

# Recent topics on Mesic atoms and Mesic nuclei

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*Satoru Hirenzaki (Nara Women's University, Japan)*



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## **【0】 Introduction**

**【1】  $\eta'$ (958) mesic nucleus -- after GSI exp.--**

**【2】  $\bar{p}$  : Hadronic systems from Anti-Protonic atom**

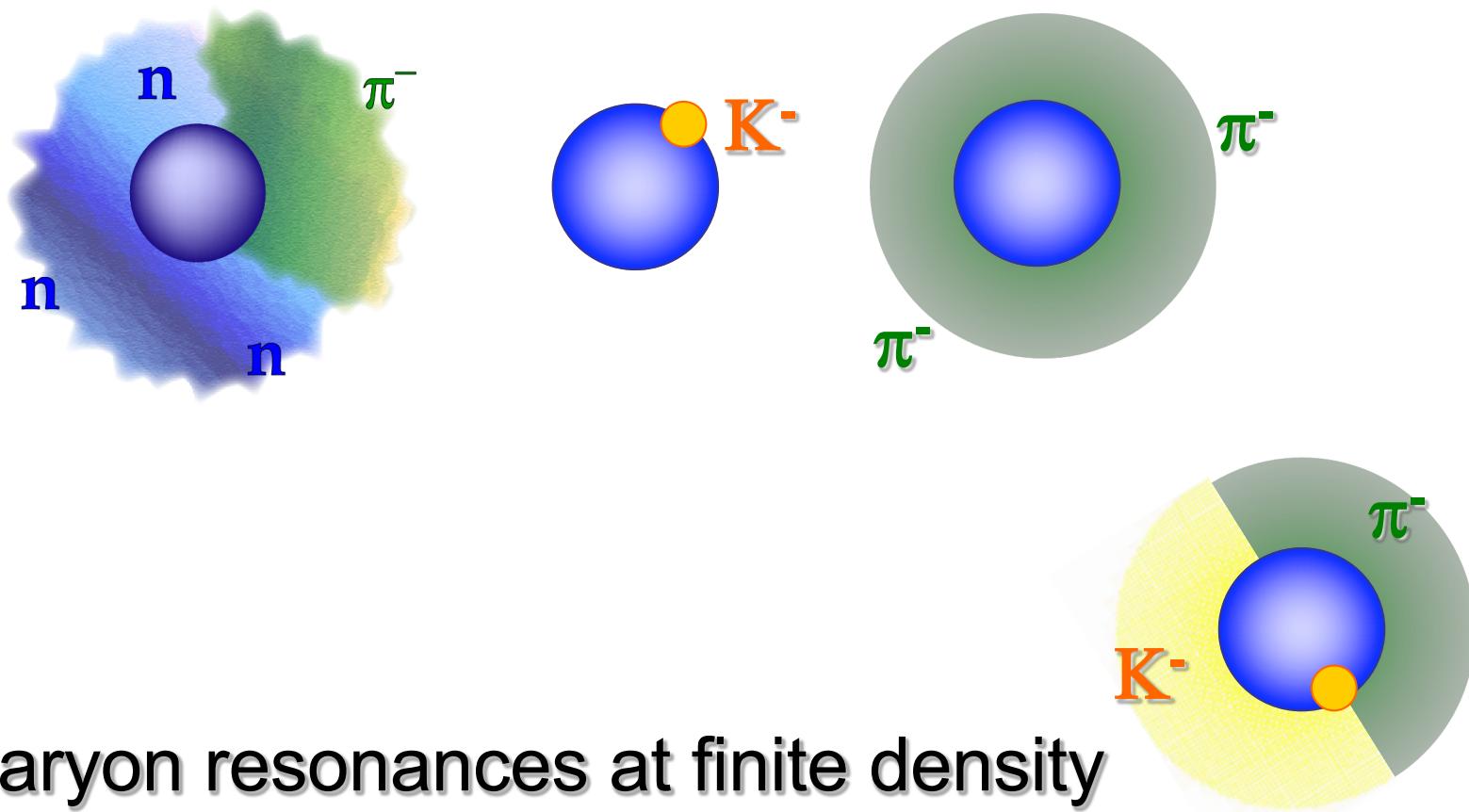
**【3】  $\eta$  mesic nucleus**

**【4】 Summary**

## 【0】Introduction

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### (1) New exotic Hadron many body systems



