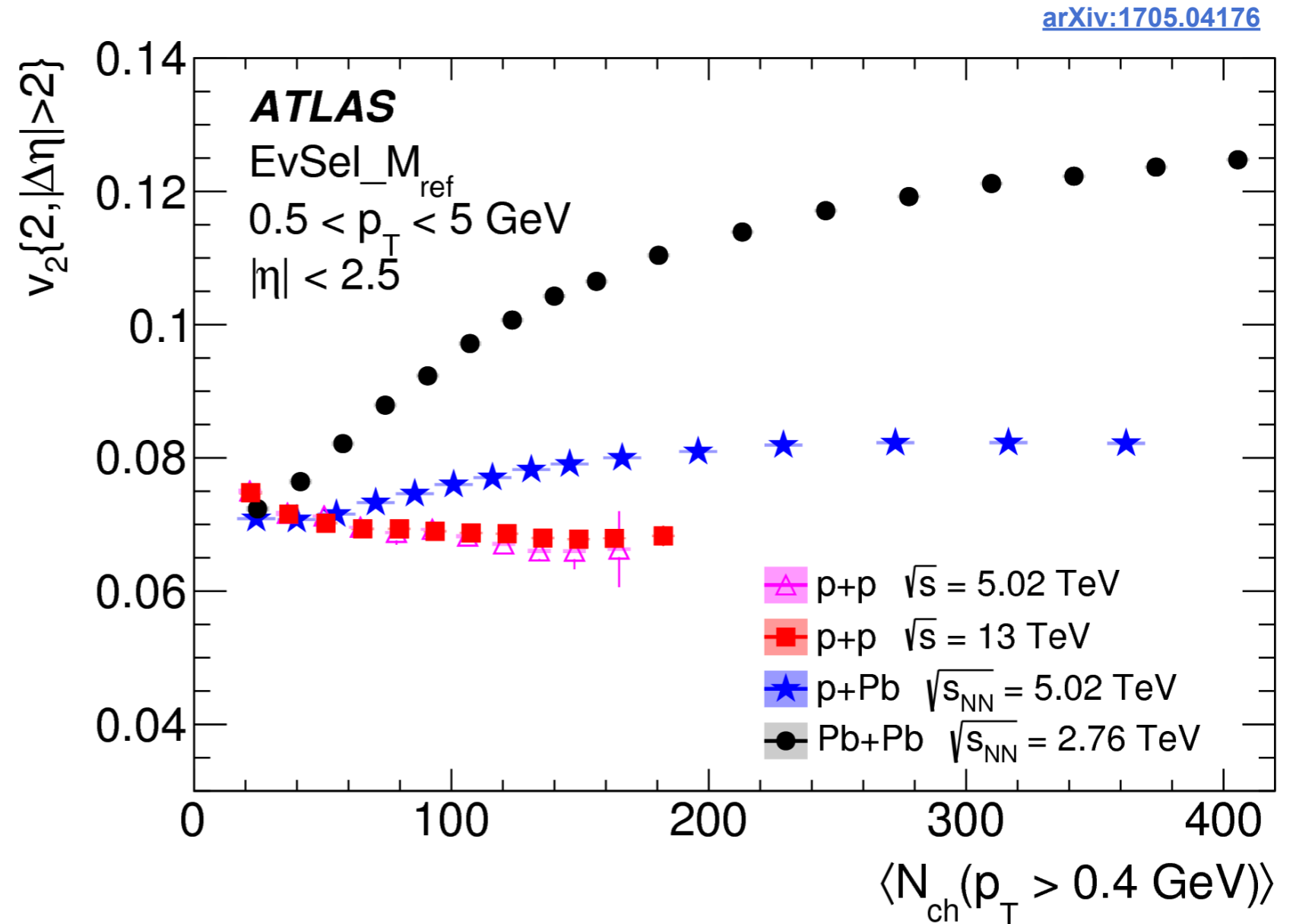
LHC data are used to constrain shear ( $\eta$ ) and bulk ( $\zeta$ ) viscosity of the QGP ....

.... a lot of assumptions on the form of initial state, exact hydrodynamic model, transition to hadrons ...

More data, different collisions systems helps

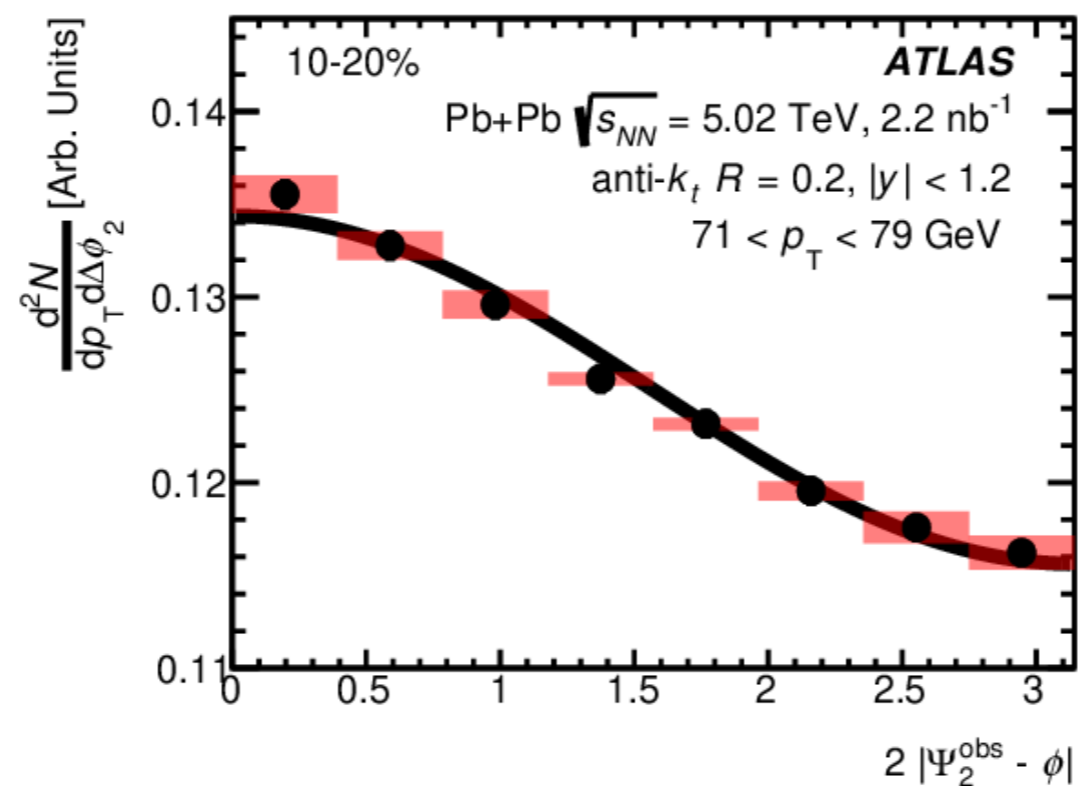
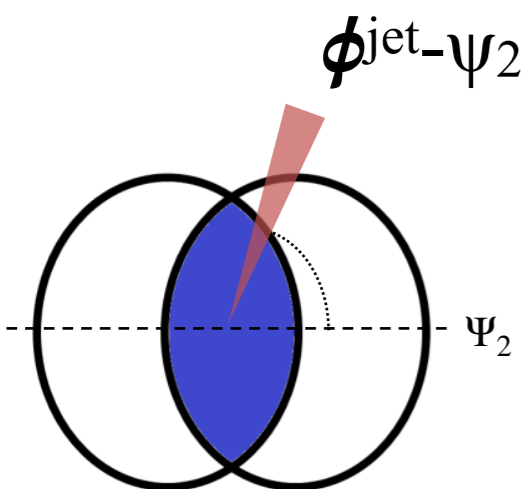
# Collectivity in small systems

One of the most striking observations from LHC data is a presence of collective effects in p+Pb, pp and even smaller collision systems.

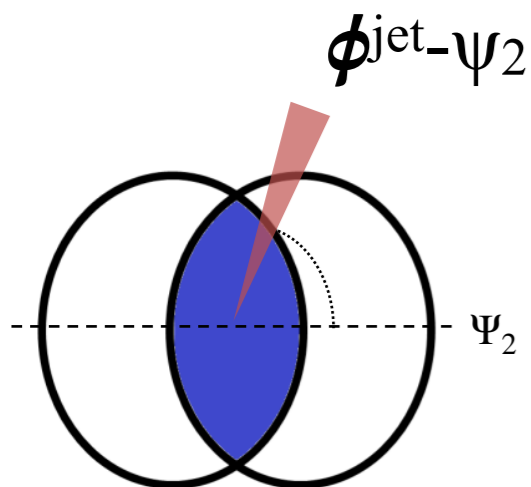


- It is still unclear what is the smallest droplet of QGP that behaves hydrodynamically
- Are there any other QGP like effects present in small systems ?

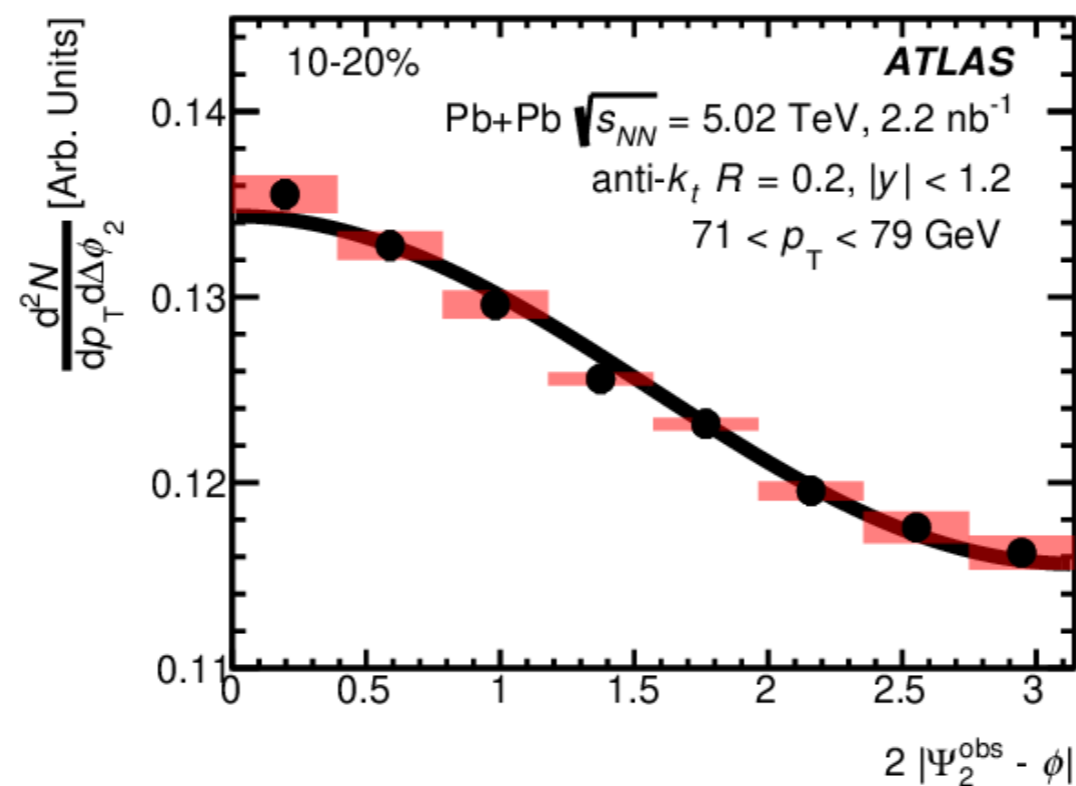
[Phys. Rev. C 105 \(2022\) 064903](#)



measurement of the jets  
 wrt. collision geometry  
 this varies the amount of  
 QGP that the jet sees

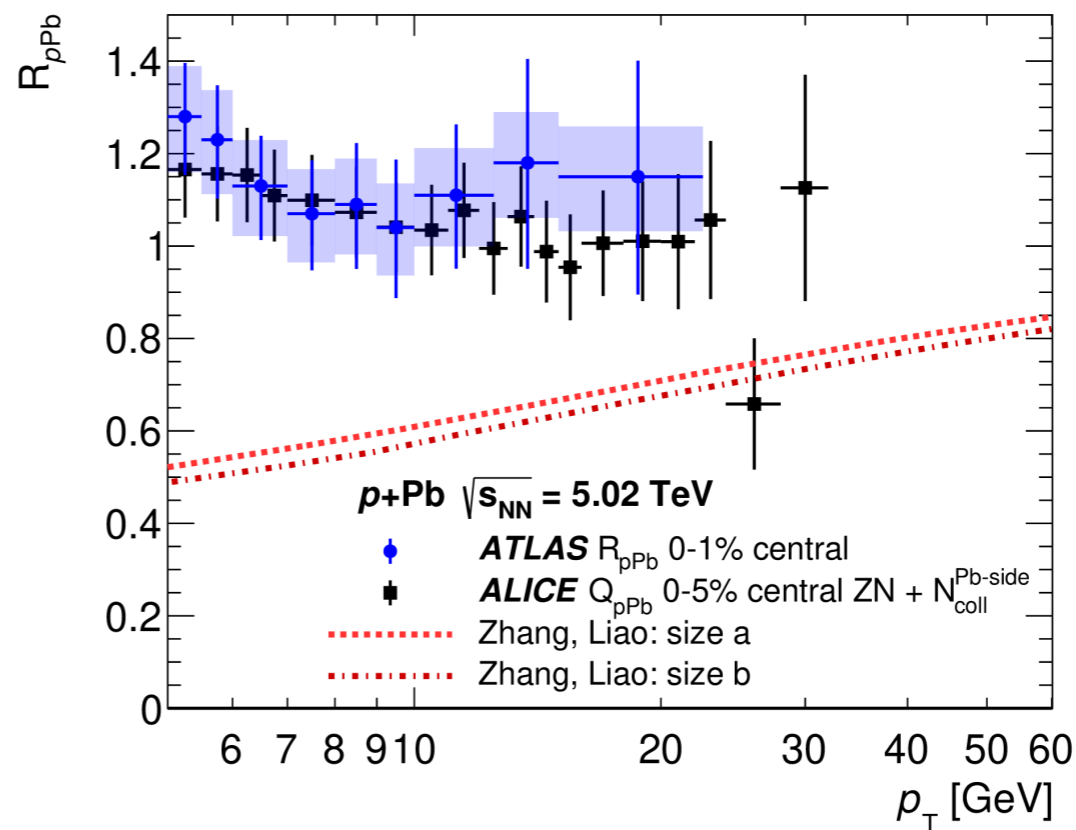
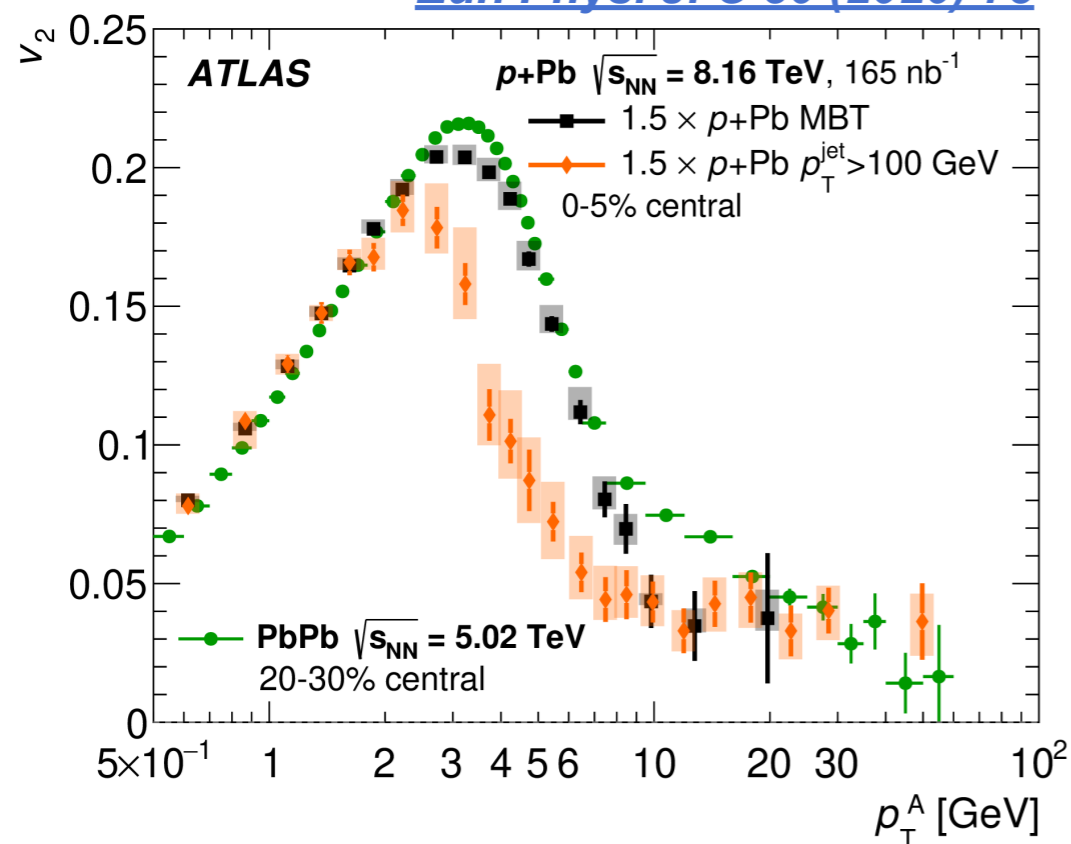


*Phys. Rev. C 105 (2022) 064903*



measurement of the jets  
 wrt. collision geometry  
 this varies the amount of  
 QGP that the jet sees