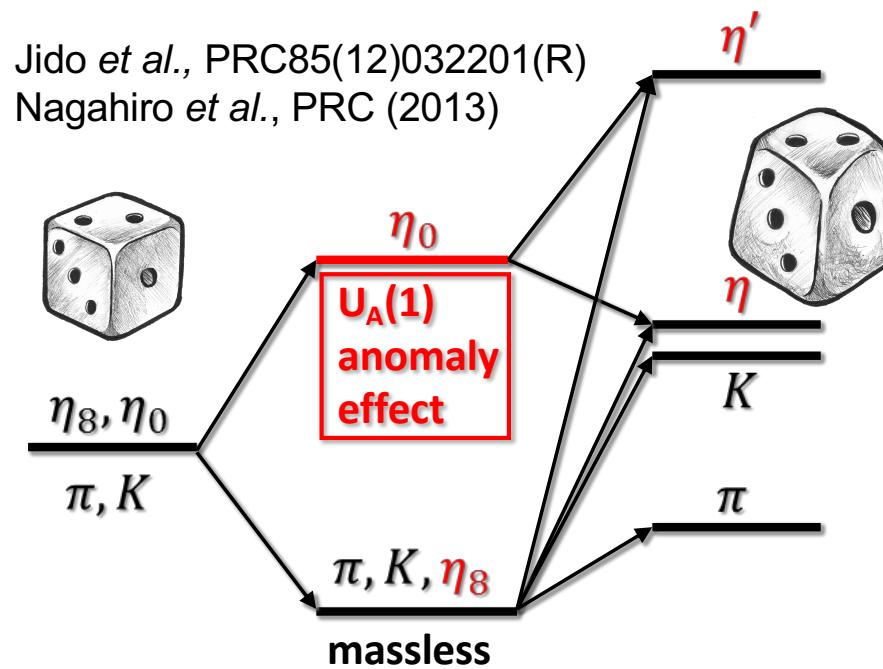
### (2) Baryon resonances at finite density

$N^*(1535)$ 、 $\Lambda(1405)$ 、、

# [0] Introduction

schematic view of the mass of  $\pi, K, \eta$  &  $\eta'$

## (3) Aspects of the Strong Int. Symmetry



$$m_q, m_s = 0$$

$$\langle \bar{q}q \rangle = 0$$

ChS  
manifest

$$m_q, m_s = 0$$

$$\langle \bar{q}q \rangle \neq 0$$

dynamically  
broken

$$m_q, m_s \neq 0$$

$$\langle \bar{q}q \rangle \neq 0$$

dyn. & explicitly  
broken

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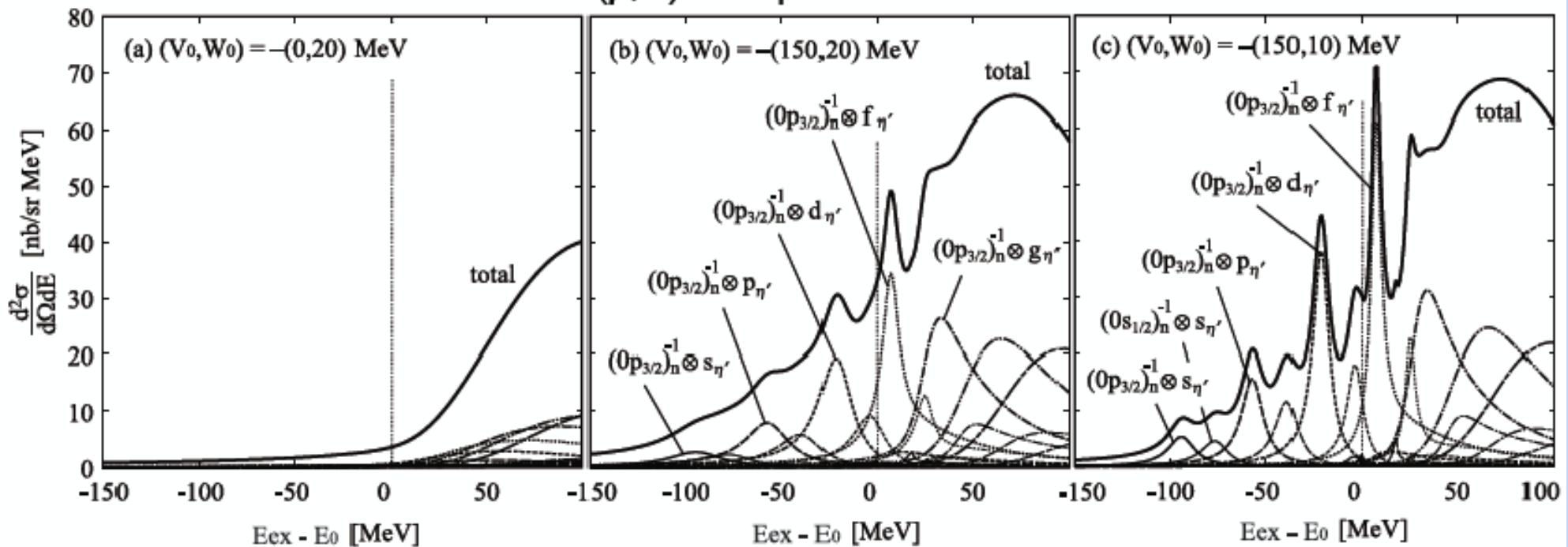
## 【0】 Introduction