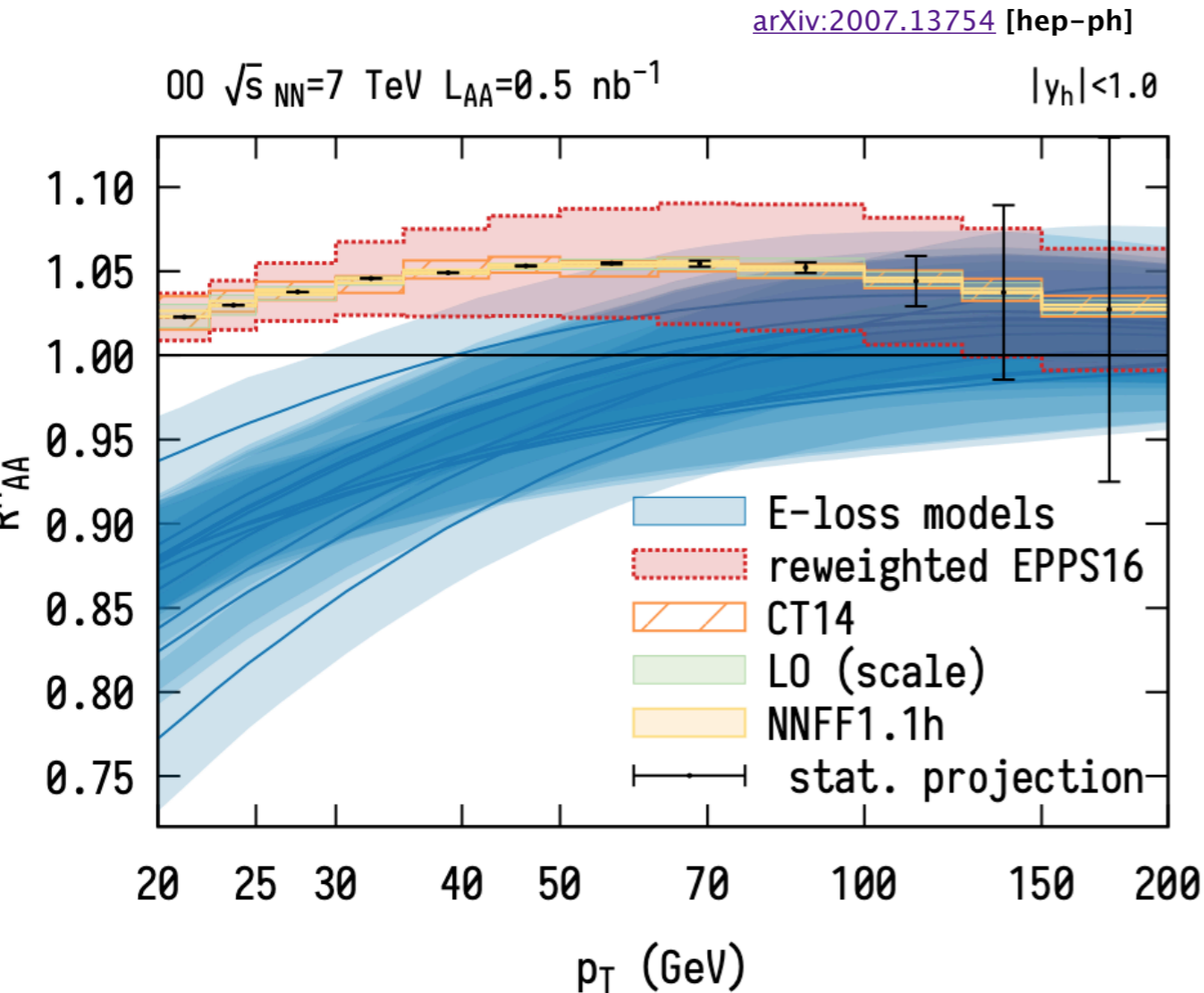
*Eur. Phys. J. C 80 (2020) 73*



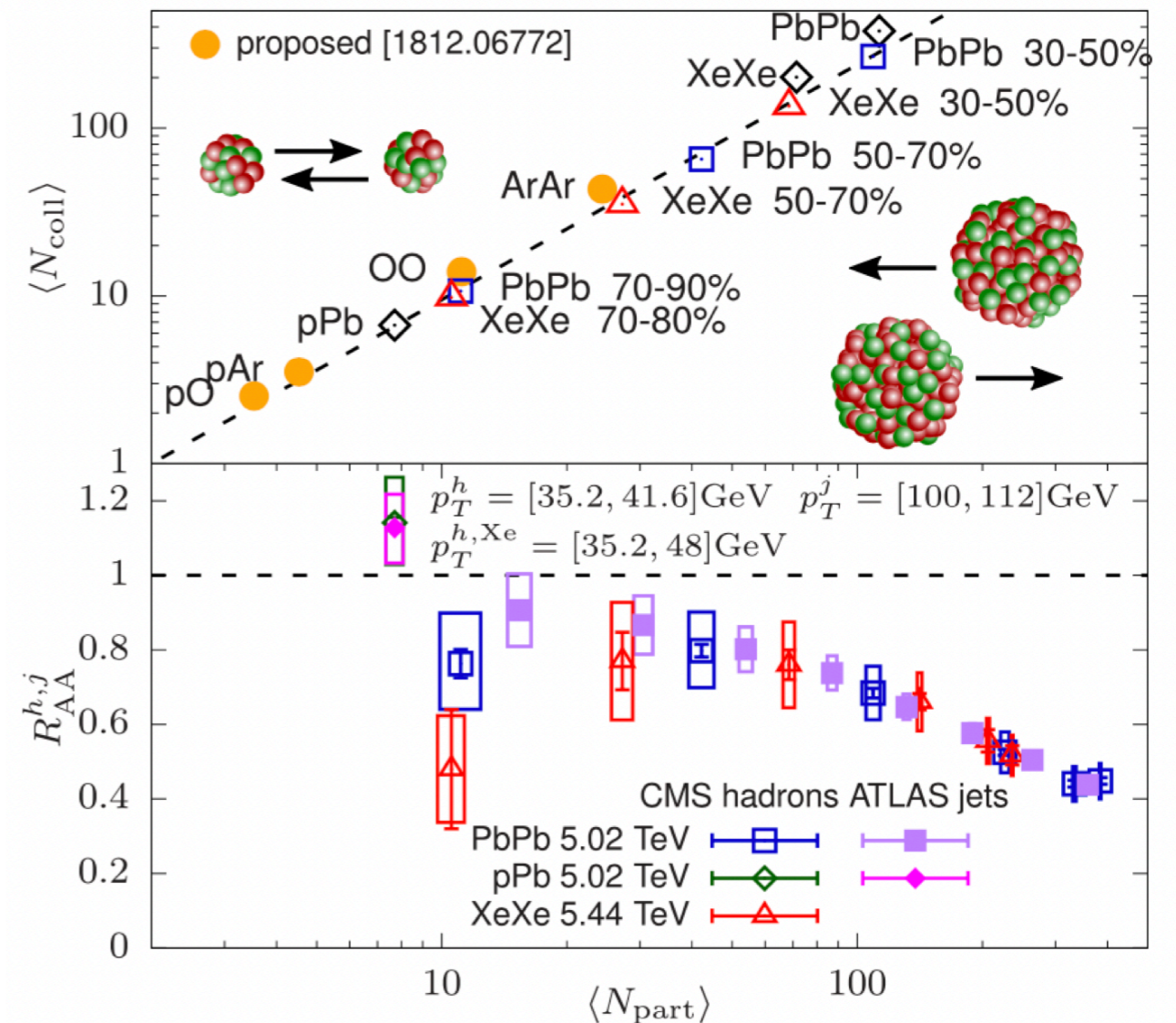
$p+Pb$ :  
 Non zero  $v_2$   
 No signs of  $R_{AA}$   
 modification

# Jet quenching in light ions collisions

Measurements of p+A type of collisions at LHC and RHIC left us with unresolved problem: **how to connect soft QGP with lack of modification in the hard sector?**



OO is a symmetric system about the size of p+Pb

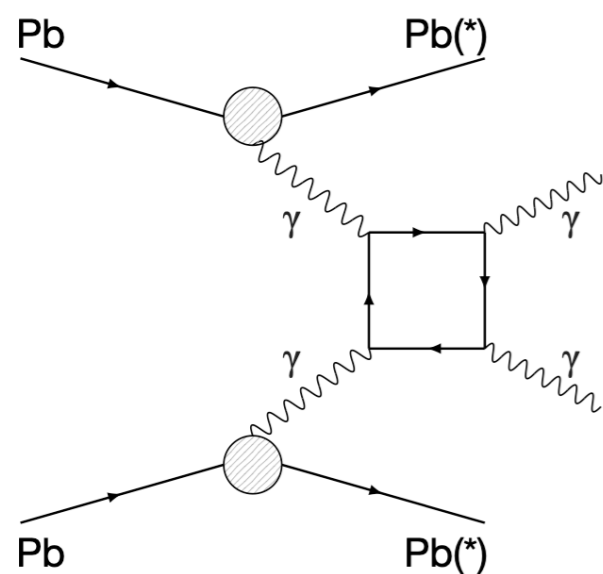


# Photon induced process

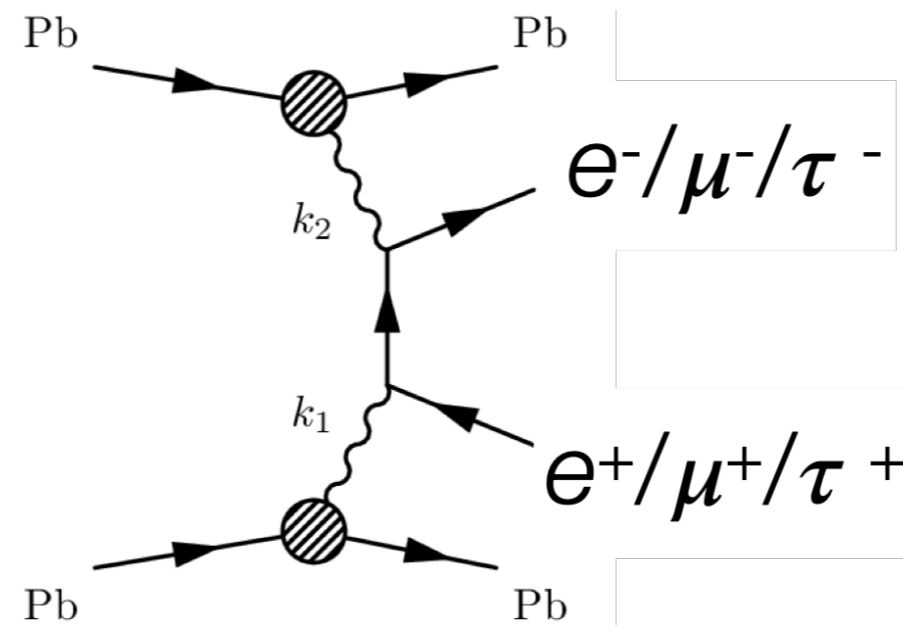
Boosted nuclei are intense source of photons

- Coherent photon flux with  **$E_{\text{max}} \sim 80 \text{ GeV @ LHC}$**

Various types of interactions possible:



Light-by-light scattering



Di-lepton production

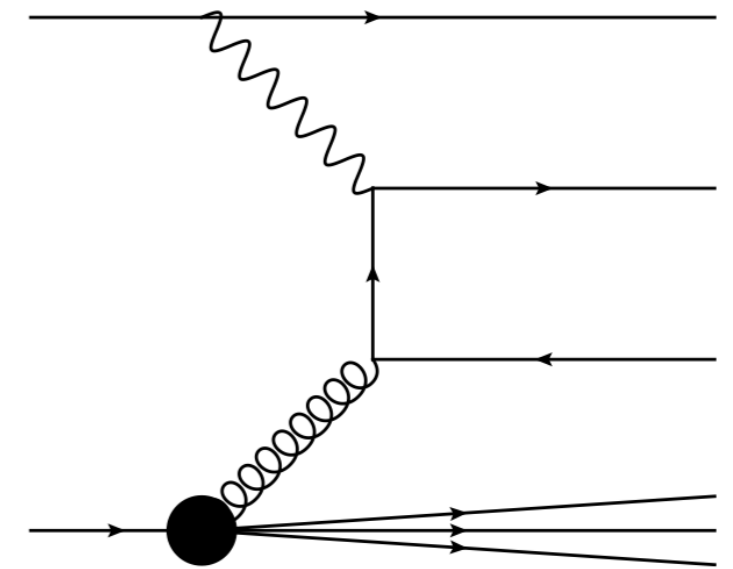
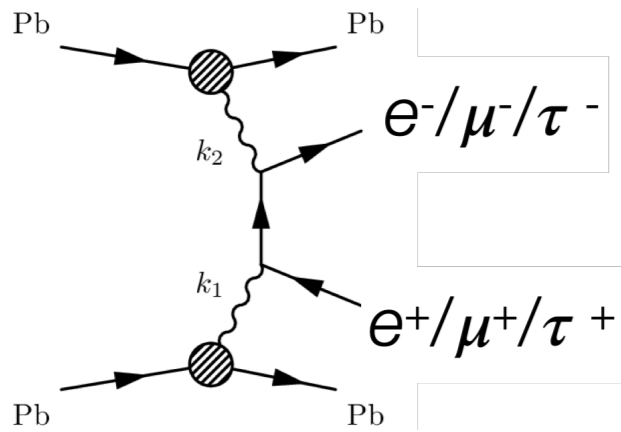


Photo nuclear

- Using ZDC to categorize events ( $0n, Xn$ )
- Rapidity gap helps to distinguish photo nuclear from inelastic Pb+Pb events

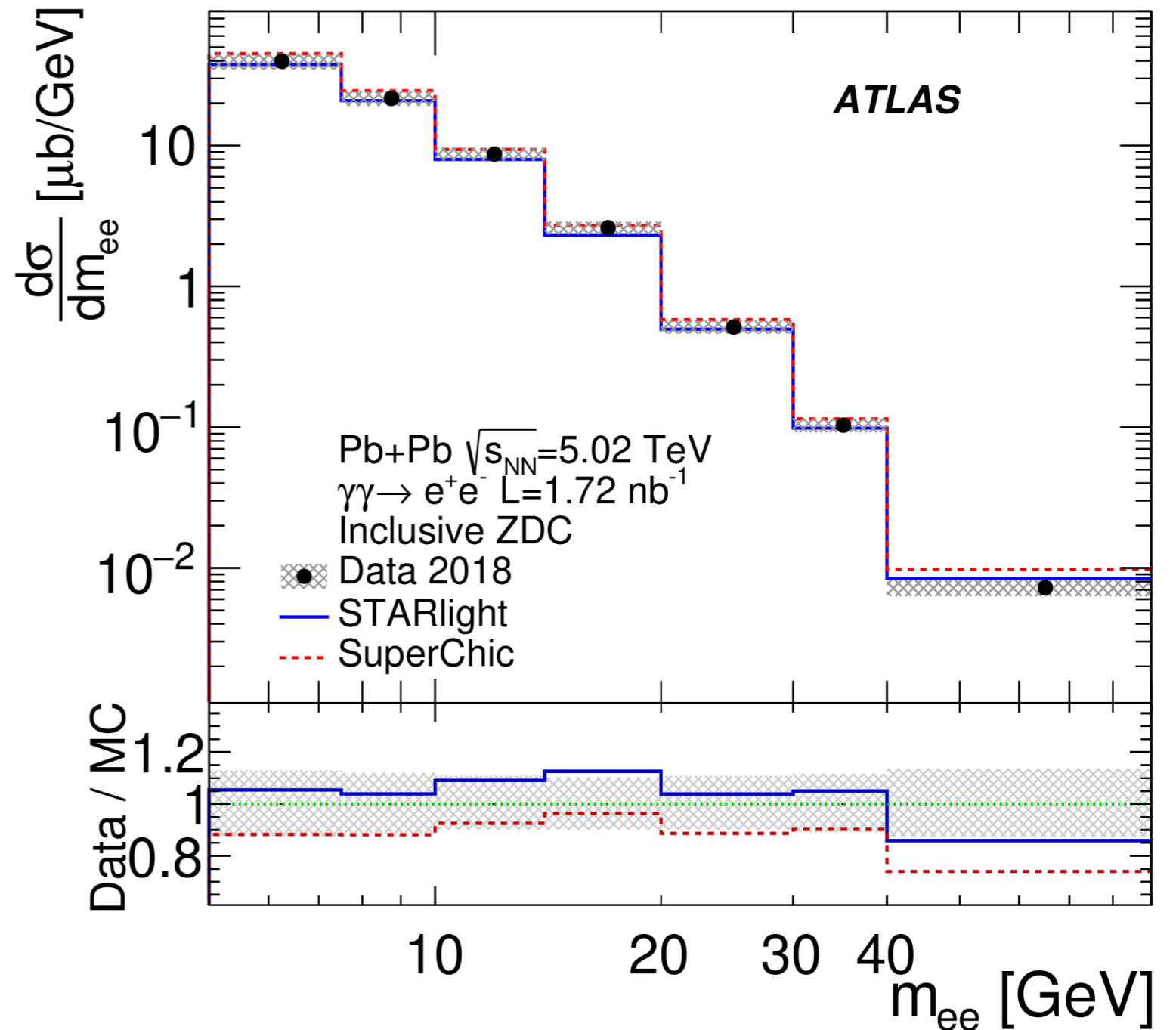
# Exclusive di-lepton production in UPC



[arXiv:2207.12781 \[nucl-ex\]](https://arxiv.org/abs/2207.12781)

Exclusive di-lepton production is one of the fundamental processes in  $\gamma+\gamma$  interaction  $\rightarrow$  benchmark process to understand other  $\gamma$  induced processes.

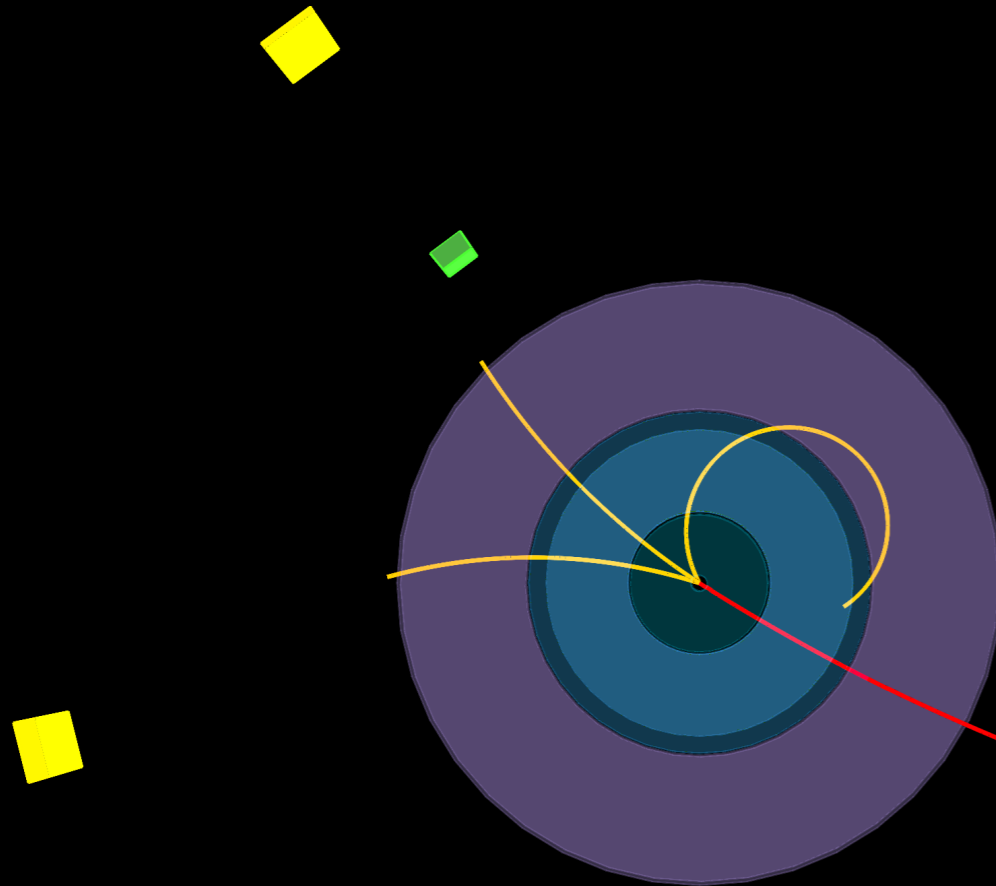
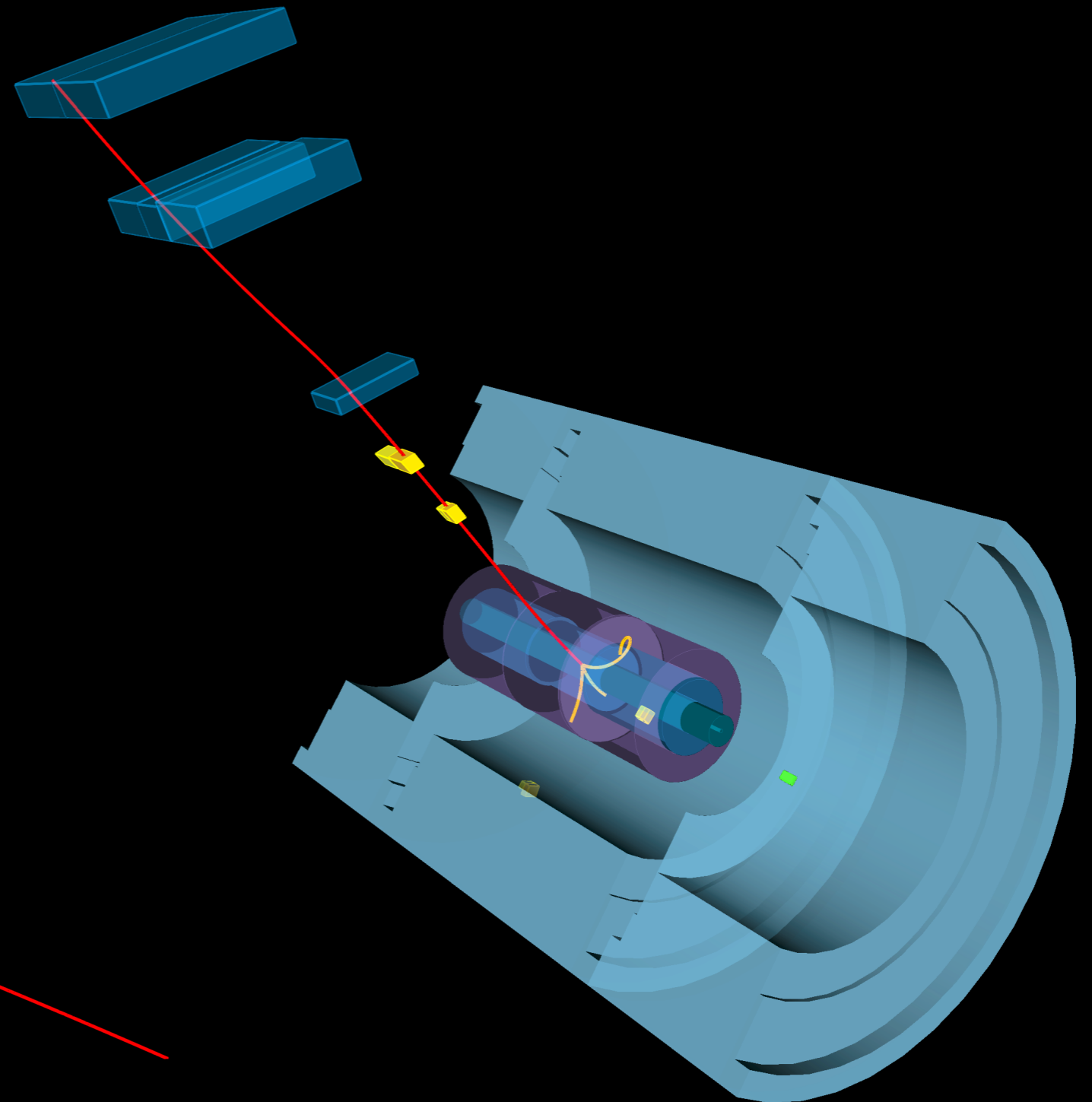
Cross section for di-lepton production depends on the photon fluxes from each nucleons.



# Exclusive di-lepton production in UPC

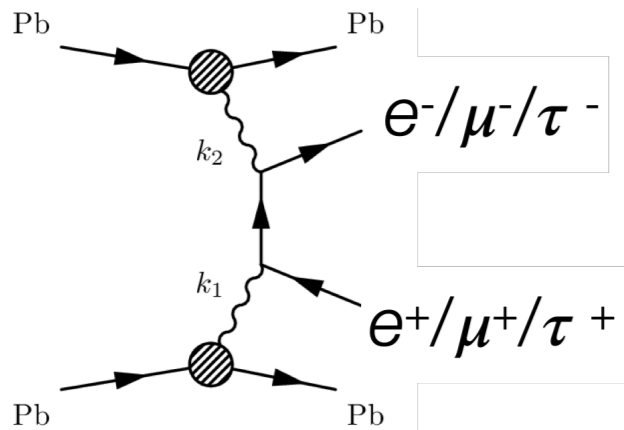


Run: 366268  
Event: 3305670439  
2018-11-18 16:09:33 CEST





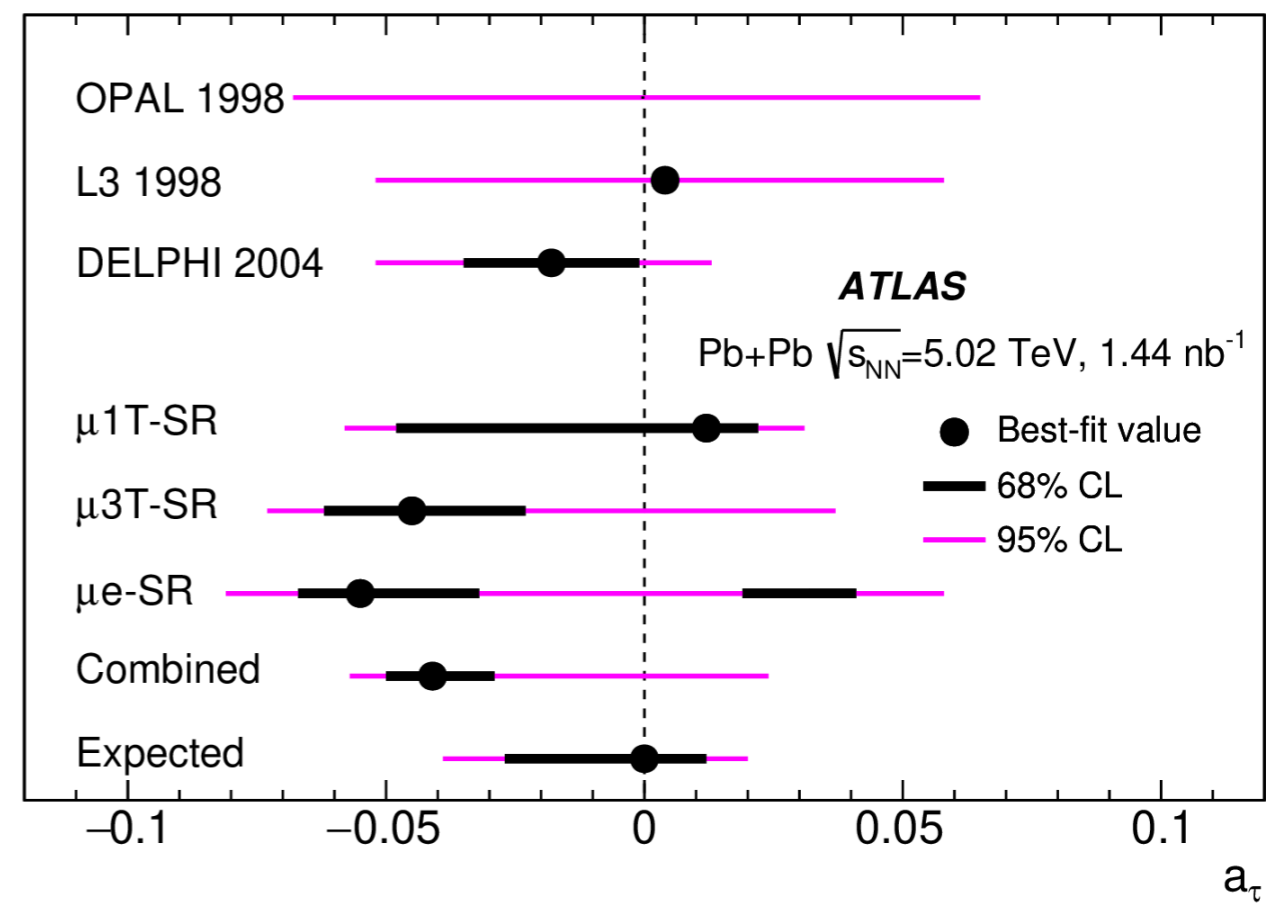
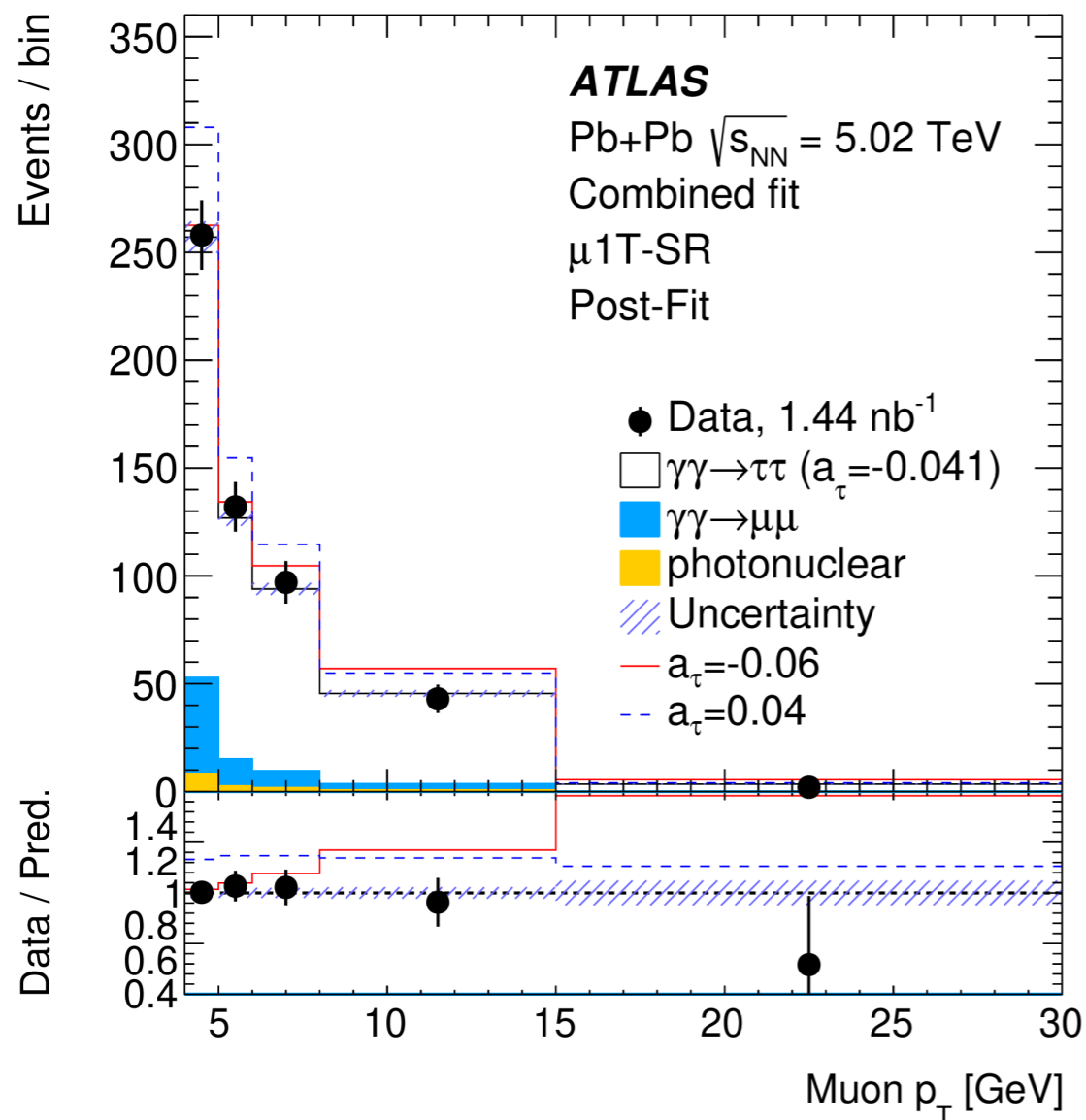
# Exclusive di-lepton production in UPC



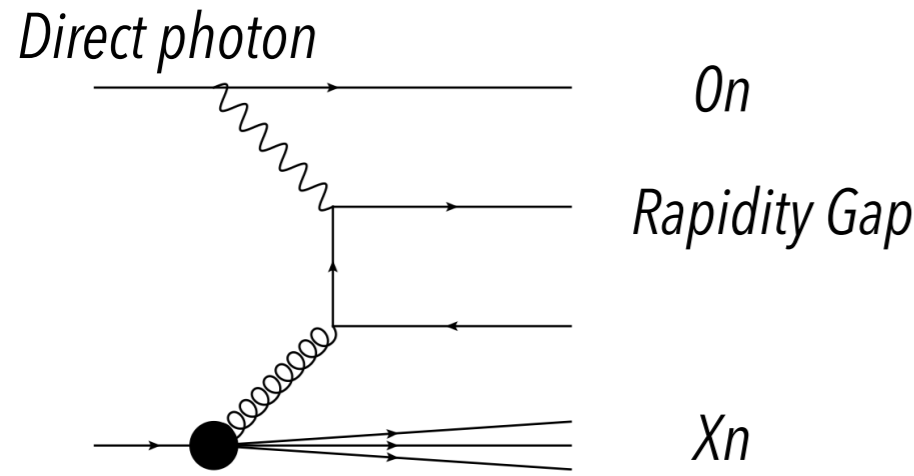
**Using  $\gamma\gamma \rightarrow \tau\tau$  process to set limits on anomalous magnetic moment of tau lepton ( $a_\tau$ )**

(Much less constrained than  $a_e$   $a_\mu$ )

[arXiv:2204.13478 \[hep-ex\]](https://arxiv.org/abs/2204.13478)