## 【1】 $\eta'$ (958) mesic nucleus -- after GSI exp.--

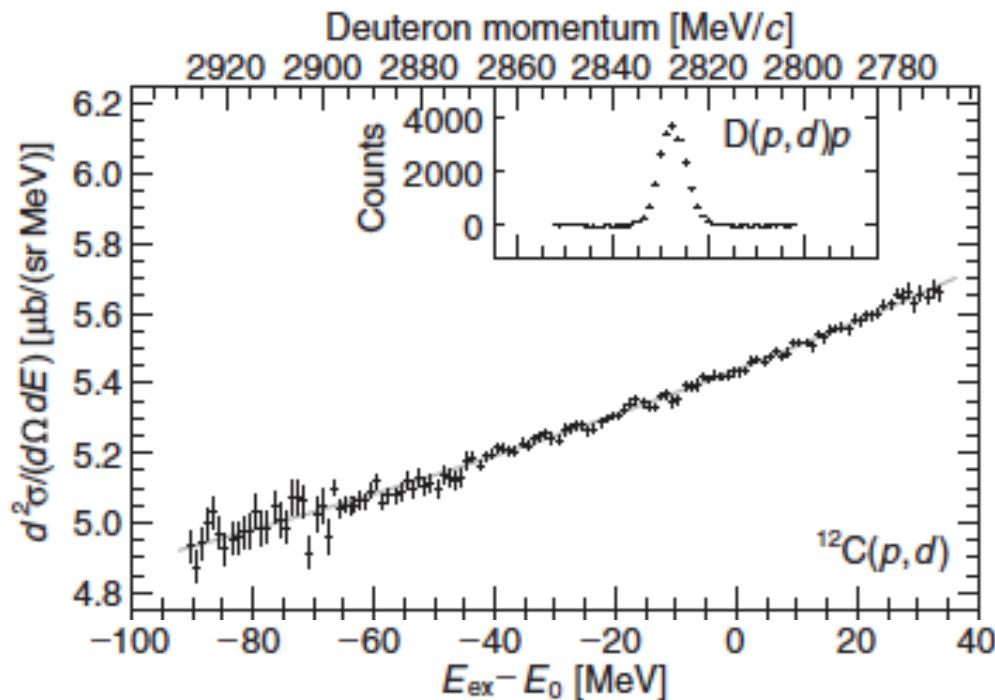
# Theoretical predictions

$$V_{\eta'}(r) = (V_0 + iW_0) \frac{\rho(r)}{\rho_0}$$

$^{12}\text{C}(p,d)$  at  $T_p = 2.50$  GeV



Nagahiro et al., PRC87(13)045201.



## Experimental Results

**FIG. 2.** (Top panel) Excitation spectrum of  $^{11}\text{C}$  measured in the  $^{12}\text{C}(p, d)$  reaction at a proton energy of 2.5 GeV. The abscissa is the excitation energy  $E_{ex}$  referring to the  $\eta'$  emission threshold  $E_0 = 957.78$  MeV. The overlaid gray solid curve displays a fit of the spectrum with a third-order polynomial. The upper horizontal axis shows the deuteron momentum scale. (Inset) Deuteron