# Photo-nuclear production of di-jets



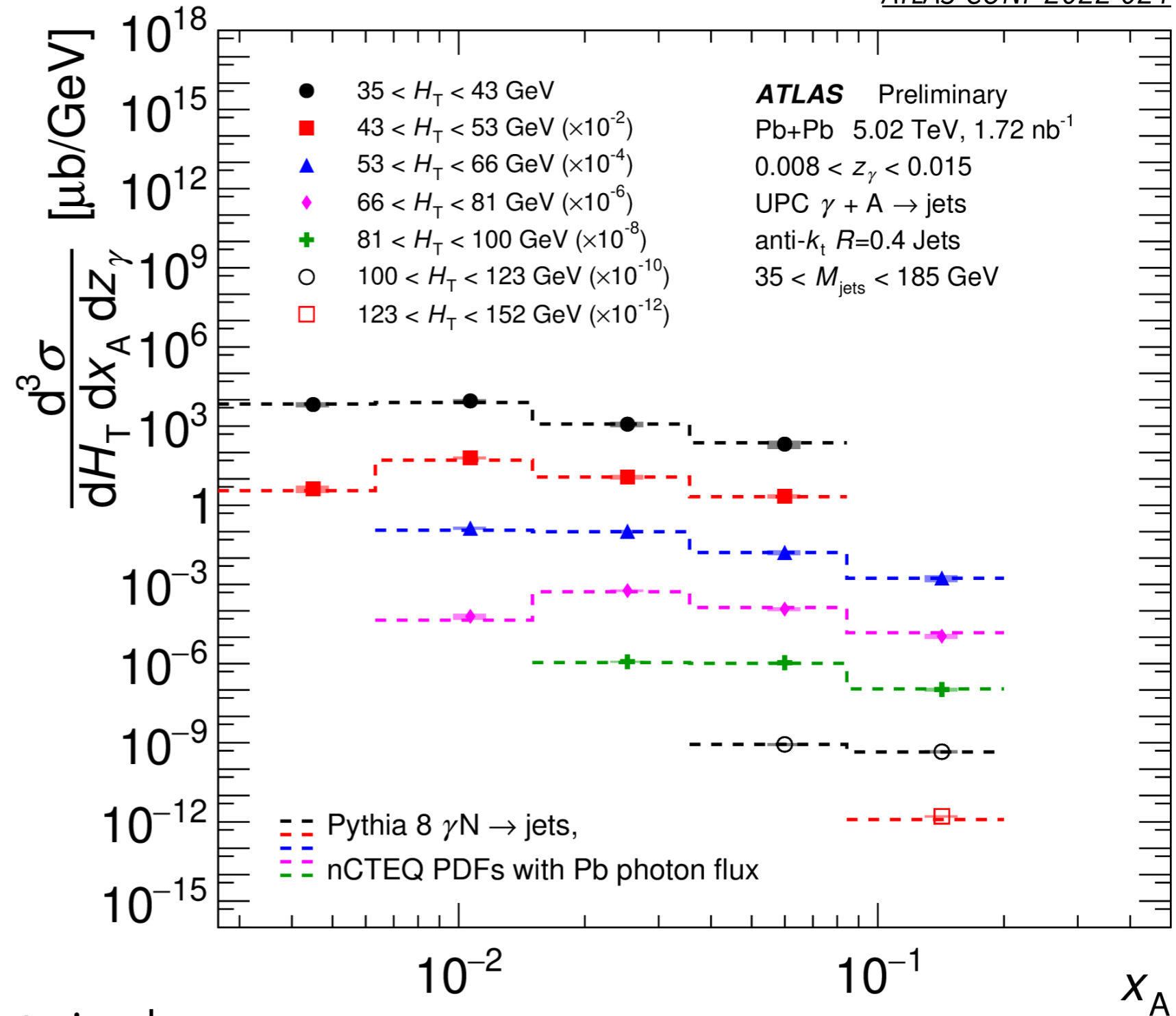
Di-jet kinematics corresponds to the hard scattering kinematics

$$H_T \equiv \sum_i p_{Ti}$$

$$z_\gamma \equiv \frac{M_{\text{jets}}}{\sqrt{s}} e^{+y_{\text{jets}}}$$

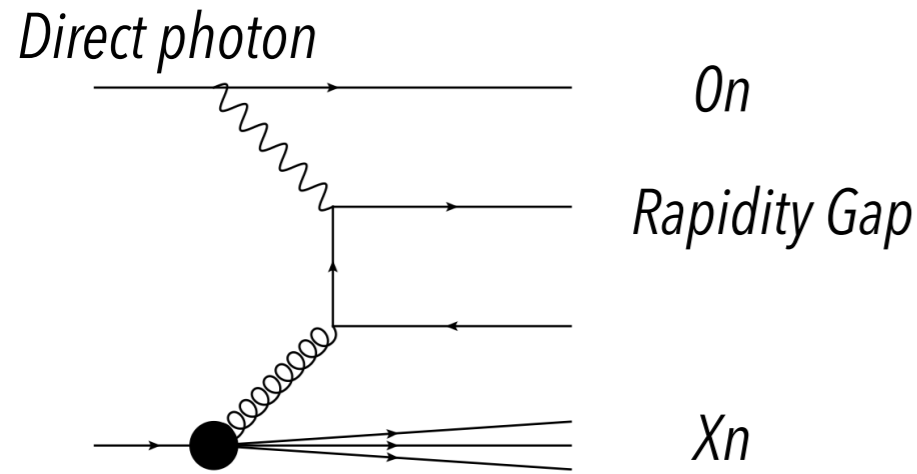
$$x_A \equiv \frac{M_{\text{jets}}}{\sqrt{s}} e^{-y_{\text{jets}}}$$

ATLAS-CONF-2022-021

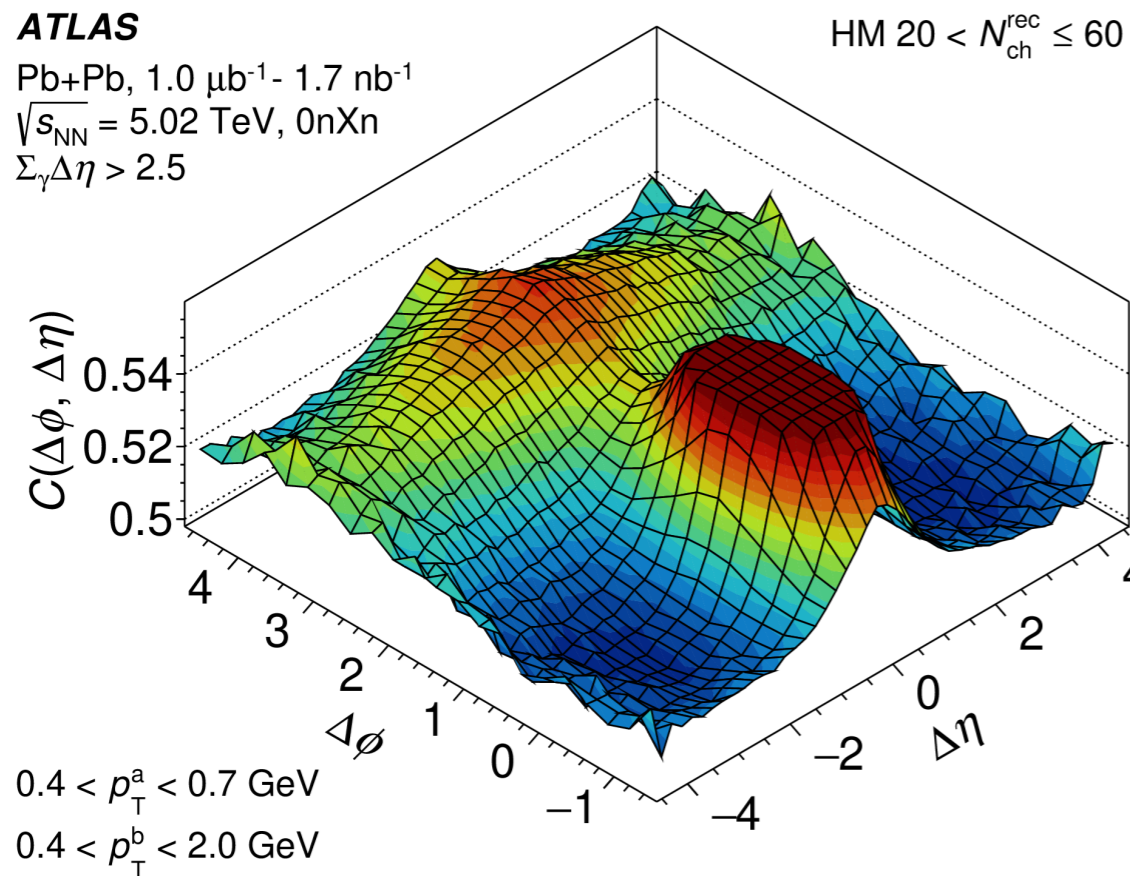


Clean probe to explore poorly constrained region at low- $x$  and intermediate  $Q^2$

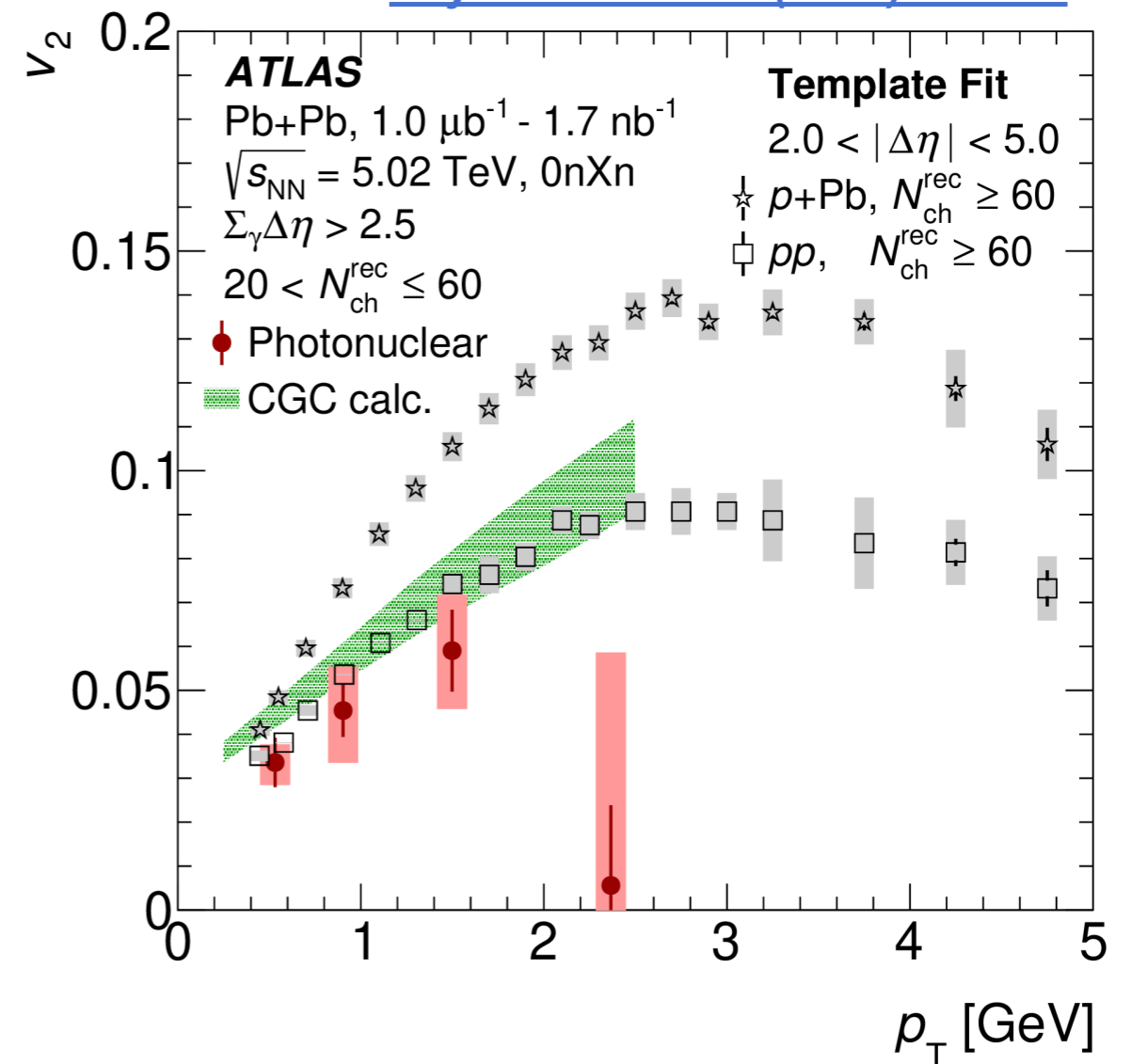
# Flow in photo-nuclear collisions



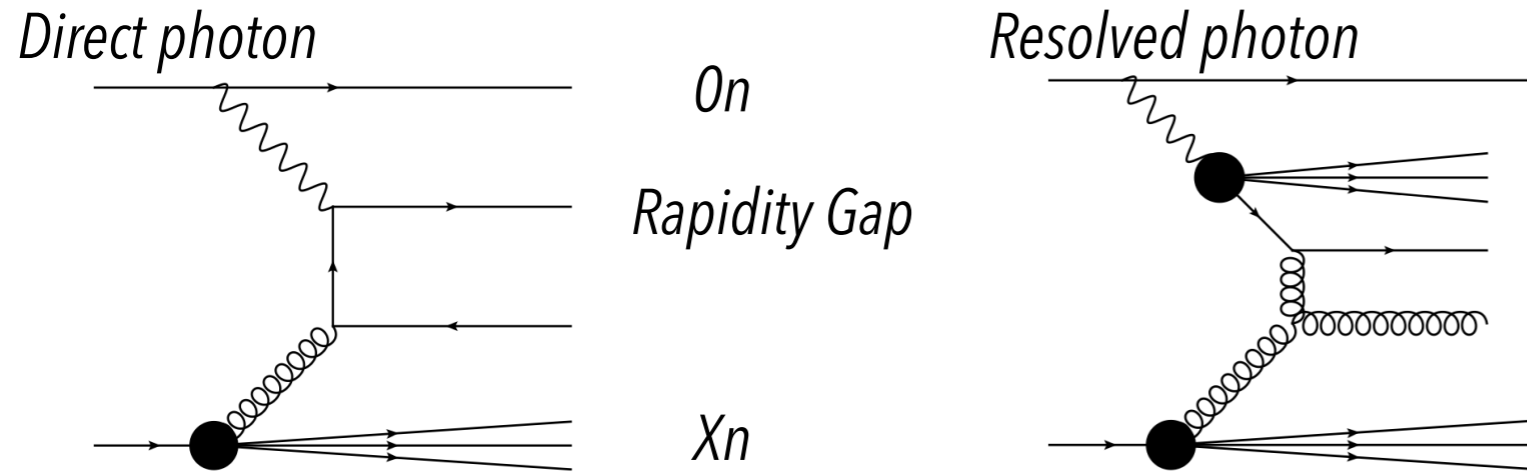
Comparable  $v_2$  what is seen in pp



[Phys. Rev. C. 104 \(2021\) 014903](#)

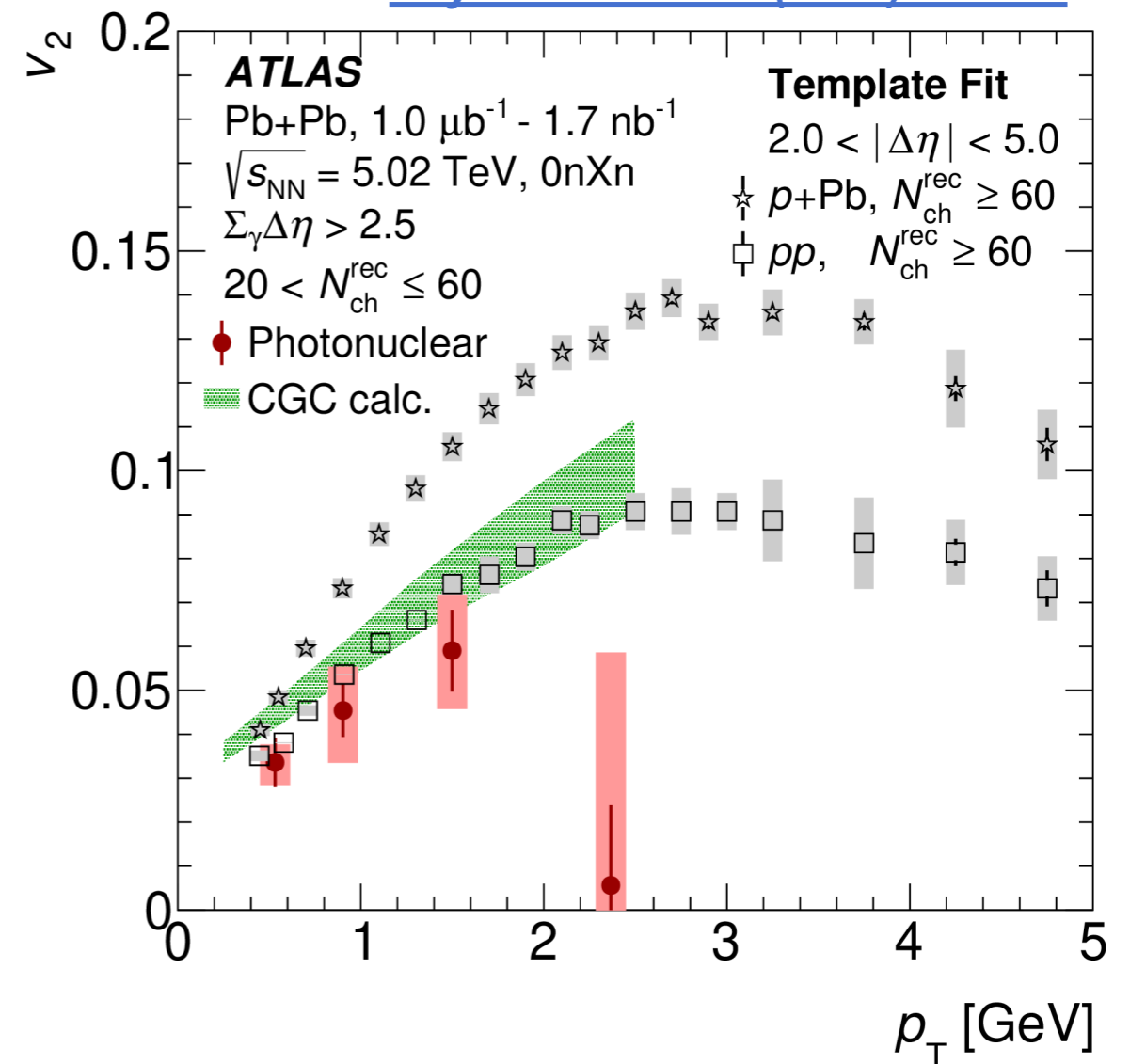
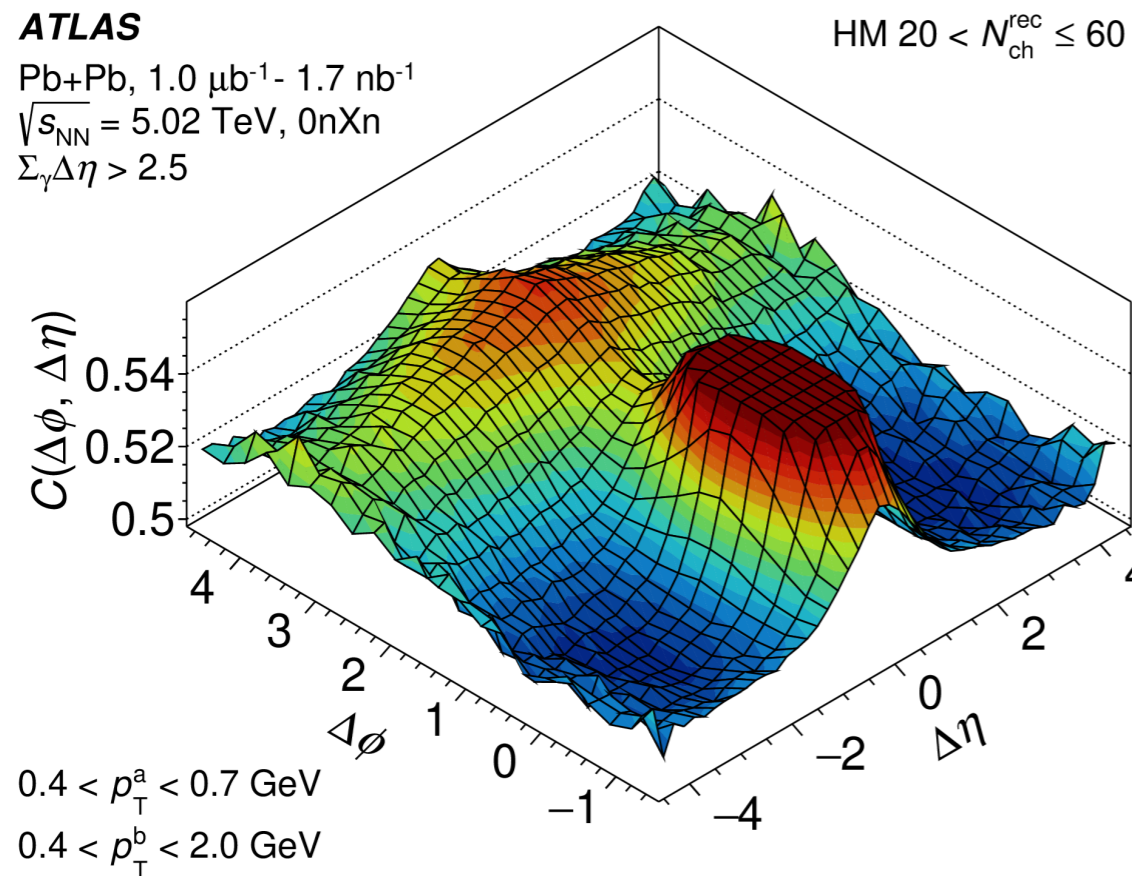


# Flow in photo-nuclear collisions



Comparable  $v_2$  what is seen in pp

[Phys. Rev. C. 104 \(2021\) 014903](#)



Using high statistics LHC data and new techniques bring us to era of precise measurements of QGP produced in heavy-ion collisions

We still need new measurements

- Each observable is sensitive to different aspect of probing the QGP
- Some observables are statistics hungry - looking forward for more data
- Interesting opportunity to study new collision system(s) (O+O LHC Run3, future of heavy ion program?)
- Growing interest in UPC physics

More details on ATLAS public results page:

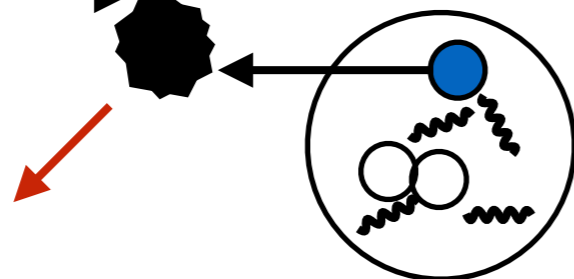
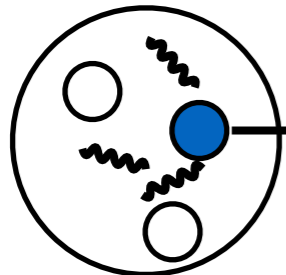
<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/HeavyIonsPublicResults>

# Jets in vacuum

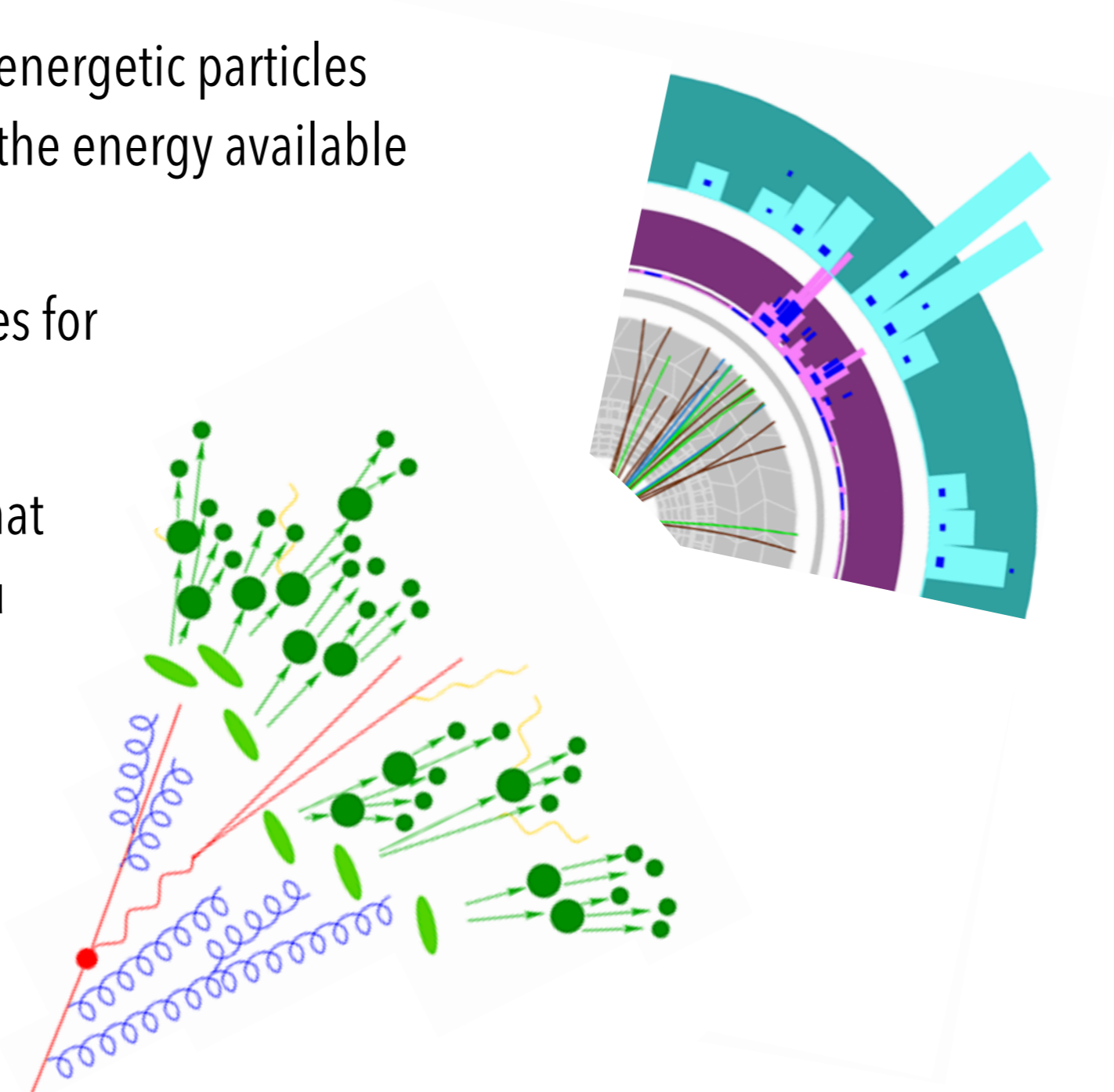
Jets: collimated showers of energetic particles that carry a large fraction of the energy available in the collisions

- In Theory: jets are proxies for hard-scattered partons
- In Experiment: jet is what your jet-finder gives you

Quark or gluon  
inside of nucleon



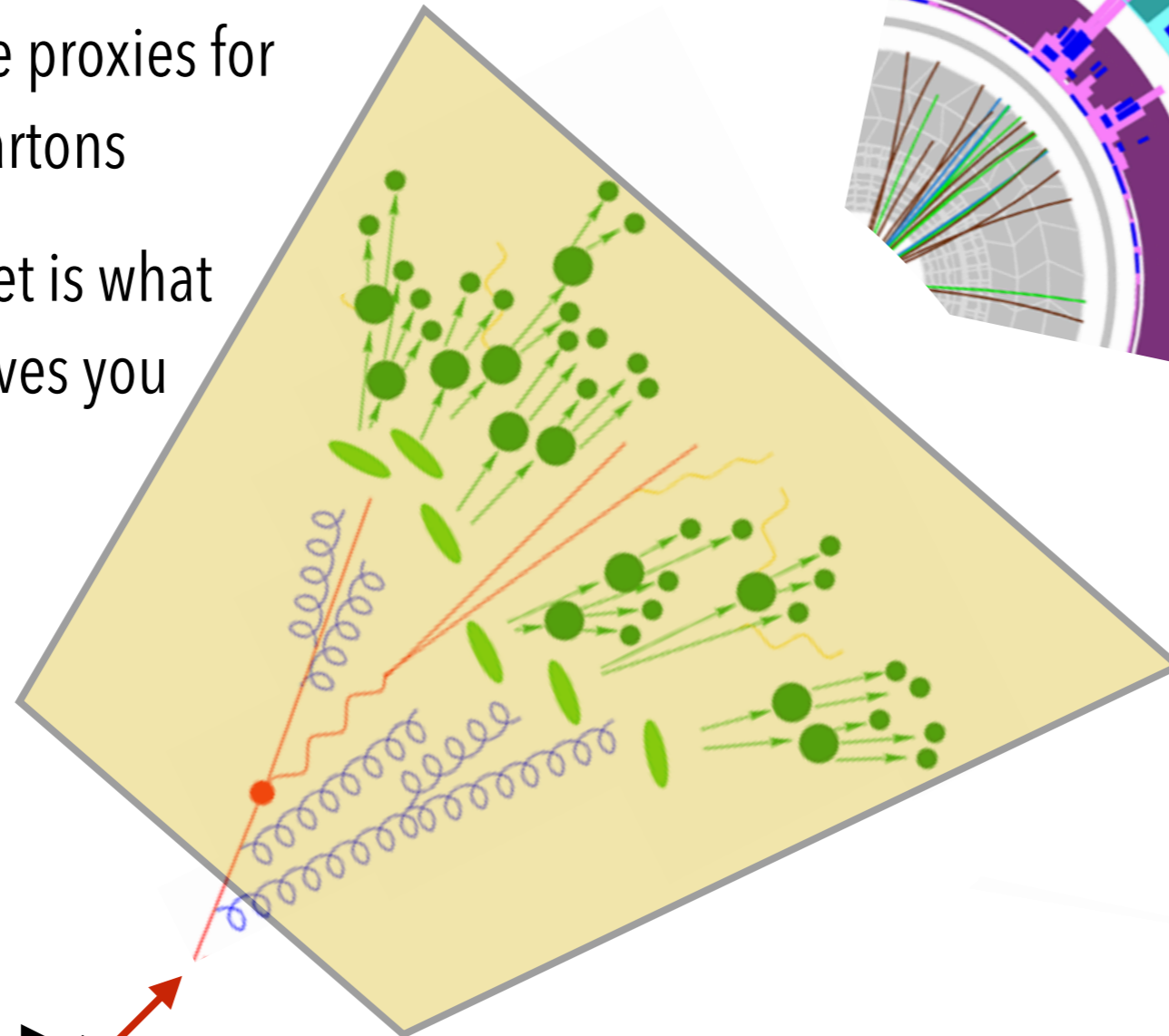
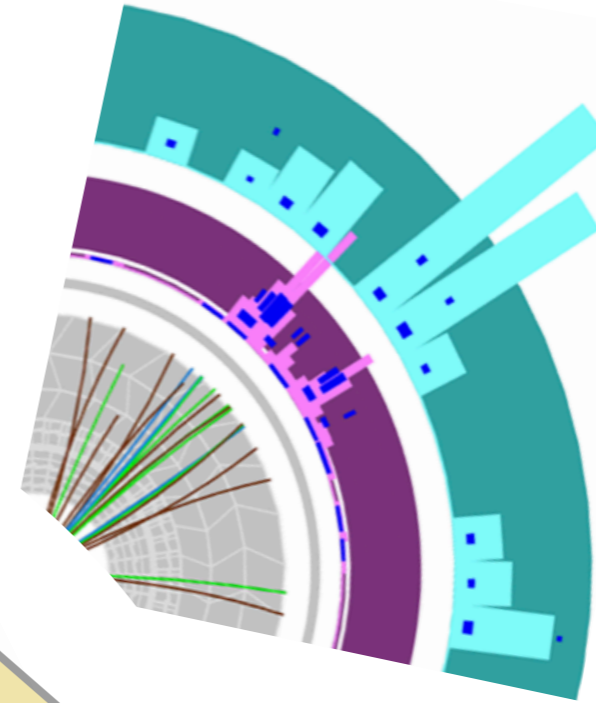
Quark or gluon  
inside of nucleon



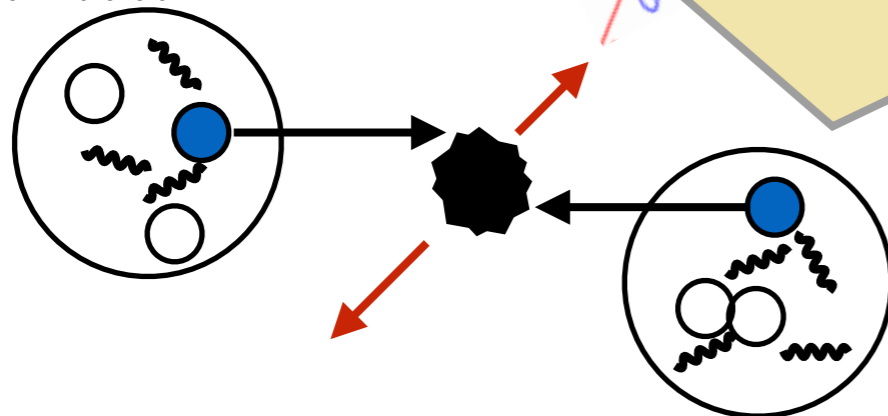
# Jets in the medium

Jets: collimated showers of energetic particles that carry a large fraction of the energy available in the collisions

- In Theory: jets are proxies for hard-scattered partons
- In Experiment: jet is what your jet-finder gives you