Contour on  
a potential  
parameter plane

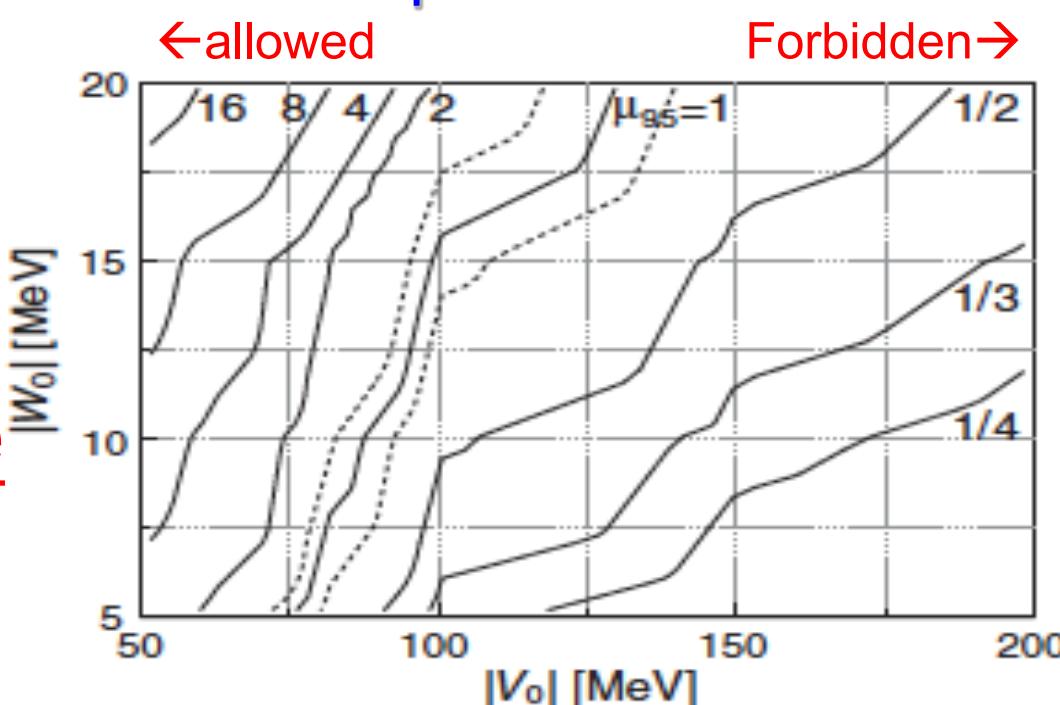


FIG. 4. A contour plot of  $\mu_{95}$  (the solid curves), upper limit of the scale parameter  $\mu$  at the 95% C.L., on a plane of real and imaginary potential parameters ( $|V_0|, |W_0|$ ). The limits have been evaluated for the potential-parameter combinations  $(V_0, W_0)$  in  $\{-50, -100, -150, -200\} \times \{-5, -10, -15, -20\}$  and  $\{-60, -80\} \times \{-5, -10, -15\}$  MeV and linearly interpolated in between. Dashed curves show a band of  $\mu_{95} = 1$  contour indicating the systematic errors. Regions for  $\mu_{95} \leq 1$  are excluded by the present analysis.

- No clear peak. Upper limit is obtained.
- After GSI exp., for further studies,
  - Theoretical eta(958)-N and -Nucleus interactions
  - Semi-exclusive exp. ( $p, dp$ ) at FAIR/J-PARC
  - Self-consistent structure of eta(958)-nucleus  
=> Possible Deformation (Compression) of Nucleus  
and Effects to the spectra.
- Meanings of the exp. upper limit may be changed.

# Structure of $\eta'$ mesonic nuclei in a relativistic mean field theory

Daisuke Jido<sup>1,3,\*</sup>, Hanayo Masutani<sup>1</sup>, and Satoru Hirenzaki<sup>2</sup> Prog. Theor. Exp. Phys. **2019**, 053D02 (22 pages)

$$\begin{aligned} \mathcal{L} = & \bar{\psi} \left[ i\gamma_\mu \left\{ \partial^\mu + ig_\omega \omega^\mu + ig_\rho \rho^\mu \frac{\tau^3}{2} + ieA^\mu \frac{1+\tau^3}{2} \right\} \right] \psi - \bar{\psi} (m - g_\sigma \sigma) \psi \\ & + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - \frac{1}{2} m_\sigma^2 \sigma^2 - \frac{1}{3} b m g_\sigma^3 \sigma^3 - \frac{1}{4} c g_\sigma^4 \sigma^4 \\ & - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega^2 - \frac{1}{4} R_{\mu\nu} R^{\mu\nu} + \frac{1}{2} m_\rho^2 \rho_{03}^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + \frac{1}{2} \partial_\mu \eta' \partial^\mu \eta' - \frac{1}{2} m_{\eta'}^2 \eta'^2 + g_{\sigma\eta'} m_{\eta'} \eta'^2 \sigma. \end{aligned}$$