Quark or gluon  
inside of nucleon

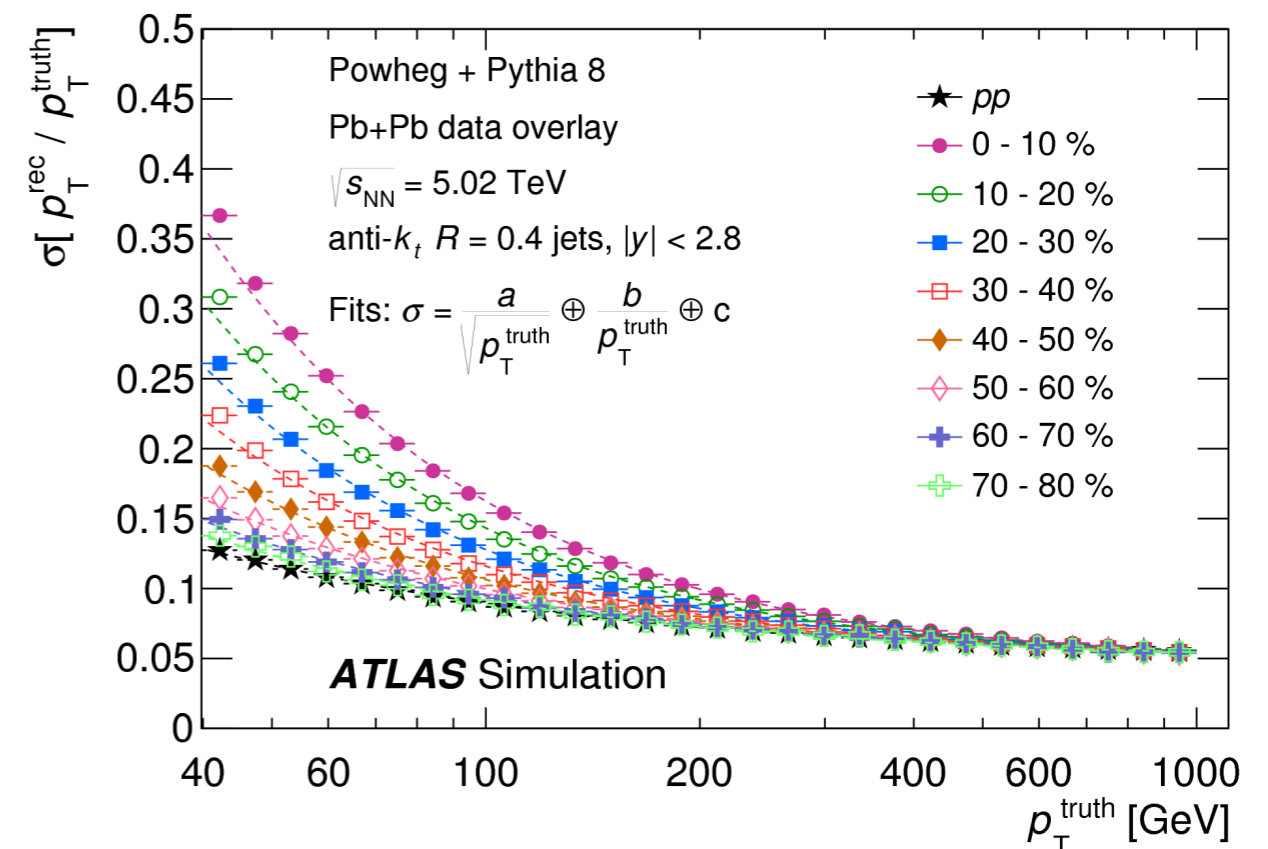
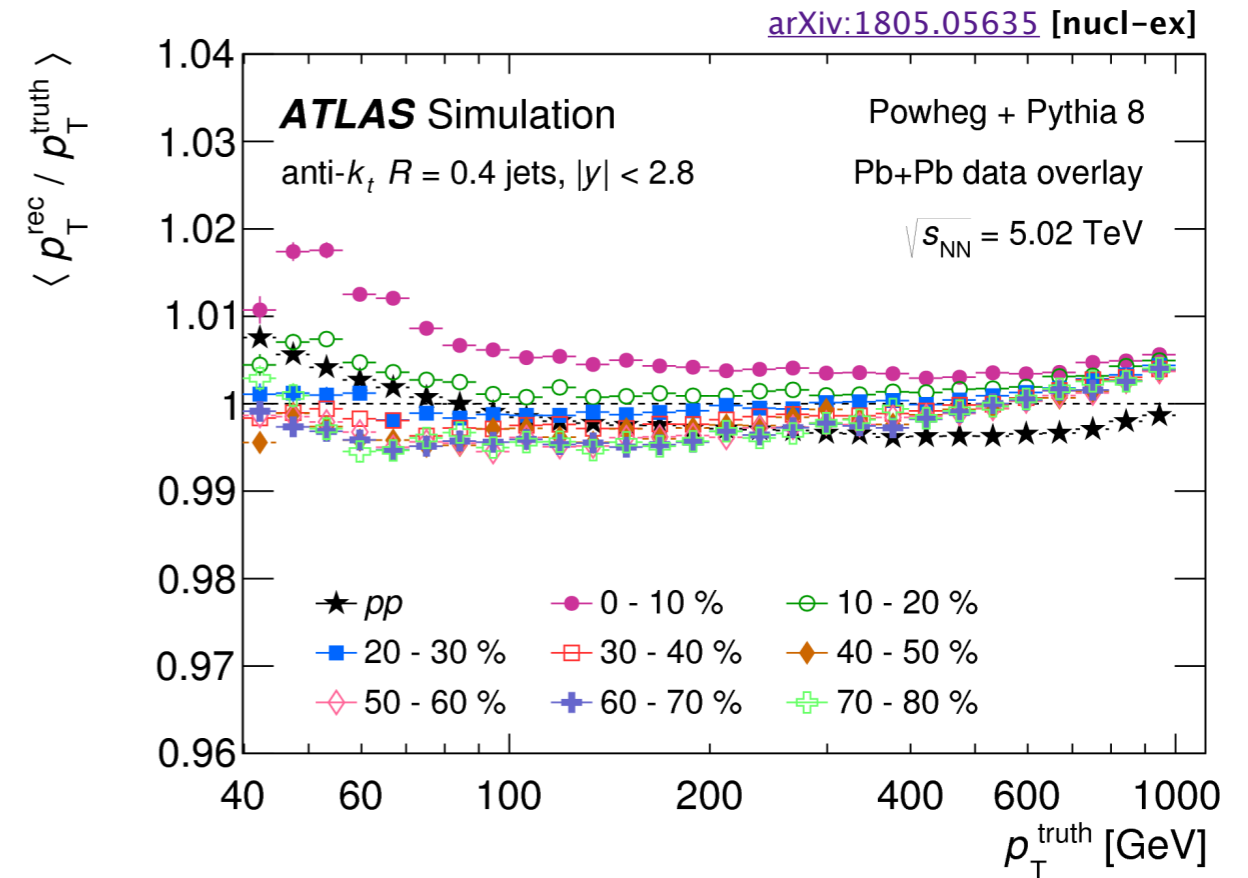
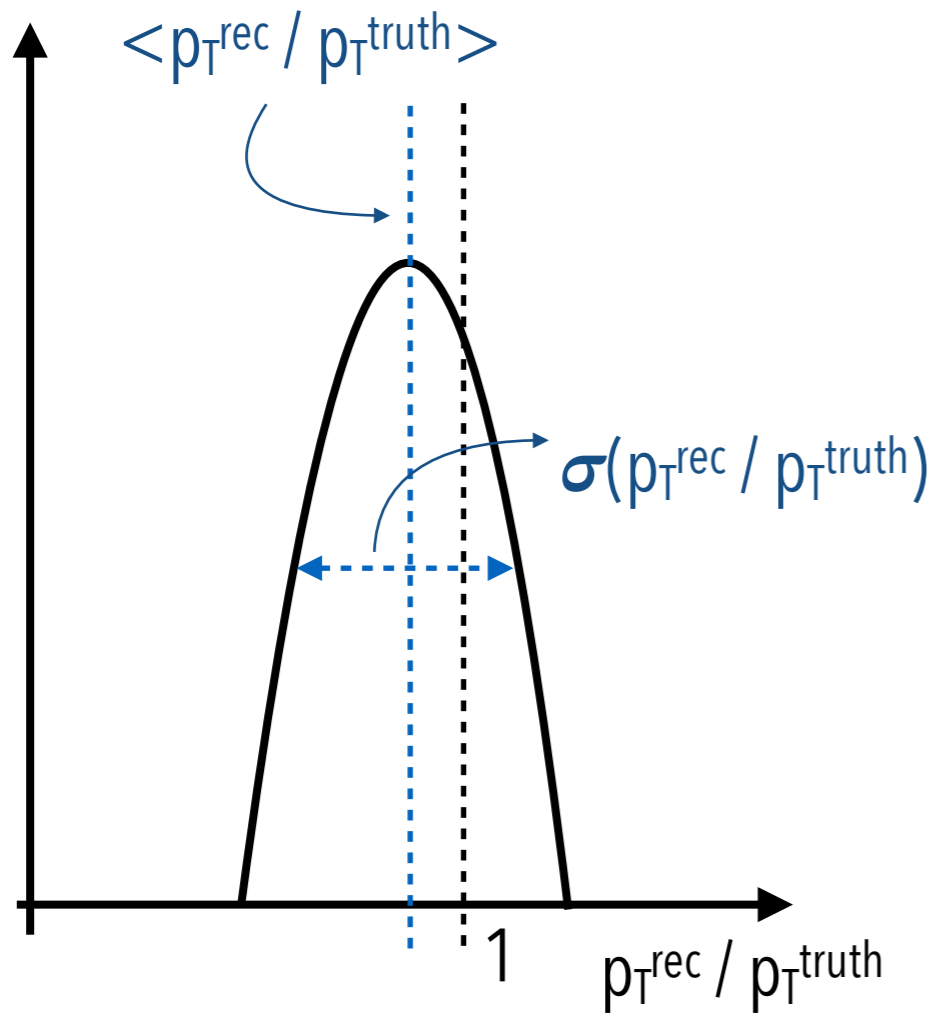


Quark or gluon  
inside of nucleon

Interactions of medium and colored probe: elastic scattering, medium induced radiation, "drag force", medium excitation ?

Initial cross-section unchanged by presence of medium (modulo change in nPDF)

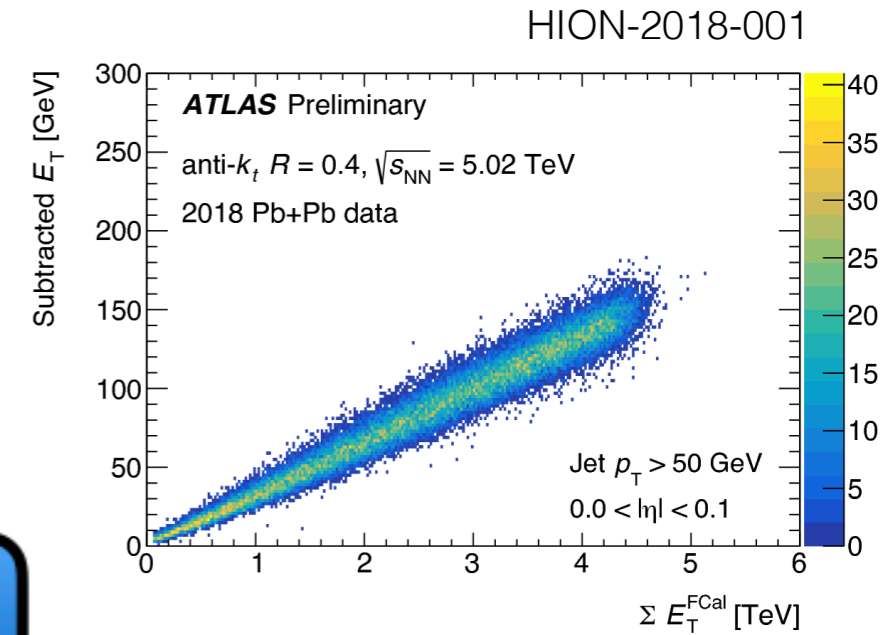
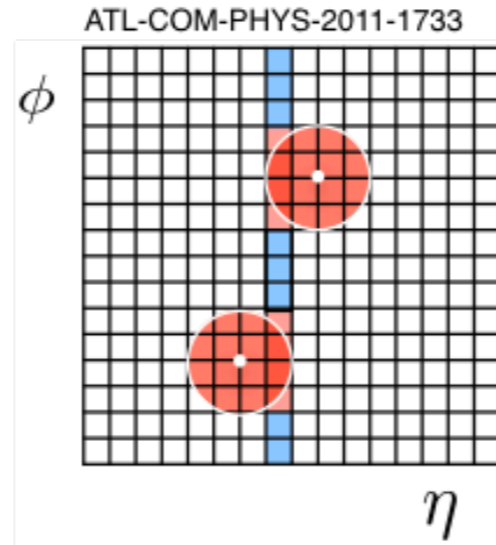
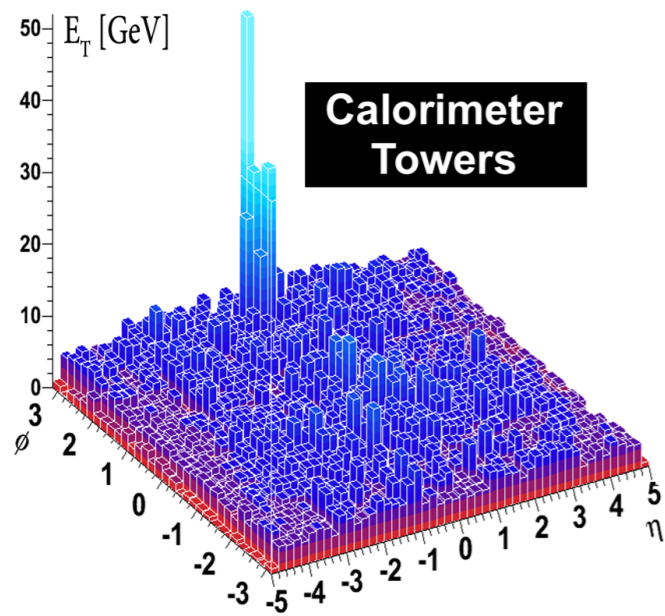
# Jets reconstruction - ATLAS heavy ion style



- Average response ( $\langle p_T^{\text{rec}} / p_T^{\text{truth}} \rangle$ ) within 1% from unity almost independent on centrality
- Jet energy resolution ( $\sigma(p_T^{\text{rec}} / p_T^{\text{truth}})$ ) dominated by the underlying event fluctuations.



# Jets reconstruction - ATLAS heavy ion style



Calorimeter towers

Average  $E_T$  density:  $\rho(\eta, \text{layer})$

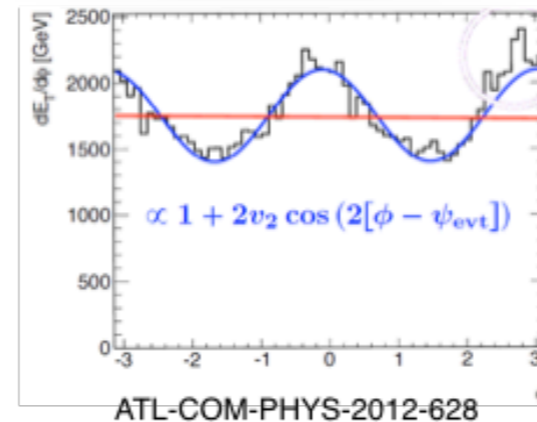
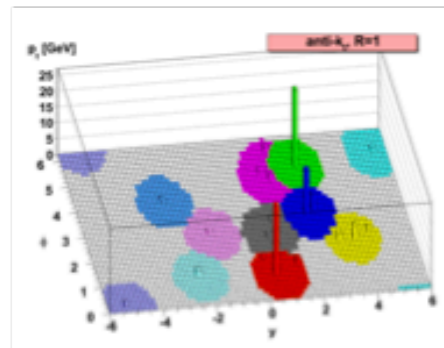
Jets

Reconstructed Jets

Anti- $k_t$  Algorithm

Flow modulation ( $v_2, v_3, v_4$ )