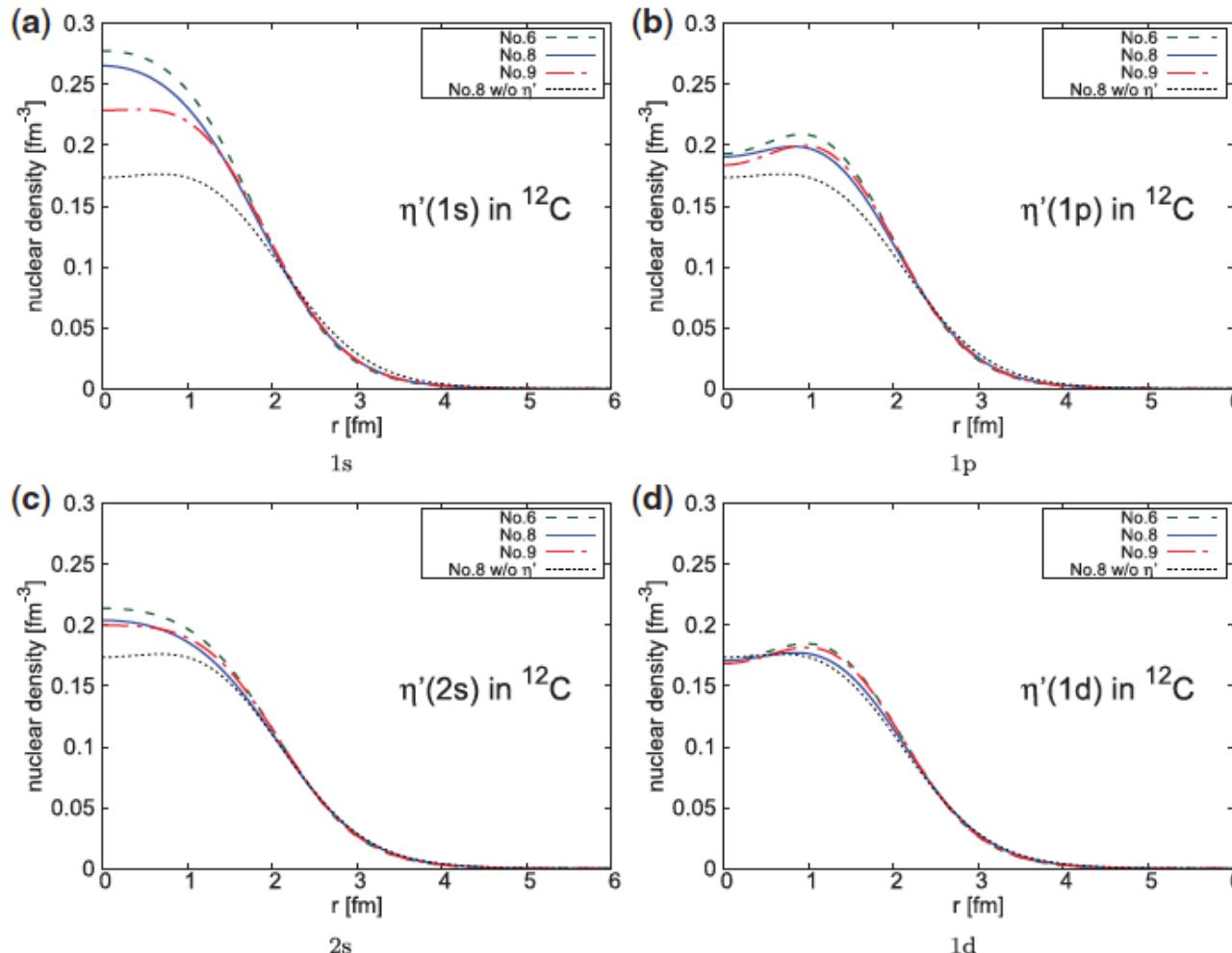
**Table 4.** Coupling constants of the  $\sigma$ - $\eta'$  interaction. These are determined so as to reproduce the effective  $\eta'$  mass 80 MeV smaller than the in-vacuum mass at the saturation density.

No.	1	2	3	4	5	6	7	8	9
$g_{\sigma\eta'}$	2.21	2.51	2.94	2.17	2.44	2.82	2.13	2.38	2.70

\* No width (absorptive process) for eta(958) states

\* Large deformation (compression) for 1s eta(958)

Prog. Theor. Exp. Phys. 2019, 053D02 (22 pages)



**Fig. 8.** Nuclear density profiles of the  $\eta'$  bound systems in  $^{12}\text{C}$  obtained with parameter sets 6, 8, and 9. The dotted line stands for the density distribution of the normal nucleus  $^{12}\text{C}$  without  $\eta'$  calculated with parameter set 8.

- No clear peak. Upper limit is obtained.
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Suzuna Kinutani started some discussions  
on Formation Spectra => Poster presentation

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## **(0) Introduction**

**(1)  $\eta'$ (958) mesic nucleus -- after GSI exp.--**

**(2)  $\bar{p}$  : Hadronic systems from Anti-Protonic atom**

# Hadronic systems from Anti-Protonic atom

## ■ An idea proposed to JPARC

Letter of Intent for J-PARC

### Double Anti-kaon Production in Nuclei by Stopped Anti-proton Annihilation

*dated on 17 / 06 / 2009*

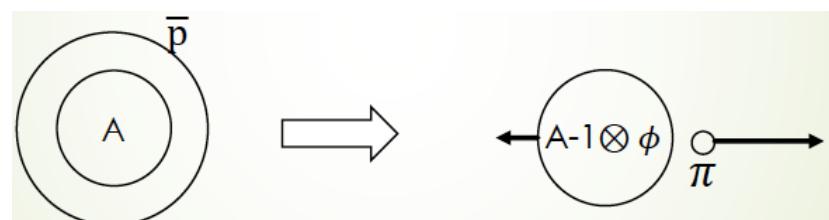
M. Iwasaki<sup>1</sup>, P. Kienle<sup>2,3</sup>, H. Ohnishi<sup>1</sup>, F. Sakuma<sup>1\*</sup>, and J. Zmeskal<sup>2</sup>

## ■ Other channels and/or Theoretical evaluation

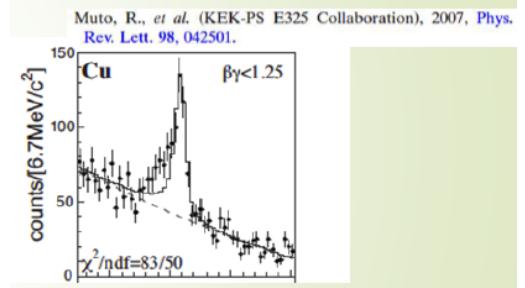
based on the formula for hypernucleus formation by stopped K

A. Matsuyama and K. Yazaki. NPA477 (1988) 673

Ex.)  $\phi$  from  $\bar{p} + N \rightarrow \phi + \pi$  by Mai Hirao and Suzuna Kinutani



cf.)



■ Based on A. Matsuyama and K. Yazaki. NPA477 (1988) 673

Probability to form bound states

$$P \left( (nl)_{\bar{p}} \rightarrow \pi j_N^{-1} (nl)_\phi \right) = \frac{\Gamma \left( (nl)_{\bar{p}} \rightarrow \pi j_N^{-1} (nl)_\phi \right)}{\Gamma_{\text{total}} \left( (nl)_{\bar{p}} \rightarrow X \right)}$$

Numerator

$$\begin{aligned} & \Gamma \left( (nl)_{\bar{p}} \rightarrow \pi j_N^{-1} (nl)_\phi \right) \\ &= \frac{k_\pi \omega_\pi}{4\pi^2} \frac{1}{2l_{\bar{p}} + 1} |t(\bar{p}N \rightarrow \pi\phi)|^2 \int d\Omega_\pi N_{\text{eff}} \end{aligned}$$

$$N_{\text{eff}} = \sum_{m_{\bar{p}} M} \left| \int d\vec{r} \chi_\pi^* (\vec{r}) \left[ \phi_\phi^* (\vec{r}) \bigotimes \psi_N (\vec{r}) \right]_{JM} \varphi_{\bar{p}}^{nl} \right|^2$$

Abs(t)<sup>2</sup> term – from experimental BR of elementary process.

$$B(N^{-1}\phi) = \frac{\sum_{N,\phi,\pi} \Gamma \left( (nl)_{\bar{p}} \rightarrow \pi j_N^{-1} (nl)_\phi \right)}{\Gamma_{\text{total}} \left( (nl)_{\bar{p}} \rightarrow X \right)}$$

## ■ An Estimation by Hirao and Kinutani

BR

\* ASTERIX collabo., PLB267 (1991) 299

$$B(\bar{p} + p \rightarrow \phi + \pi^0) = 40 \times 10^{-5}$$

$$B(\bar{p} + n \rightarrow \phi + \pi^-) = 51 \times 10^{-5}$$

For 12C,  $(1s)_{\bar{p}} \rightarrow \pi s_{1/2}^{-1}(0s)_\phi$  case,

$$\Rightarrow P\left((1s)_{\bar{p}} \rightarrow \pi s_{1/2}^{-1}(0s)_\phi\right) = 5.7 \times 10^{-11}$$

so small !!

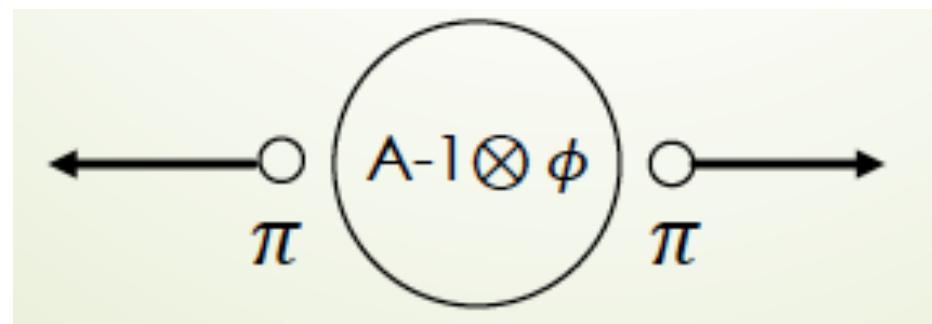
\* 2 origins for small P

(1) Small BR,

(2) Large momentum transfer ( $\sim 1\text{GeV}$  in this case)

Thus,

- Decay modes with Relatively large BR
- Decay modes into 3 quasi-stable mesons with Back to back kinematics like



will be interesting.

# Proton-antiproton annihilation and meson spectroscopy with the Crystal Barrel

Claude Amsler

Reviews of Modern Physics, Vol. 70, No. 4, October 1998

TABLE IX. Branching ratios for  $\bar{p}p$  annihilation at rest into three narrow mesons. Mesons in parentheses were not detected.