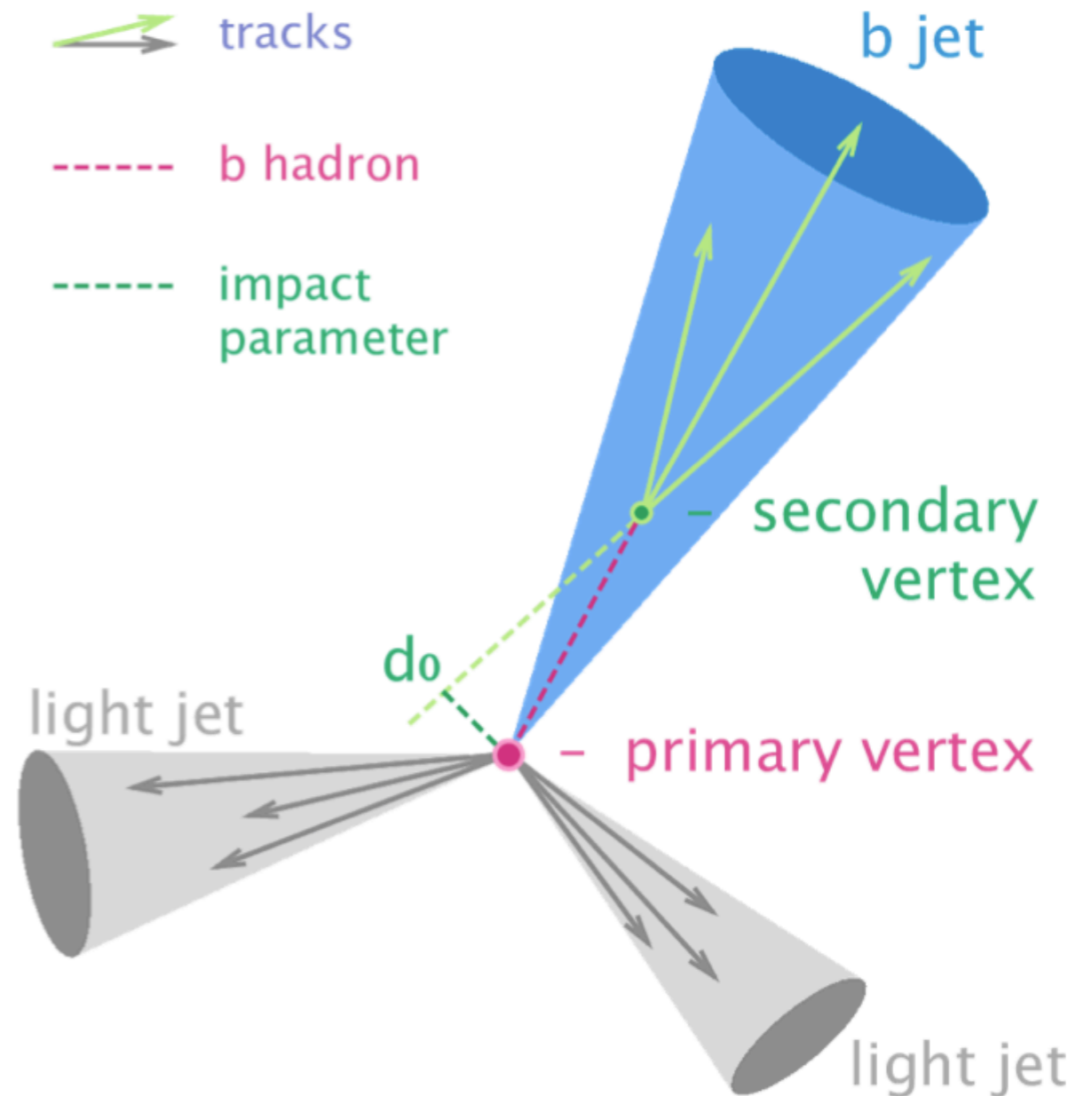
Iterative subtraction



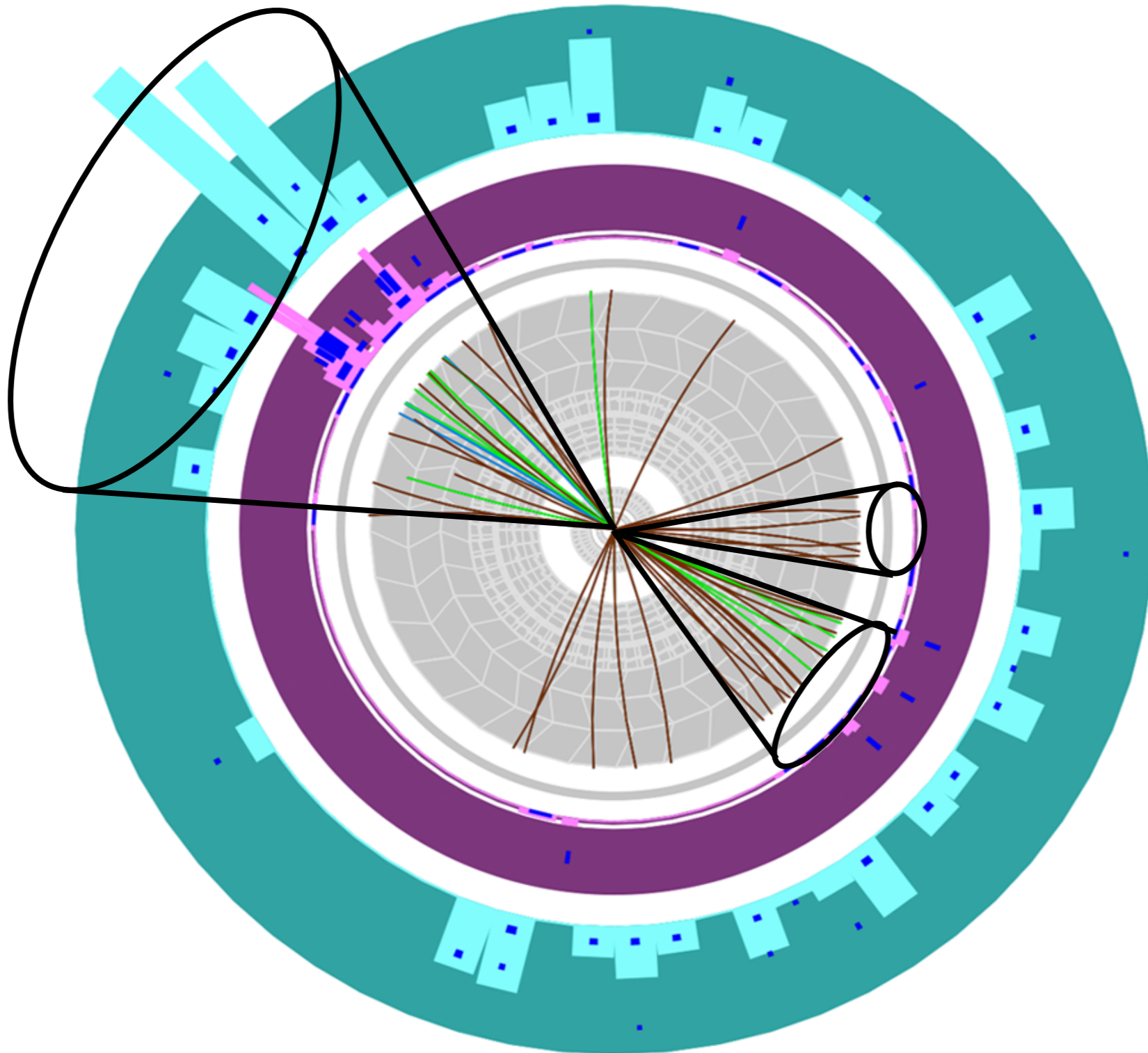
# Quenching of b-jets

The dependence of quenching on the type of parton that initiates the jet may provide insight into the underlying dynamics

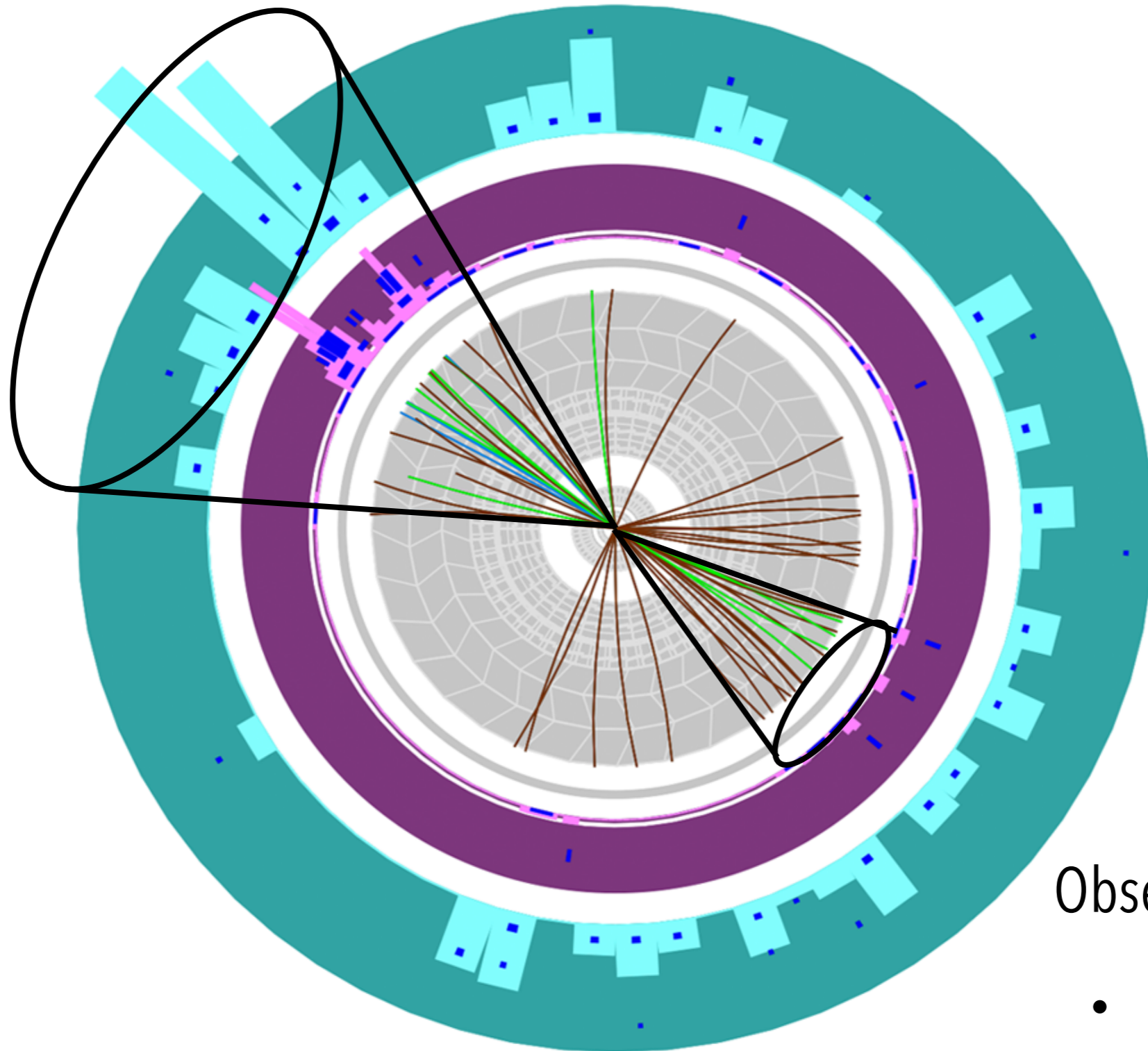
- Type of parton that initiates the jet is difficult to determine experimentally
- Machine learning techniques used in experiments to identify b-jets



# Inclusive jet spectra

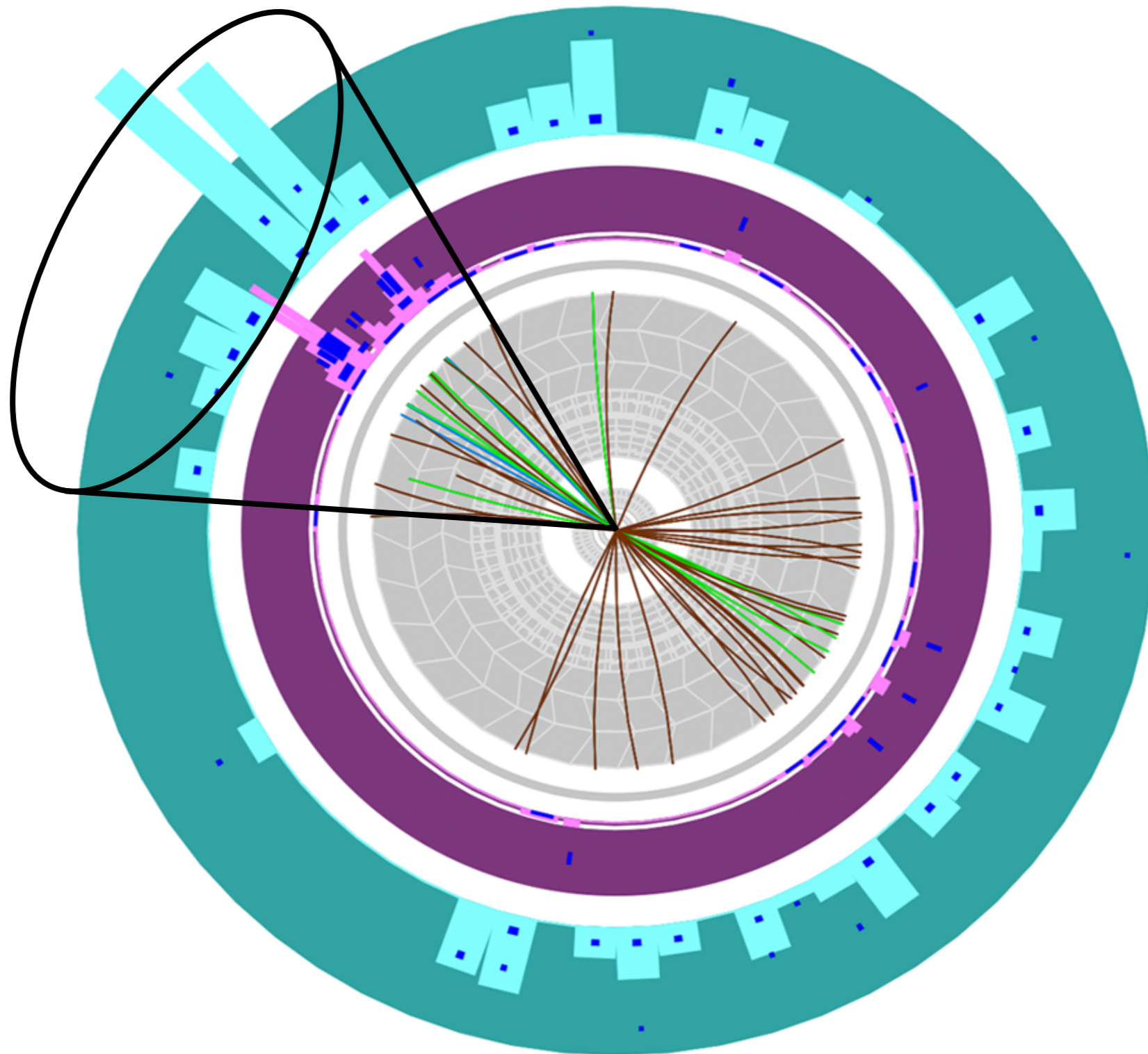


# Path length dependent energy loss



Observables:

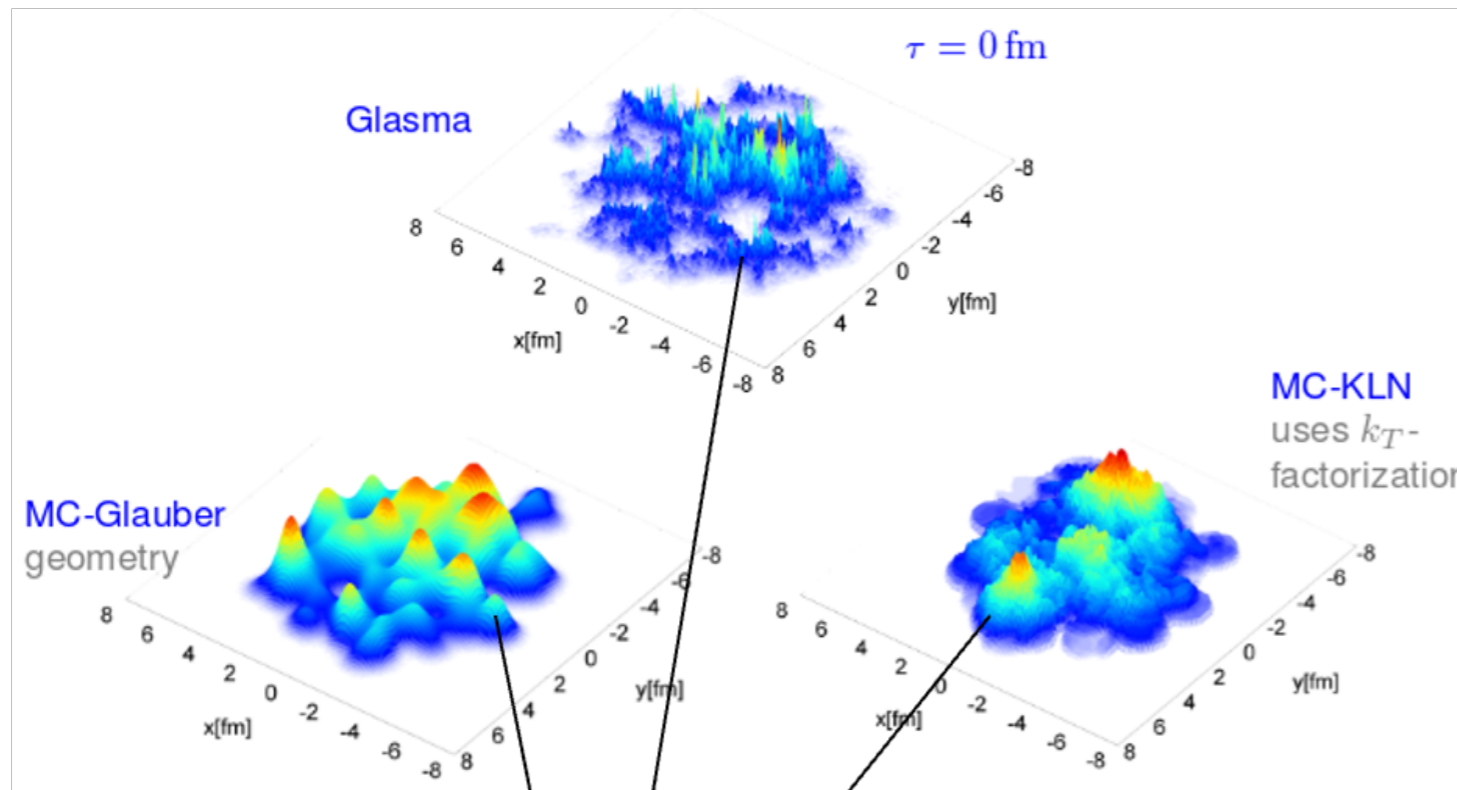
- Dijet asymmetry
- Jet  $v_n$



# Flow measurements

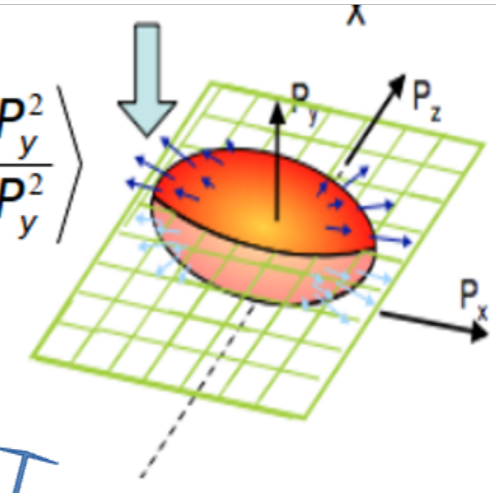
Method	Pros	Cons
<b>Event plane</b>	Simple to use Used widely by experiment and theory	Results susceptible to non-flow effects Comparison between experiment depend on event plane resolution
<b>Scalar product</b>	Simple to use Uniform treating of detector acceptance	Results susceptible to non-flow effects
<b>Two particle correlation</b>	Simple to use Used widely by experiment and theory	Results susceptible to non-flow effects
<b>Multiparticle cumulants</b>	Large reduction of non-flow effects	Difficult in computation (large cpu resources) Requires large signal
<b>Lee Yang Zeros</b>	Complete reduction of non-flow effects	Difficult in computation (large cpu resources) Requires large signal

# Stages of HI collision

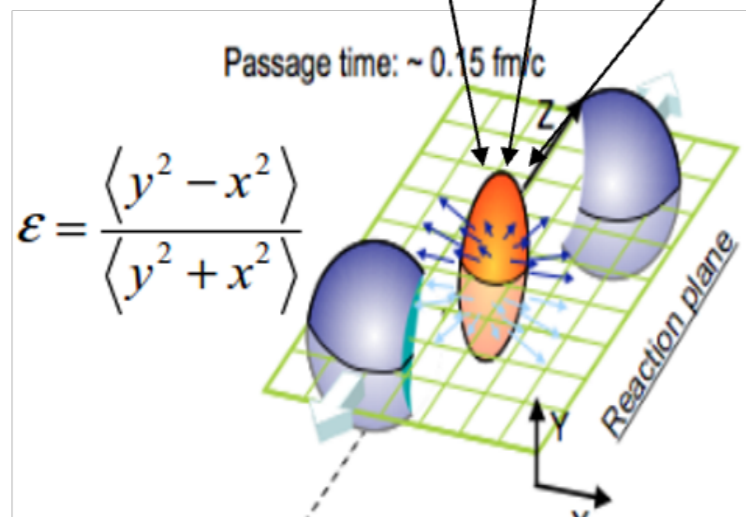


Particles observed in the detector.

$$V_2 = \left\langle \frac{P_x^2 - P_y^2}{P_x^2 + P_y^2} \right\rangle$$

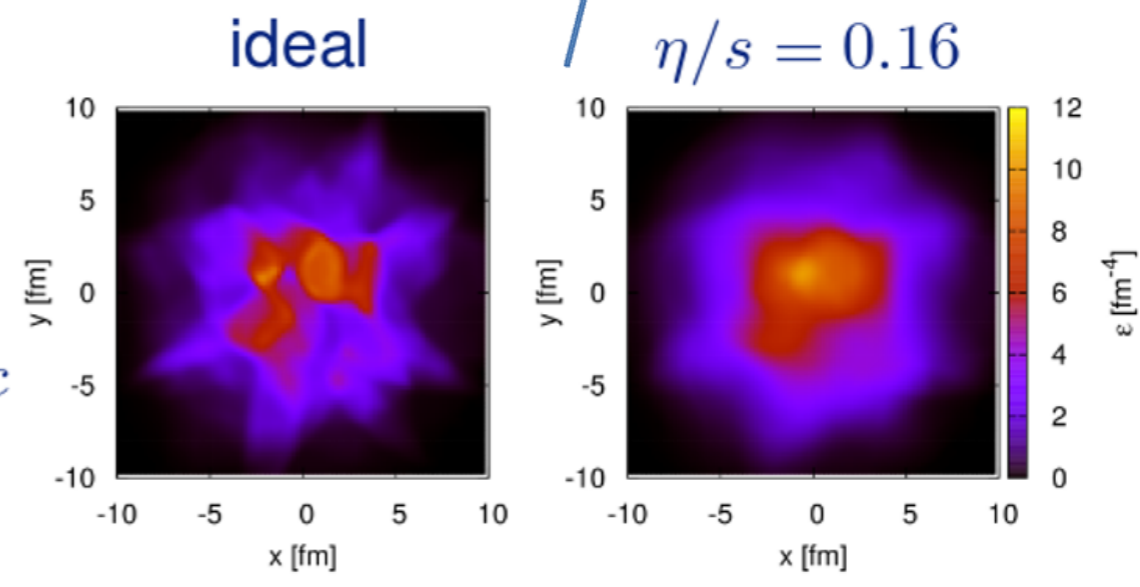


Hadronization – transition from QGP to hadronic matter.