Channel	Final state	<i>B</i>				Reference
$\pi^0\pi^0\pi^0$	$6\gamma$	6.2	$\pm$	1.0	$10^{-3}$	Amsler <i>et al.</i> (1995f)‡
$\pi^+\pi^-\pi^0$	$\pi^+\pi^-(\pi^0)$	6.9	$\pm$	0.4	$10^{-2}$	Foster <i>et al.</i> (1968b)
$\pi^0\eta\eta$	$6\gamma$	2.0	$\pm$	0.4	$10^{-3}$	Amsler <i>et al.</i> (1995e)‡
$\pi^0\pi^0\omega$	$7\gamma$	2.00	$\pm$	0.21	$10^{-2}$	Amsler <i>et al.</i> (1993a)‡
	$\pi^+\pi^-6\gamma$	2.57	$\pm$	0.17	$10^{-2}$	Amsler <i>et al.</i> (1994d)‡
$\pi^+\pi^-\omega$	$2\pi^+2\pi^-(\pi^0)$	6.6	$\pm$	0.6	$10^{-2}$	Bizzarri <i>et al.</i> (1969)
$\omega\eta\pi^0$	$7\gamma$	6.8	$\pm$	0.5	$10^{-3}$	Amsler <i>et al.</i> (1994c)‡
$\pi^0\pi^0\eta$	$6\gamma$	6.7	$\pm$	1.2	$10^{-3}$	Amsler <i>et al.</i> (1994b)‡
	$\pi^+\pi^-6\gamma$	6.50	$\pm$	0.72	$10^{-3}$	Amsler <i>et al.</i> (1994d)‡
$\pi^+\pi^-\eta$	$\pi^+\pi^-2\gamma$	1.63	$\pm$	0.12	$10^{-2}$	Abele <i>et al.</i> (1997a)‡
	$\pi^+\pi^-6\gamma$	1.33	$\pm$	0.16	$10^{-2}$	Amsler <i>et al.</i> (1994d)‡
	$2\pi^+2\pi^-(\pi^0)$	1.38	$\pm$	0.17	$10^{-2}$	Espigat <i>et al.</i> (1972)
	$2\pi^+2\pi^-(\pi^0)$	1.51	$\pm$	0.17 0.21	$10^{-2}$	Foster <i>et al.</i> (1968a)
$\pi^0\pi^0\eta'$	$10\gamma$	3.2	$\pm$	0.5	$10^{-3}$	Abele <i>et al.</i> (1997e)‡
	$6\gamma$	3.7	$\pm$	0.8	$10^{-3}$	Abele <i>et al.</i> (1997e)‡
$\pi^+\pi^-\eta'$	$\pi^+\pi^-6\gamma$	7.5	$\pm$	2.0	$10^{-3}$	Urner (1995)‡
	$3\pi^+3\pi^-(\pi^0)$	2.8	$\pm$	0.9	$10^{-3}$	Foster <i>et al.</i> (1968a)
$\pi^0\eta\eta'$	$6\gamma$	2.3	$\pm$	0.5	$10^{-4}$	Amsler <i>et al.</i> (1994f)‡
$\pi^0\pi^0\phi$	$8\gamma(K_L)$	9.7	$\pm$	2.6	$10^{-5}$	Abele <i>et al.</i> (1997c)‡
$\pi^+\pi^-\phi$	$2\pi^+2\pi^-(K_L)$	4.6	$\pm$	0.9	$10^{-4}$	Bizzarri <i>et al.</i> (1969)
$\pi^0K_SK_L$	$3\pi^0(K_L)$	6.7	$\pm$	0.7	$10^{-4}$	Amsler <i>et al.</i> (1993d) a‡
$\pi^0K_SK_S$	$2\pi^+2\pi^-(\pi^0)$	7.5	$\pm$	0.3	$10^{-4}$	b
$\pi^\pm K^\mp K_S$	$\pi^+\pi^-\pi^\pm K^\mp$	2.73	$\pm$	0.10	$10^{-3}$	b
$\pi^\pm K^\mp K_L$	$\pi^\pm K^\mp(K_L)$	2.91	$\pm$	0.34	$10^{-3}$	Abele <i>et al.</i> (1998d)‡
$\omega K_SK_S$	$3\pi^+3\pi^-(\pi^0)$	1.17	$\pm$	0.07	$10^{-3}$	Bizzarri <i>et al.</i> (1971)
$\omega K^+K^-$	$K^+K^-\pi^+\pi^-(\pi^0)$	2.30	$\pm$	0.13	$10^{-3}$	Bizzarri <i>et al.</i> (1971)

<sup>a</sup>Using  $B(\pi^0\phi)$  from Table VI and Eq. (47).

<sup>b</sup>Average between Armenteros *et al.* (1965) and Barash *et al.* (1965).

<sup>‡</sup>Crystal Barrel experiment.

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【2】  $\bar{p}$  : Hadronic systems from Anti-Protonic atom

【3】  $\eta$  mesic nucleus

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- \* Constraining the optical potential in the search for  $\eta$ -mesic  $^4\text{He}$

M. Skurzok<sup>a,\*</sup>, P. Moskal<sup>a</sup>, N.G. Kelkar<sup>b</sup>, S. Hirenzaki<sup>c</sup>, H. Nagahiro<sup>c,d</sup>, N. Ikeno<sup>e</sup>

Physics Letters B 782 (2018) 6–12

=>  $dd \rightarrow ^3\text{He}N\pi$  reaction (**strong absorption**) data

- \* Non-mesonic decay of the  $\eta$ -mesic  $^3\text{He}$  via

$pd \rightarrow (^3\text{He}-\eta)_{\text{bound}} \rightarrow ^3\text{He}2\gamma(6\gamma)$  reaction

M. Skurzok<sup>a</sup>, S. Hirenzaki<sup>b</sup>, S. Kinutani<sup>b</sup>, H. Konishi<sup>b</sup>, P. Moskal<sup>a</sup>, H. Nagahiro<sup>b,c</sup>, O. Rundel<sup>a</sup>

=>  $\eta$  decay into photons. Not strong. **Not absorption but decay.**

## \* $\eta$ momentum distribution

=> photon momenta simulation

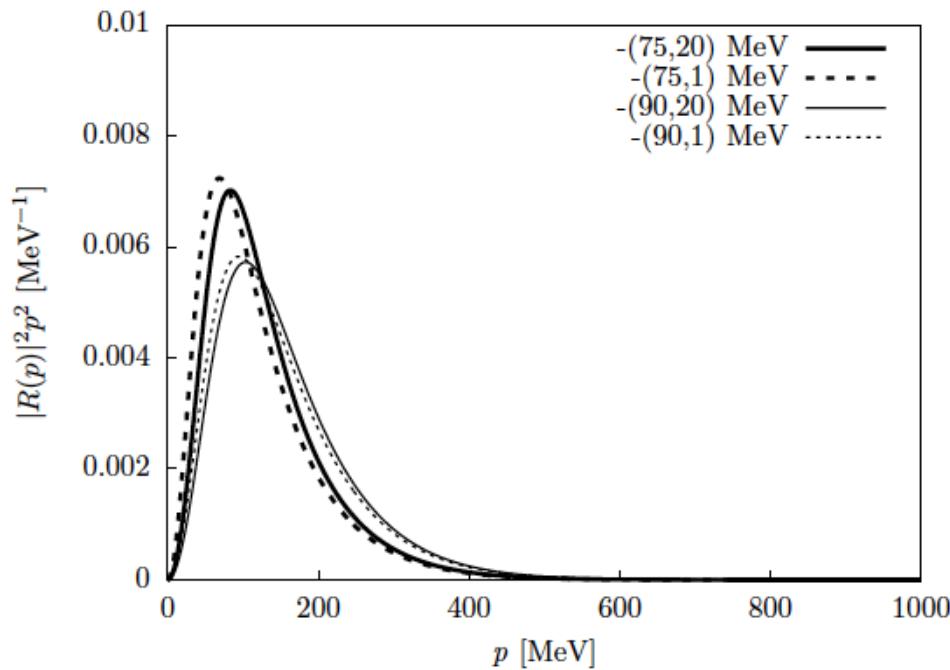


Figure 2: Fermi momentum distribution of the  $\eta$  meson in  ${}^3\text{He}-\eta$  bound system estimated for  $(V_0, W_0) = -(75, 20) \text{ MeV}$  (thick solid line),  $(V_0, W_0) = -(75, 1) \text{ MeV}$  (thick dotted line),  $(V_0, W_0) = -(90, 20) \text{ MeV}$  (thin solid line), and  $(V_0, W_0) = -(90, 1) \text{ MeV}$  (thin dotted line). The distributions are normalized to be 1 in the whole momentum range.

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**【4】 Summary**

## ■ Summary

- \*  $\eta'$ (958) : Deformation (compression) effects may change the meanings of the data (upper limit).
- \*  $\bar{p}$  : Hadronic systems production from Anti-Protonic atom may be interesting. Not only stopped but also in-flight.

J. Yamagata-Sekihara, D. Cabrera, M. J. Vicente-Vacas, S. Hirenzaki.  
Prog.Theor.Phys. 124 (2010) 147-162 (Formation by in-flight p-bar)

- \*  $\eta$  : Let's see what happen experimentally.  
(Magda's talk.)

Suppose there is nuclear deformation, nuclear deexcitation should happen simultaneously. Something could be seen.