Just after collision – large eccentricity ( $\epsilon$ ) in the nucleons distributions gives rise to the internal pressure gradients.

evolve to  
 $\tau = 6 \text{ fm}/c$

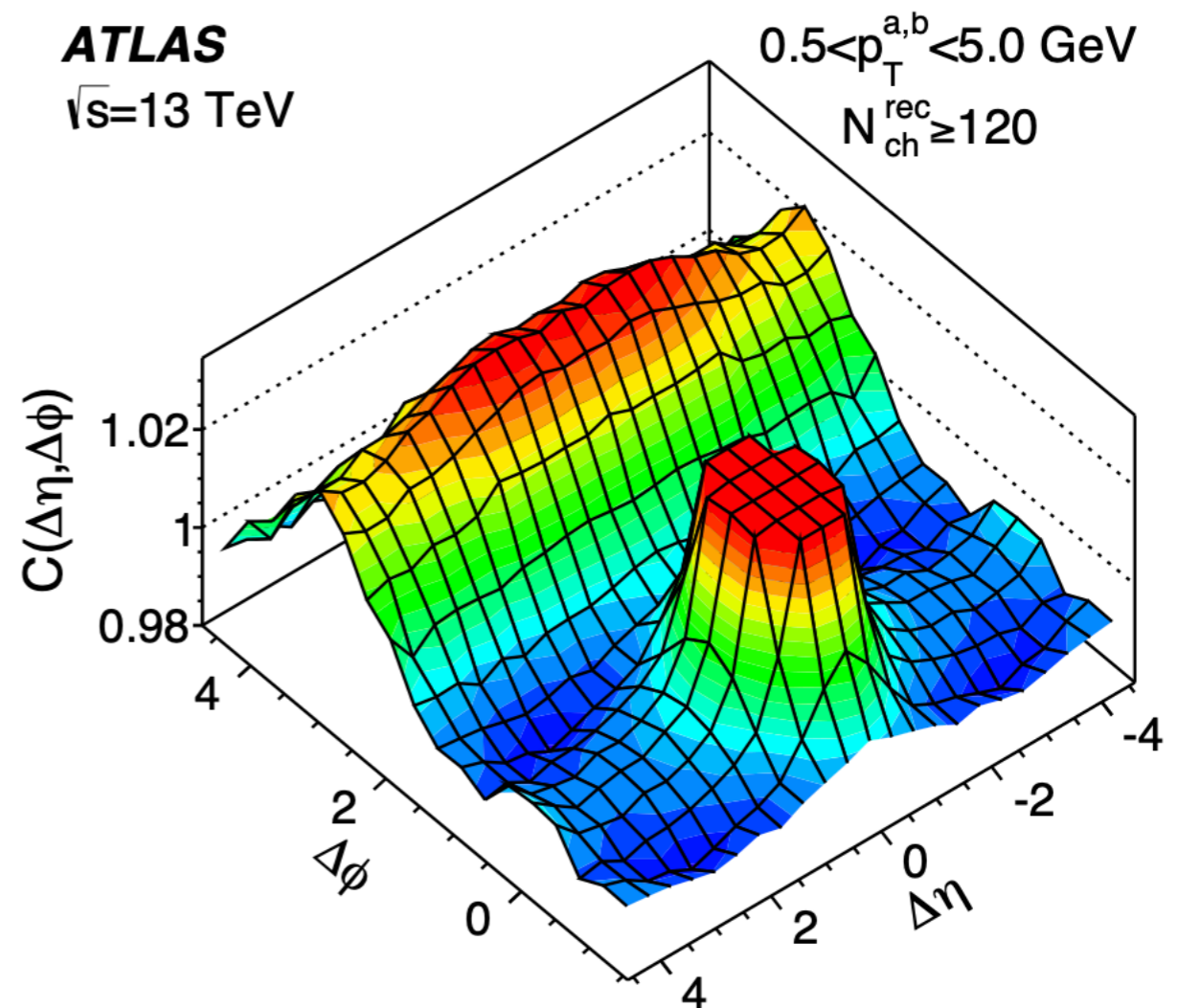
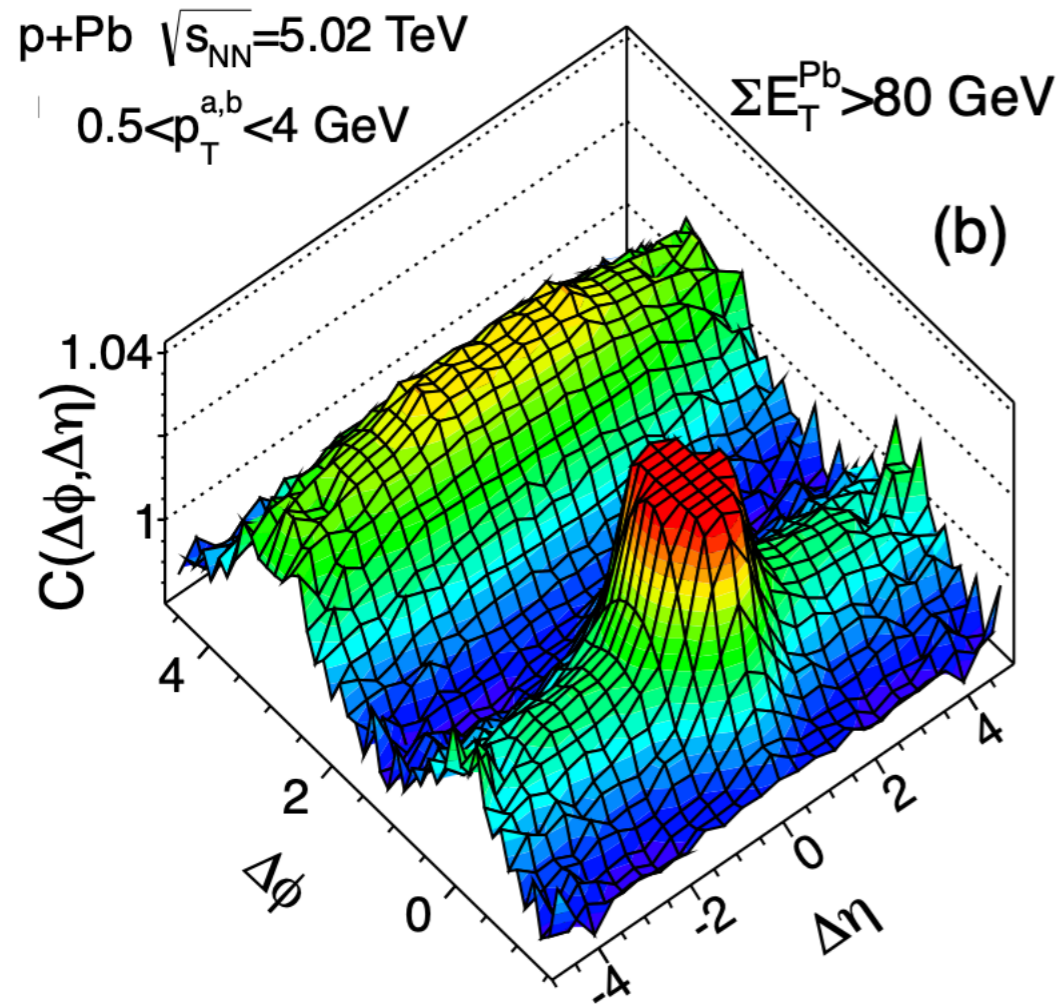


Hydrodynamic evolution; properties of QGP state.

# Magic of the template fit

PRL 110 182302 (2013)

PRL 116 172301 (2016)

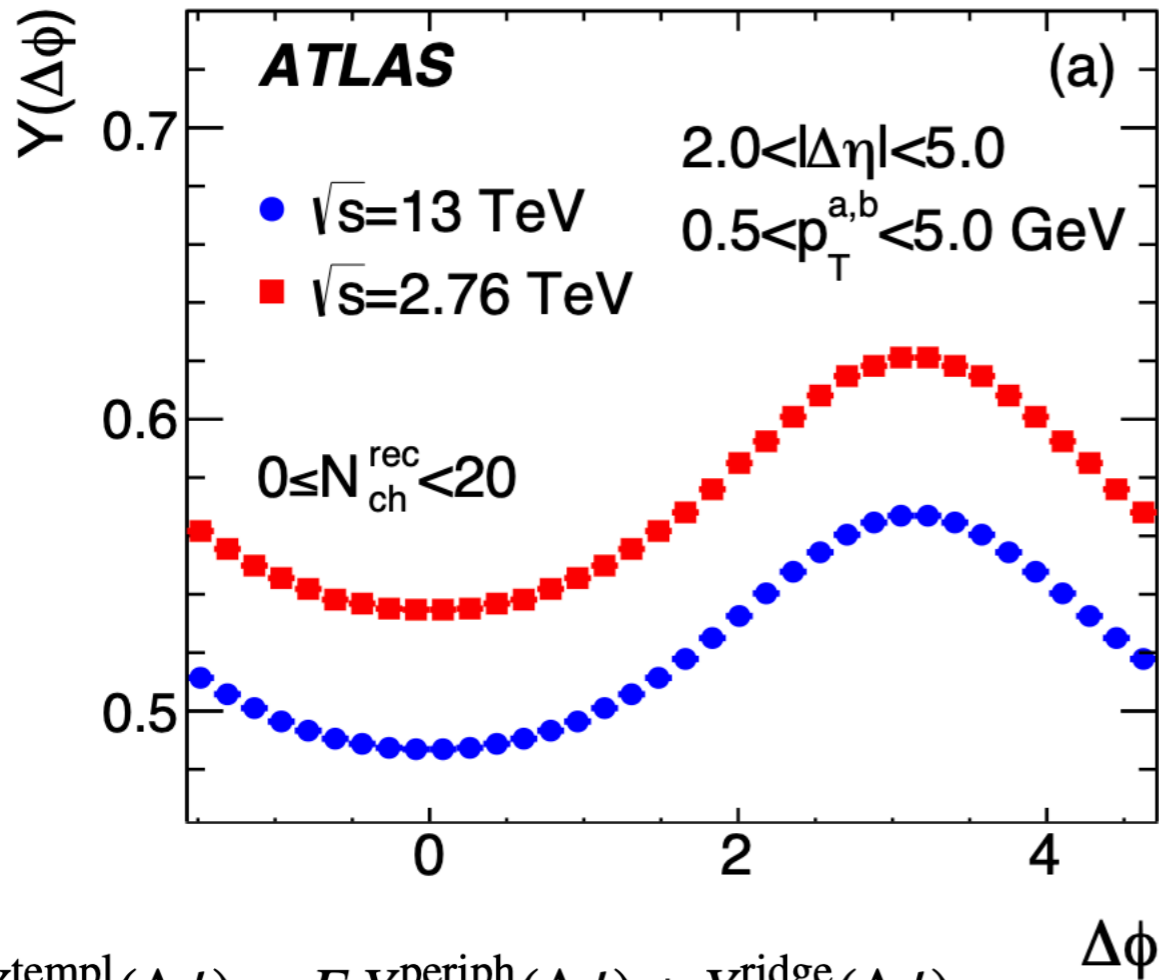


Looking for  $v_n$  in smaller systems is challenging due the presence of the large non-flow background

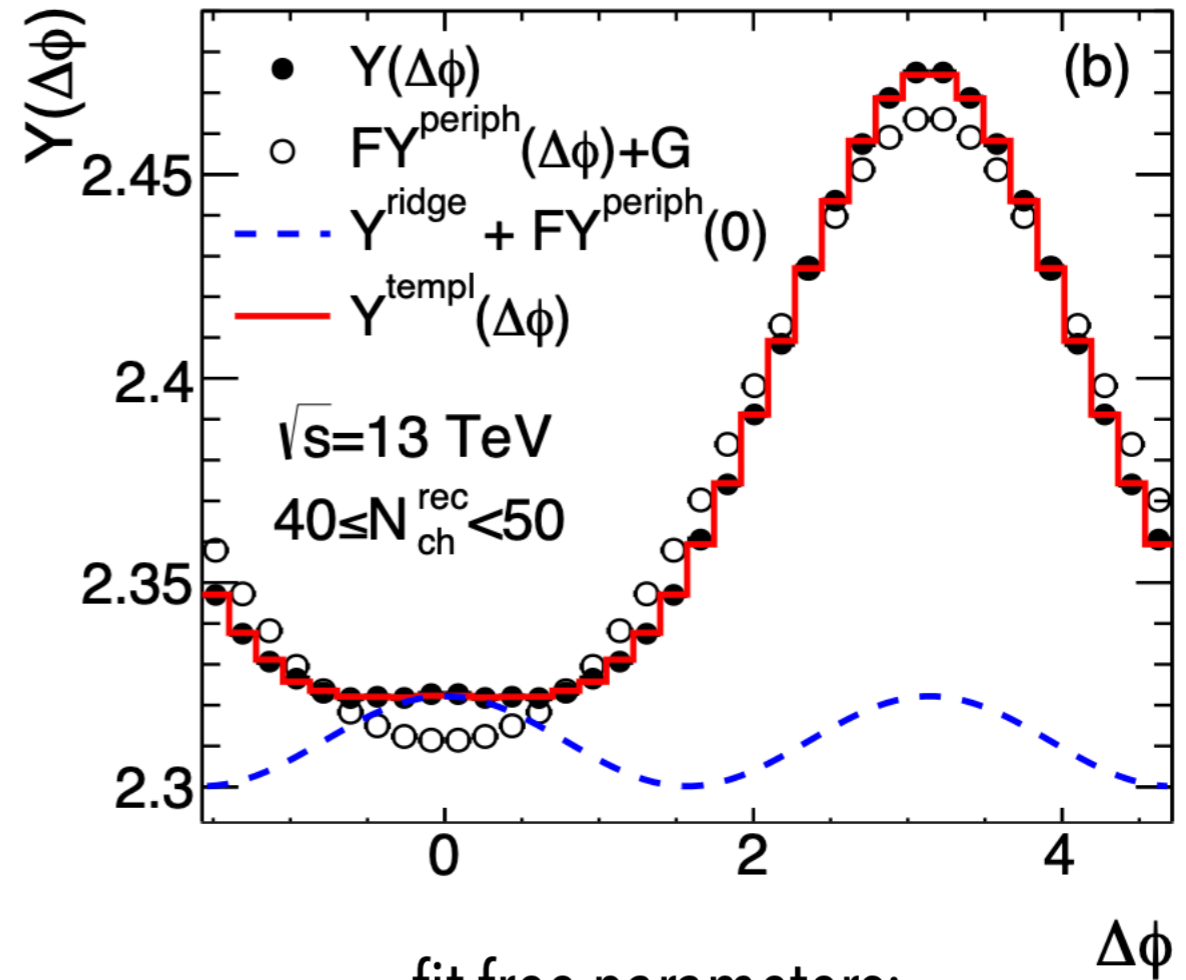


# Magic of the template fit

$Y^{\text{periph}}(\Delta\phi), N_{\text{ch}}^{\text{rec}} < 20$



$40 < N_{\text{ch}}^{\text{rec}} < 50$



$$Y^{\text{templ}}(\Delta\phi) = F Y^{\text{periph}}(\Delta\phi) + Y^{\text{ridge}}(\Delta\phi),$$

$$Y^{\text{ridge}}(\Delta\phi) = G (1 + 2v_{2,2} \cos(2\Delta\phi)),$$

fit free parameters:

$F$  &  $v_{2,2}$

$G$  fixed such that:

$$\int_0^\pi d\Delta\phi Y^{\text{templ}} = \int_0^\pi d\Delta\phi Y$$